



## FHI Seminar



# Semiconductors in presence of gases: The relationship between charge transfer and charge transport

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Berlin 01-12-2017

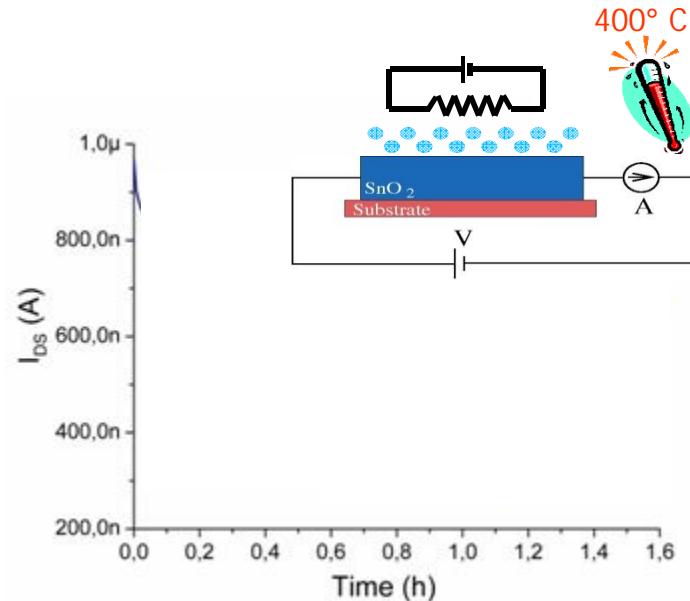
# Outline

- Motivation
- Classic semiconductor theory
  - Ideal „Bulk“
- Semiconductor theory at the bulk edge
  - Semiconductor surface-gas interactions
- One example of device operation
  - Electro-adsorptive effect
  - Thin film transistor
  - Drift of vacancies (dopants)
- Remarks

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# Motivation



- Change in the electron transport in semiconductors in presence of gases (Thin films)
  - Resistance is a macro effect, however it is affected by atomic scale interactions
  - The interaction at the solid-gas interface influences the bulk properties. Is the other way around also true?

- Why are important the semiconductors in catalysis?
  - In most case, metals are **enclosed** in a **semiconductor** coat.
  - Reactions take place at the surface, on the semiconductor.
- The electron Theory of Catalysis on Semiconductors\*.
  - Heterogeneous **catalytic process** are based on **electronic** mechanisms.
  - Elucidate the relationship between catalytic process and electronic properties of semiconductors.
- Theory limits
  - Fully applicable to dielectrics.
  - **Cannot be applied directly to metals.**
    - ⌘ Based on the “many-electron” approach.

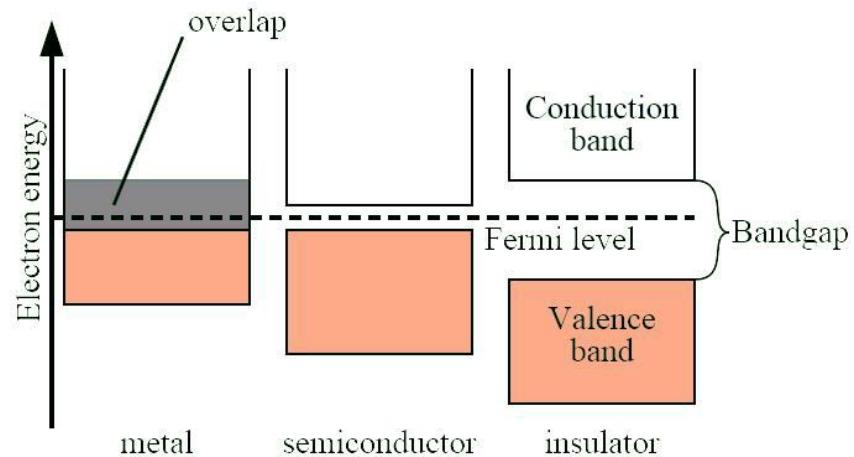
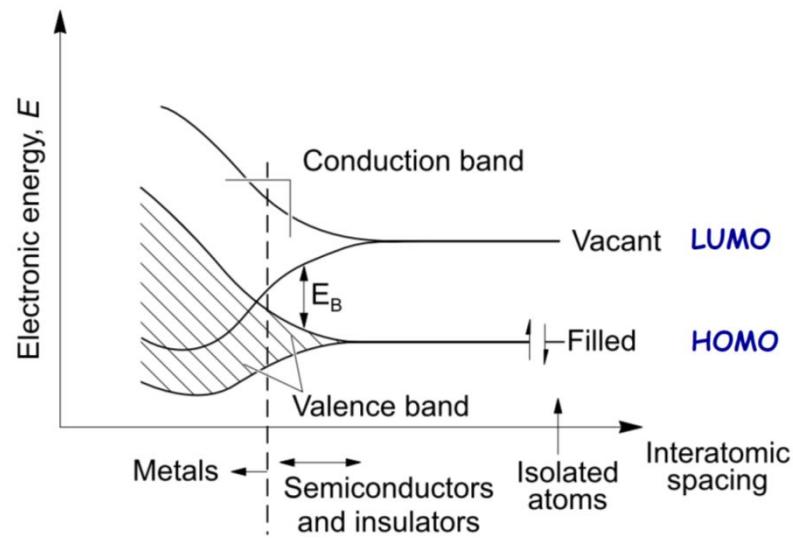
\*F. F. Wolkenstein, W. B. Sandomirski, Dokl. Acad. Nauk. SSSR, 118, pp. 980-982, 1958.

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# Semiconductors

- Is a material with an electrical conductivity between metal and insulator
  - Conductivity can be modified by different parameters:
    - ⌘ Temperature
    - ⌘ Electric field
  - Two types of carriers: electrons and holes



- Weakly bound valence electrons interact with positively charged atomic cores

- Schrödinger equations rules the motion of electrons in solids

- ⌘ Charge doesn't change with time

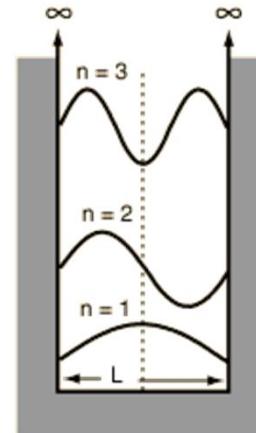
$$E \psi(\mathbf{r}) = (-\hbar^2/2m) \nabla^2 \psi(\mathbf{r}) + U(\mathbf{r})\psi(\mathbf{r})$$

- ⌘ Equipotential  $U(\mathbf{r})=0$

- $$\text{⌘ } E\psi(x) + \frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} = 0 \text{ solution } \Psi(0)=\Psi(L)=0 \quad \psi_n(x) = A \sin\left(\frac{\pi n}{L}x\right)$$

Then we obtain as Eigenvalue:

$$E_n = \frac{\hbar^2}{2m} \left( \frac{\pi n}{L} \right)^2$$

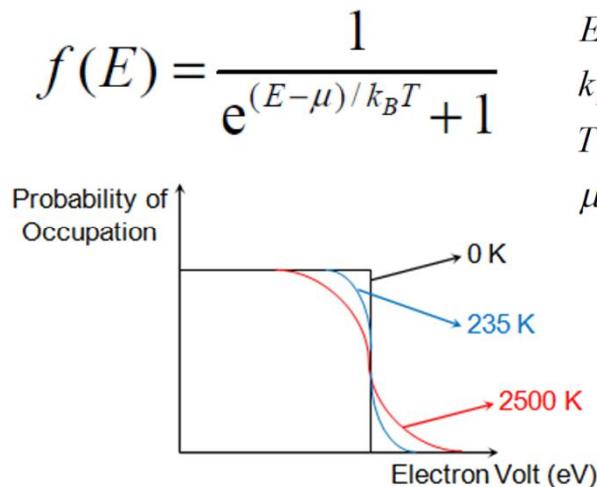


$x = 0$  at left wall of box.

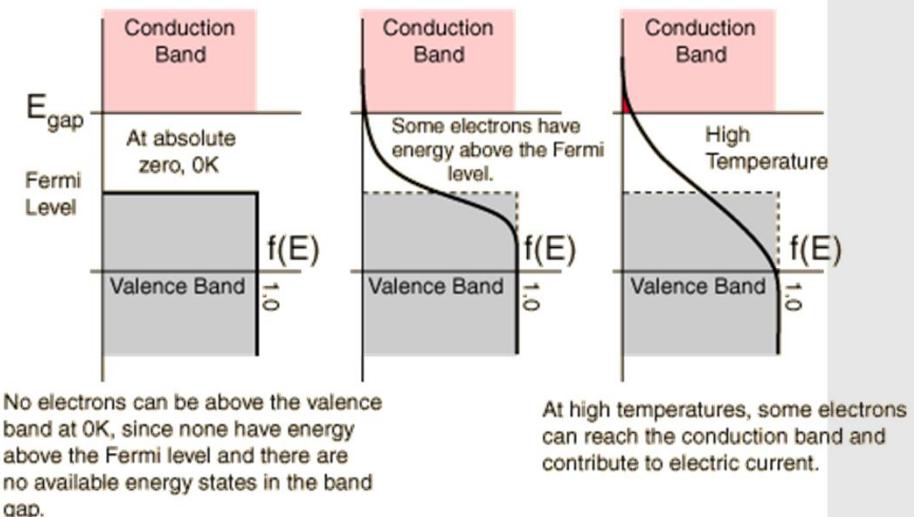
- For  $N$  valence electrons, highest occupied level energy

Fermi energy       $E_F = \frac{\hbar^2}{2m} \left( \frac{\pi N}{2L} \right)^2$

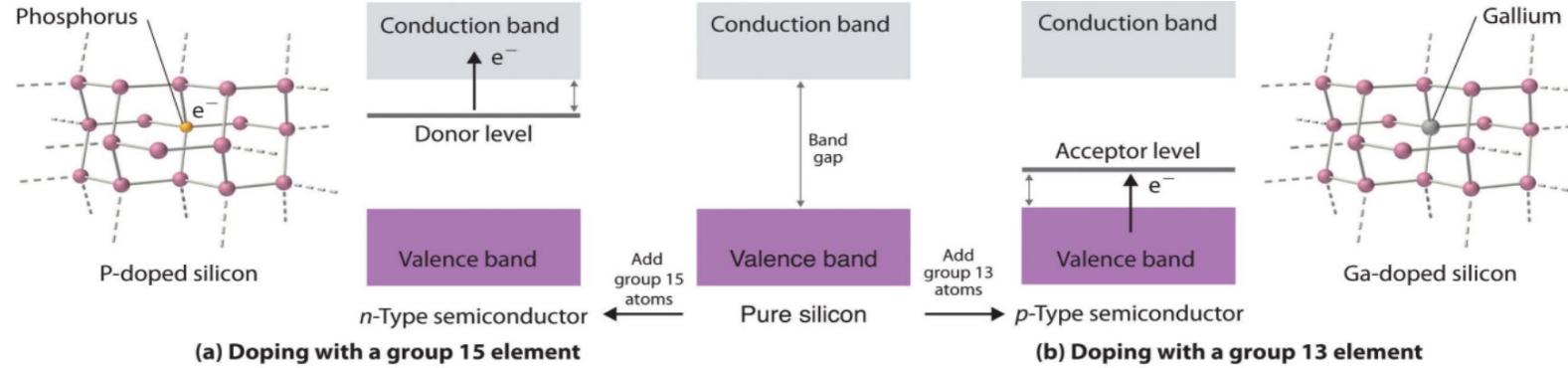
- KE of electrons increases with temperature: distribution function



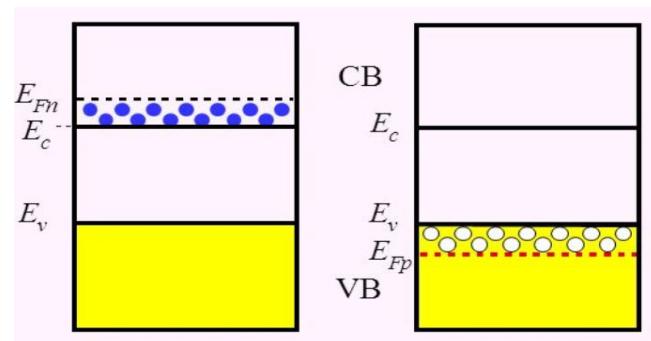
...Energy  
 $k_B$ ...Boltzmann constant  
 T...Temperature  
 $\mu$ ...(Electro-)Chemical potential



# Types of semiconductors



- Intrinsic
- Extrinsic
  - Type n
  - Type p
- Degenerate:
  - $|E_c - E_D| < 3K_B T$
  - $|E_v - E_A| < 3K_B T$



## ■ Intrinsic carriers

$$n_i = \int_{E_C}^{\infty} f(E)g(E)dE$$

$$n_i = N_C e^{-(E_C - E_F)/k_B T}$$

$$n_i = N_V e^{-(E_F - E_V)/k_B T}$$

## ■ Extrinsic carriers

### ➤ Electrons

$$n = \frac{N_D}{e^{(E_C - E_D)/k_B T} + 1} \approx N_D e^{-(E_C - E_D)/k_B T}$$

### ➤ Electron holes

$$p = \frac{N_A}{e^{(E_A - E_V)/k_B T} + 1} \approx N_A e^{-(E_A - E_V)/k_B T}$$

## ■ Resistance: $R = \frac{\rho \ell}{A}$

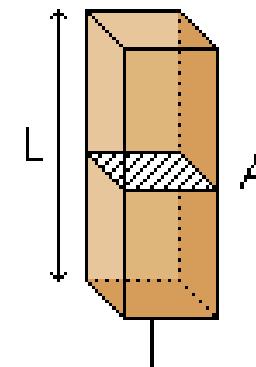
**Thermo-ionic description!!**

### ➤ Resistivity

$$\rho = 1/\sigma$$

### ➤ Conductivity $\sigma = q(\mu_p p - \mu_n n)$

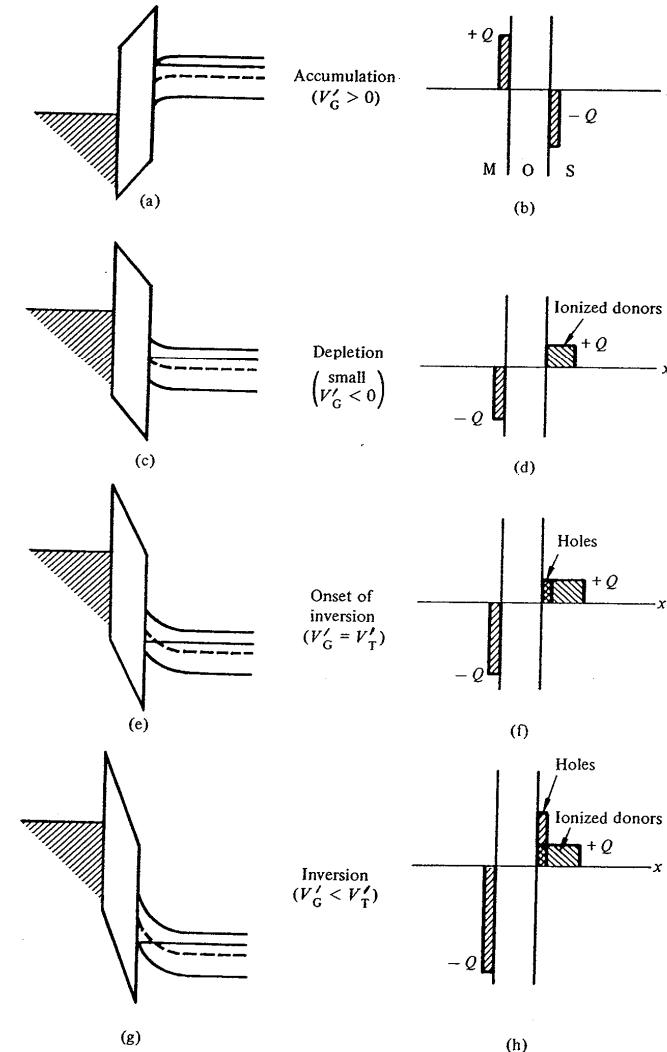
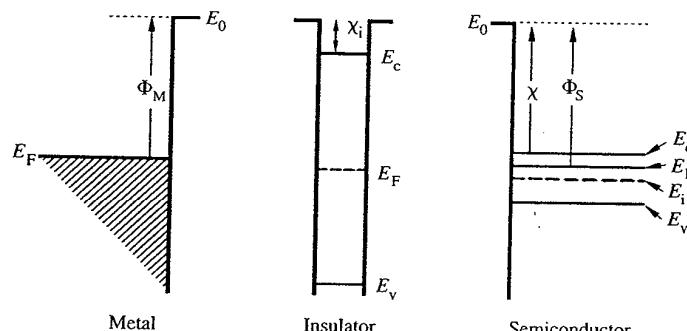
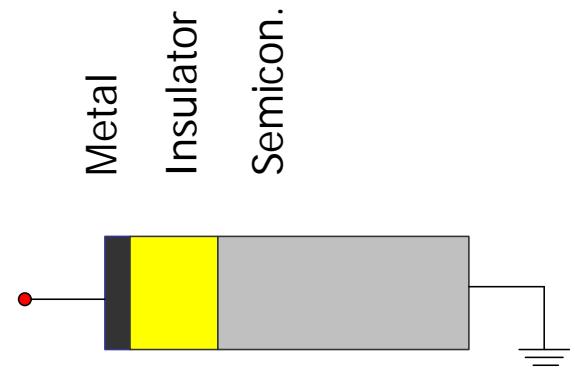
$$\mu_n k_B T = q D_n$$



## ■ Current density

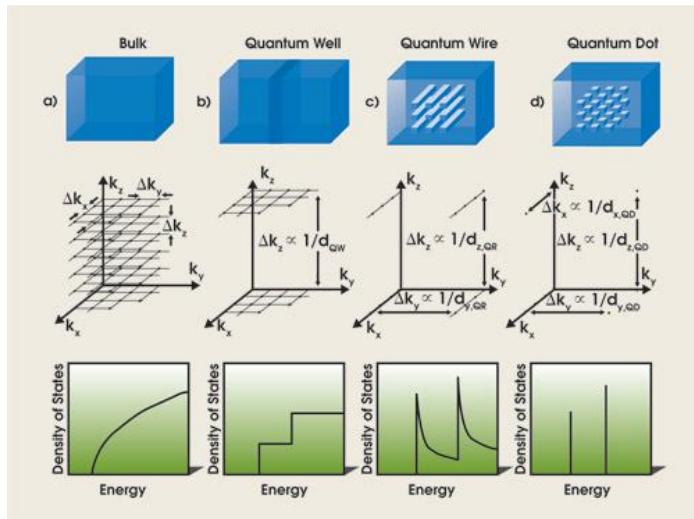
$$J = qn\mu E + qD \frac{dn}{dx}$$

# Band bending (MOS)



# Summary

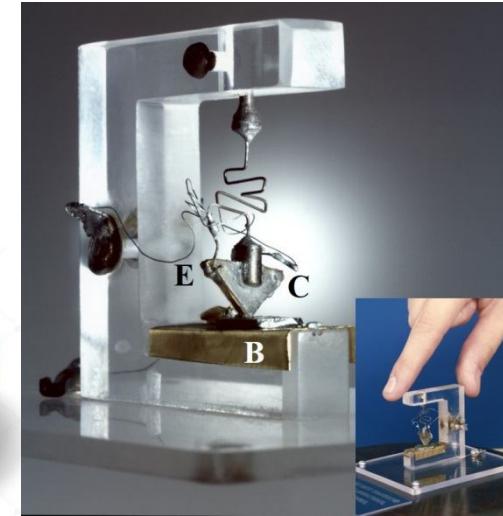
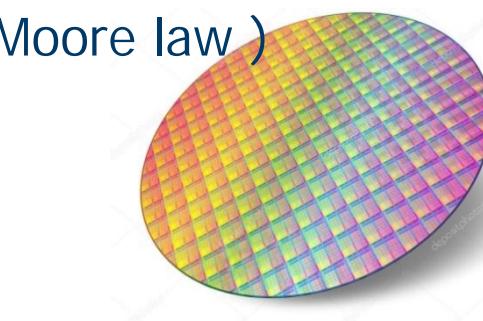
- Semiconductor technology is everywhere
  - Scaling-down Close to the quantum limit
  - $\approx L_D/2$ : quantum confinement:
  - 5n node (end of Moore law )



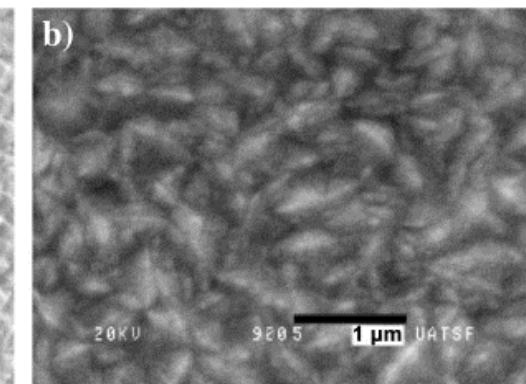
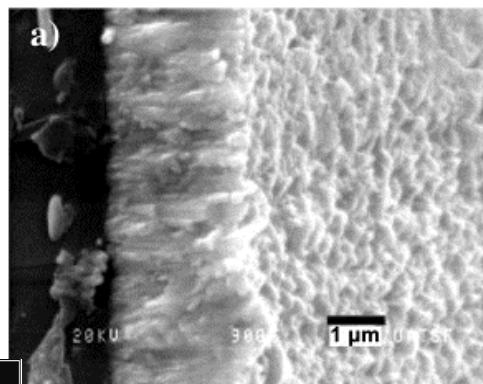
- Defects grain domains
- Surface??



God made the bulk; the surface was invented by the devil.  
—Wolfgang Pauli—



Shockley, Bardeen, Brattain 1948

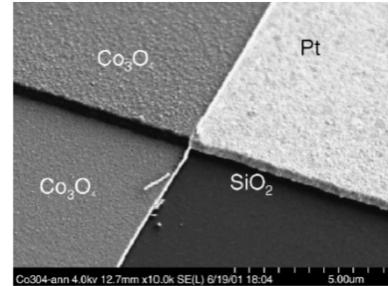


Thin Solid Films Volume 531, 15 March 2013, Pages 172-178

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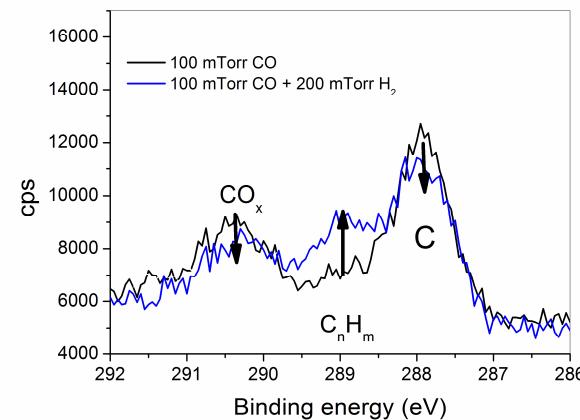
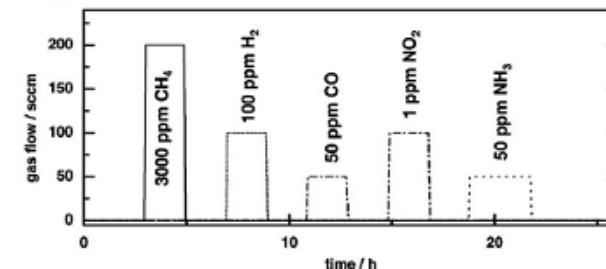
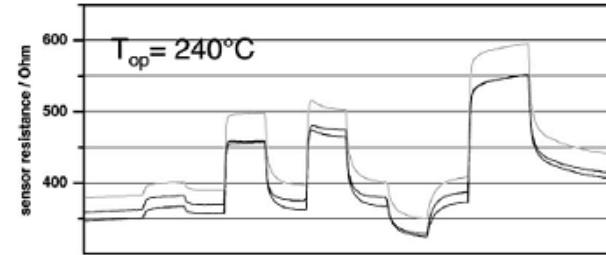
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- Thin film  $\text{Co}_3\text{O}_4$ 
  - Charge transfer information
  - ⌘ Macro-scale

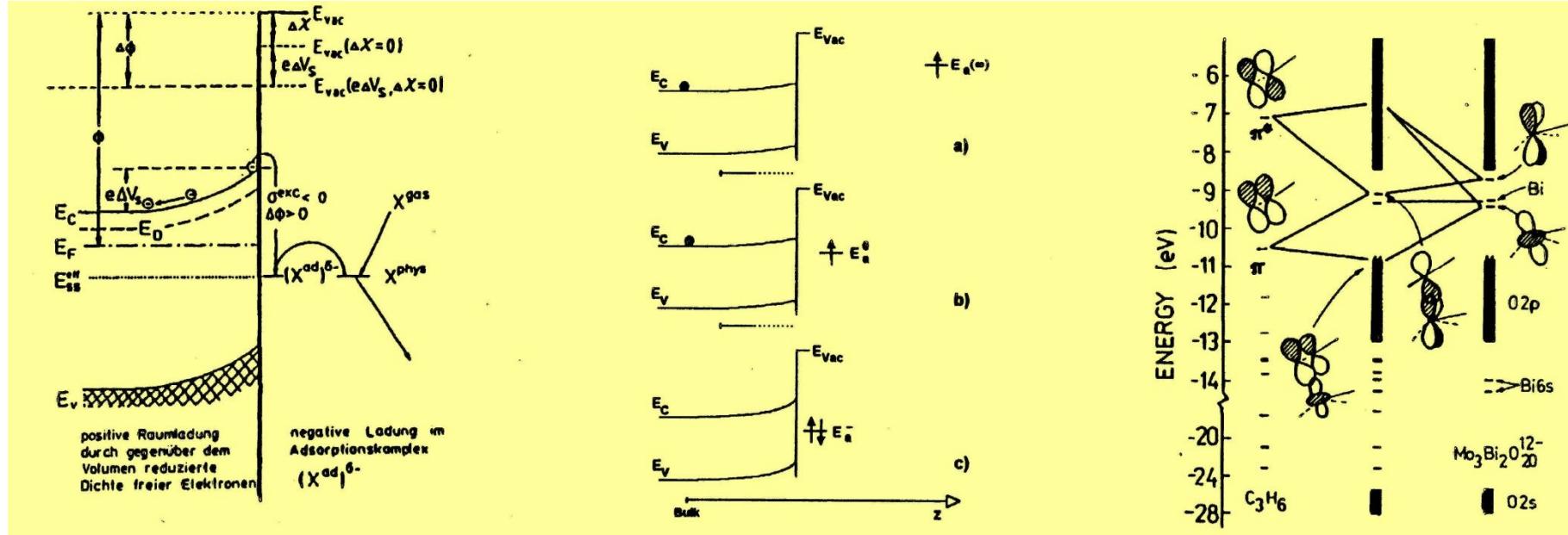


\*J. Wöllenstein, Sensors and Actuators B 93 (2003) 442–448

- Cobalt thin film 240°C
  - Chemical information
  - ⌘ Atomic scale



# Links to Theory

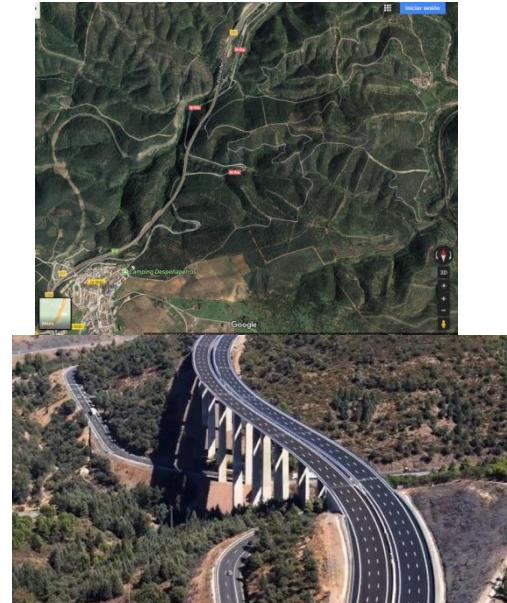
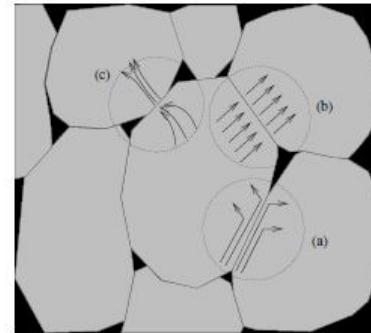


Charge Transfer Model:  
Needs Electron Transfer  
Band Bending  
Thermodynamics

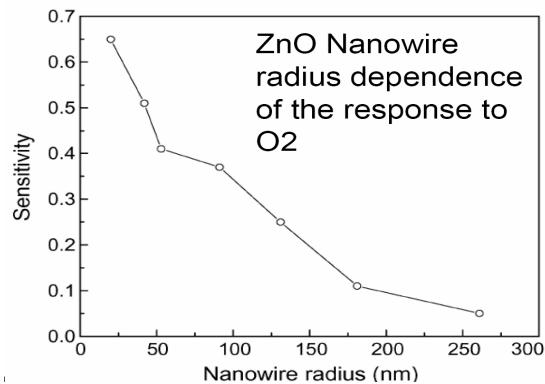
Wolkenstein Model:  
Fully Quantum Mechanical  
Easy Access to  $E_F$ ,  
Semiconductor Device  
Modeling

LCAO:  
Quantum Chemistry  
d, p and s orbitals  
Overlap  
Stereochemistry  
Some Access to  $E_F$

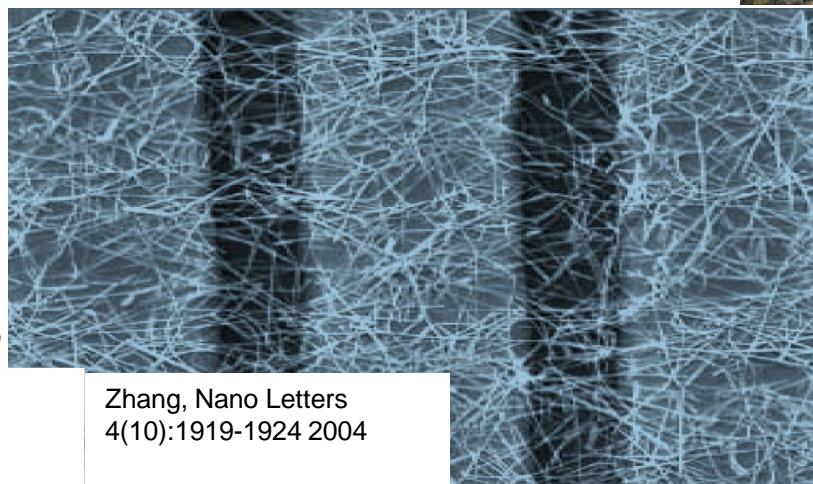
- It can not describe macro effects:
  - Different particle size
  - Or tortuosity



$$E \psi(\mathbf{r}) = (-\hbar^2/2m) \nabla^2 \psi(\mathbf{r}) + U(\mathbf{r})\psi(\mathbf{r})$$



Z.Fan and Jia G. Lu IEEE Trans Nanotech.  
5(4),393,2006



Zhang, Nano Letters  
4(10):1919-1924 2004

# Types of adsorptions

■ Heterogeneous catalysis begins with the act of adsorption.

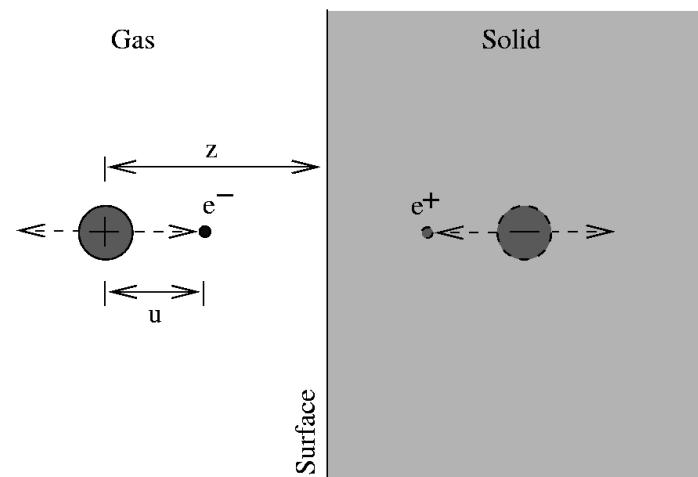
➤ Physisorption

- ⌘ Treated as dipol-dipol interaction (**van der Waals** interaction).
- ⌘ “Long distance interactions”, from 3E-10m to 5E-10m.
- ⌘ Novel-gas adsorbed on a metal.

➤ Chemisorption:

- ⌘ Weak or **neutral**.
- ⌘ Strong or **charged**.

➤ Ionsorption



$$V = \frac{e^-}{4\pi\epsilon_0 2z}$$

- Chemisorbed particle simultaneously both **affinity**:

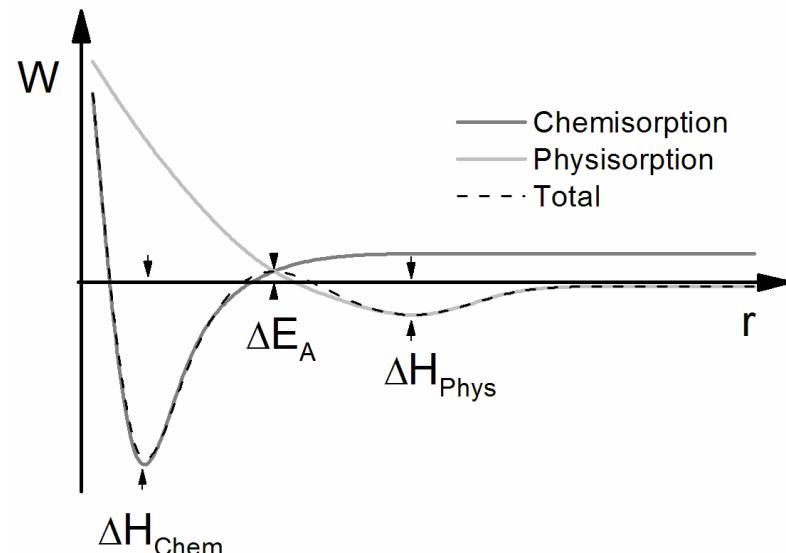
- Acceptor
  - Donor

- Weak chemisorption (CL)

- Particle remains **electrically neutral**.
  - Lattice electrons or holes do not participate in the bond.

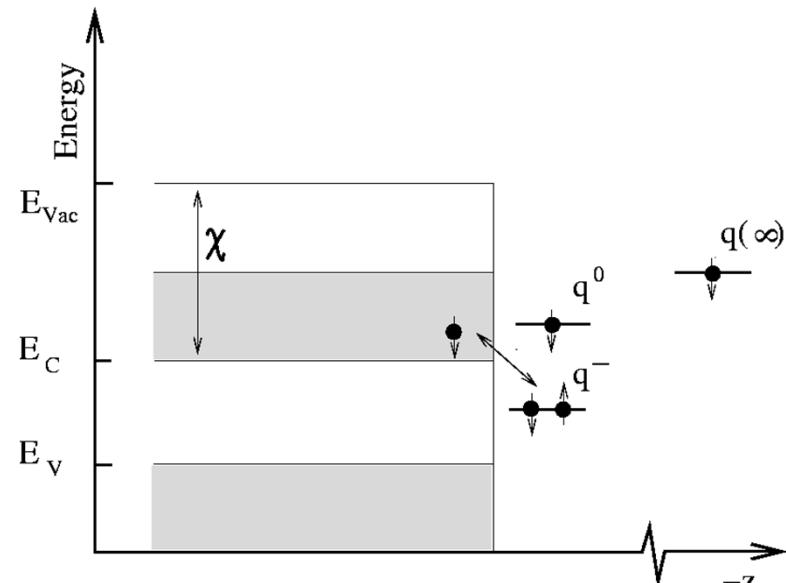
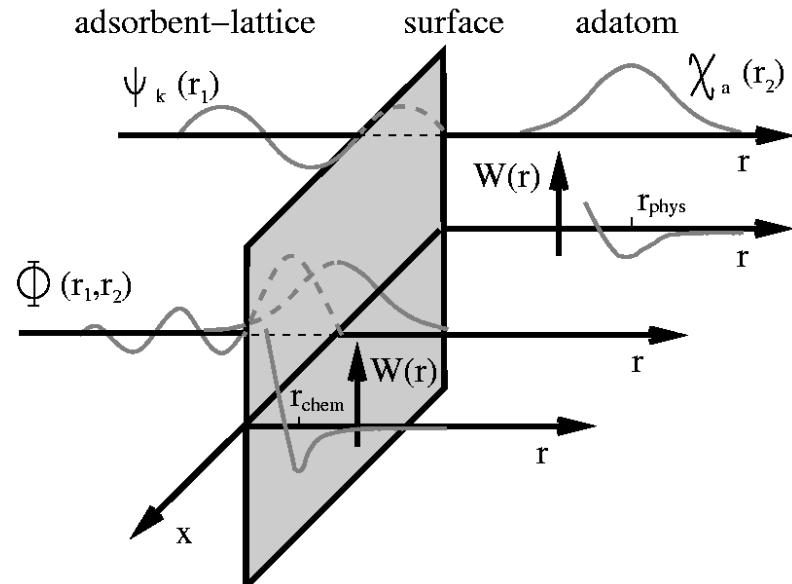
- Strong chemisorption

- Particle **adsorbs** free electrons or holes
    - ⌘n-bond (acceptor bond), CeL.
    - ⌘p-bond (donor bond), CpL.



# Geistlinger theory

- If the electron at the conduction band and adatom has parallel spin, then a repulsive potential is produced. Therefore due to Van der Waals forces, particles are localized at a minimum potential resulting in a physisorbed state. Alternatively, if they have antiparallel spin, the wave functions are overlapped, which leads to chemisorption. Because of this, Wolkenstein postulated a subdivision into two states: Neutral or "weak chemisorbed" and charged or "strong chemisorbed". They are represented as discrete states, meaning that.



**Only strong chemisorbed molecules produce change in the conductivity**

Geistlinger, H. (1993). Electron theory of thin-film gas sensors.  
*Sensors and Actuators B: Chemical*, 17(1), 47-60.

# Fermi statistics

- When still valid the Fermi statitics?

- Non-degenerated semiconductor,  
 $E_V + 3kT < E_F < E_C - 3kT$ .

- Fermi level

- Controls the catalytics properties at the surface.
  - Once electronic equilibrium is established:
    - ⌘ Surface and bulk have the same Fermi level.

$$f^0 = \frac{1}{1 + \frac{1}{2} \exp((E_F + e\Delta V_S - E_a^- + E_a^0)/kT)}; \quad f^- = \frac{1}{1 + 2 \exp((E_a^- - E_a^0 - E_F - e\Delta V_S)/kT)}$$

- Many factors control the Fermi level position.

# Wolkenstien Isotherm

- The Fermi level determines the probability of weak and strong occupancy

$$f^0 + f^- = 1$$

$$\frac{N^-}{N} = f^- = \frac{1}{1 + 2 \cdot \exp\left(\frac{E_t - E_F - e\Delta V_S}{kT}\right)},$$

$$\frac{N^0}{N} = f^0 = 1 - f^- = \frac{1}{1 + \frac{1}{2} \exp\left(\frac{E_F + e\Delta V_S - E_t}{kT}\right)}$$

- The number of chemisorbed particles and adsorption sites are related to the coverage:

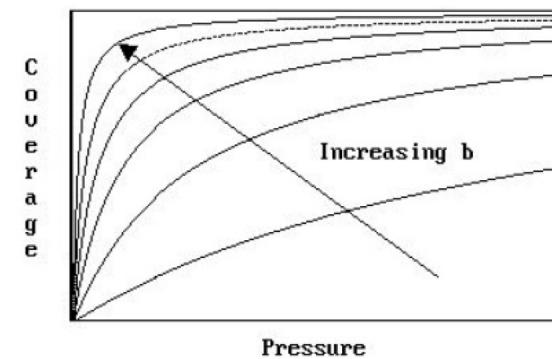
$$\theta = \frac{N}{N^*} = \frac{N^0 + N^-}{N^*} = f^0\theta + f^-\theta = \theta^0 + \theta^-$$

- The Langmuir isotherm is defined as

$$b = \frac{\alpha}{\nu^0} \exp\left(\frac{q^0}{kT}\right)$$

- The Wolkenstein isotherm is expressed as

$$\beta = b \left\{ f^0 \left[ 1 + 2 \cdot \frac{\nu^- f^-}{\nu^0 f^0} \exp\left(\frac{E_t - E_c}{kT}\right) \right] \right\}^{-1}$$

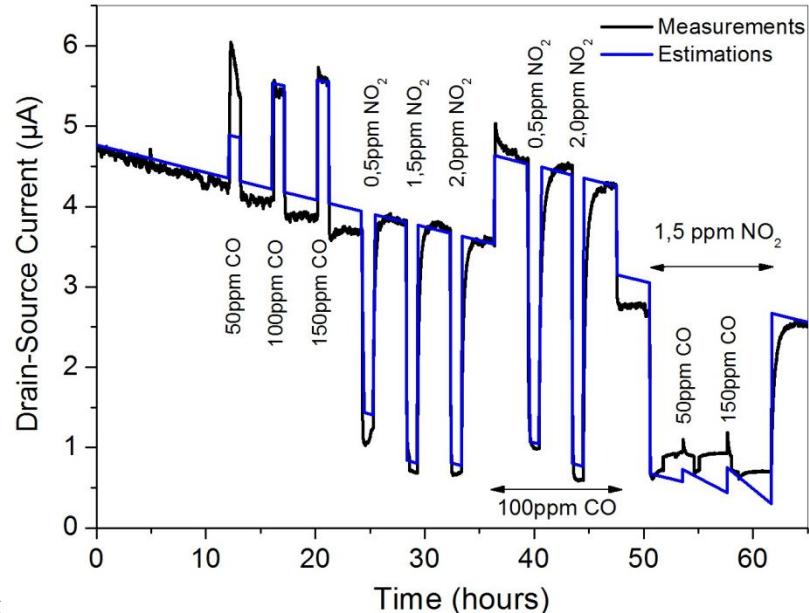


$$\theta(\tilde{p}) = \frac{\beta \tilde{p}}{1 + \beta \tilde{p}}$$

- They don't reach with each other
- Certain interaction
  - Adsorbates compete for the free adsorption centres

$$\theta_j^- = \frac{\beta_j \cdot p_j}{1 + \sum_{i=0}^n \beta_i \cdot p_i} \cdot f^- \quad N^- = \sum_{i=0}^n N_{Surf} \cdot \theta_i^-$$

- Problem to extend this model to catalysis
  - Same problem with Langmuir isotherm



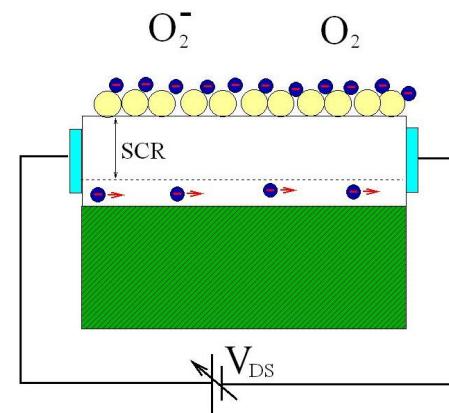
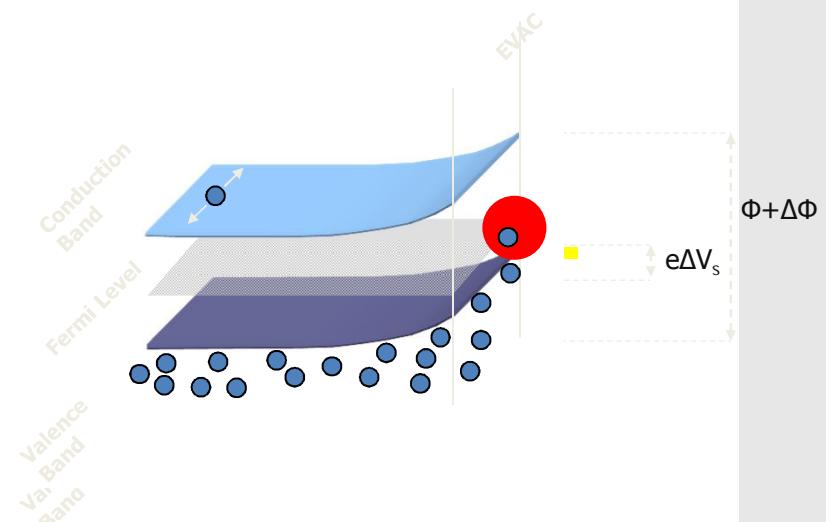
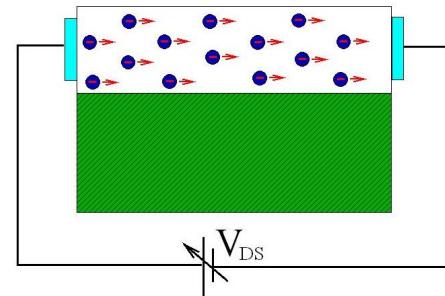
## ■ Band bending

- Surface potential.
- Work function change
- Depleted zone (SCR).
- Electrical neutrality

$$Q_{SCR} = Q_{SS}$$

## ■ Electric properties

- Changes in the bulk conductivity



# Charge transfer

- Charge neutrality  $Q_{SCR} = Q_{SS}$

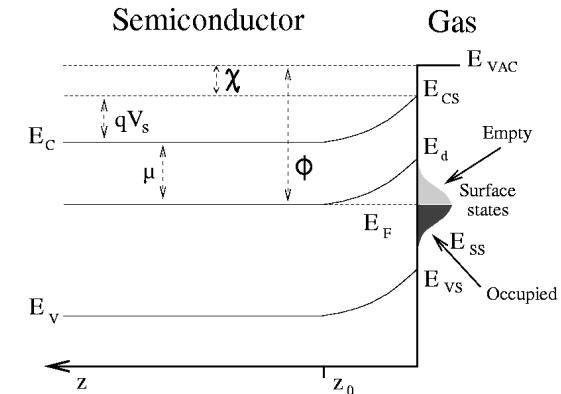
- Charge surface  $Q_{SS} = e^- \cdot \int_{-\infty}^{+\infty} D_{SS}(E) f_{SS}(E) dE$

- Distribution function  $f_{SS}(E, E_F) \rightarrow f_{SS}(E, E_{F,bulk} + e\Delta V_S(\theta^-))$

- Poisson equation  $\frac{d^2V}{dz^2} = -\frac{\rho}{\epsilon_0 \epsilon_r}$

- Surface potential  $V_S = \frac{eN_D z_{SCR}^2}{2\epsilon_0 \epsilon_r} \quad V_S = \frac{e(N^-)^2}{2\epsilon_0 \epsilon_r N_D}$

- Movility thermally activated  $\mu = \mu_0 \cdot \exp(-e^- V_S / kT)$



# Weisz limit

$$(N^- / N) = f^-(E_a^{0/-}, E_F, e\Delta V_S) \quad f^- = \frac{1}{1 + 2 \cdot \exp(\frac{E_t - E_F - e\Delta V_S}{kT})}$$

$$V_S = \frac{e(N^-)^2}{2\varepsilon_0\varepsilon_r N_D}$$

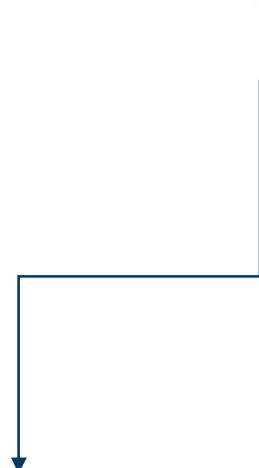
$$Z_{SCR} = (N_s / N_D) = 1E14 \text{ cm}^{-2} / 1E18 \text{ cm}^{-3} = 1 \mu\text{m}$$

$$V_S = 696 \text{ V} !!!!!!!$$

$$V_S = \frac{e N_D z_{SCR}^2}{2\varepsilon_0\varepsilon_r}$$

Self Limitation of Strong  
(=charged) Chemisorption

This is the Weisz Effect



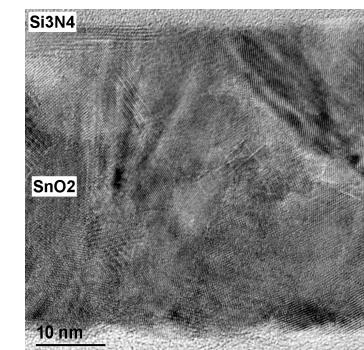
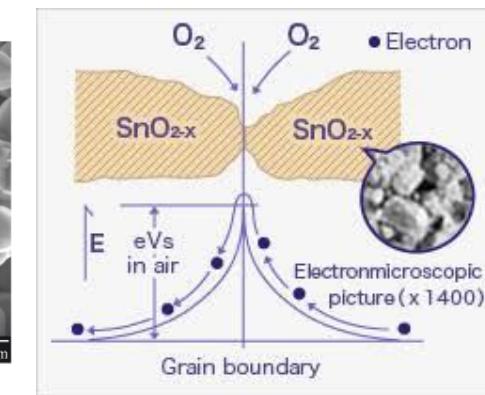
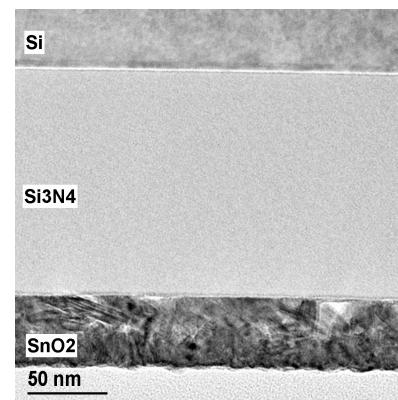
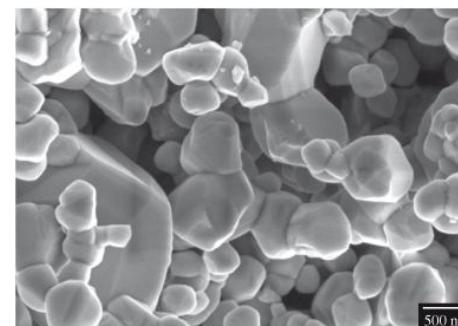
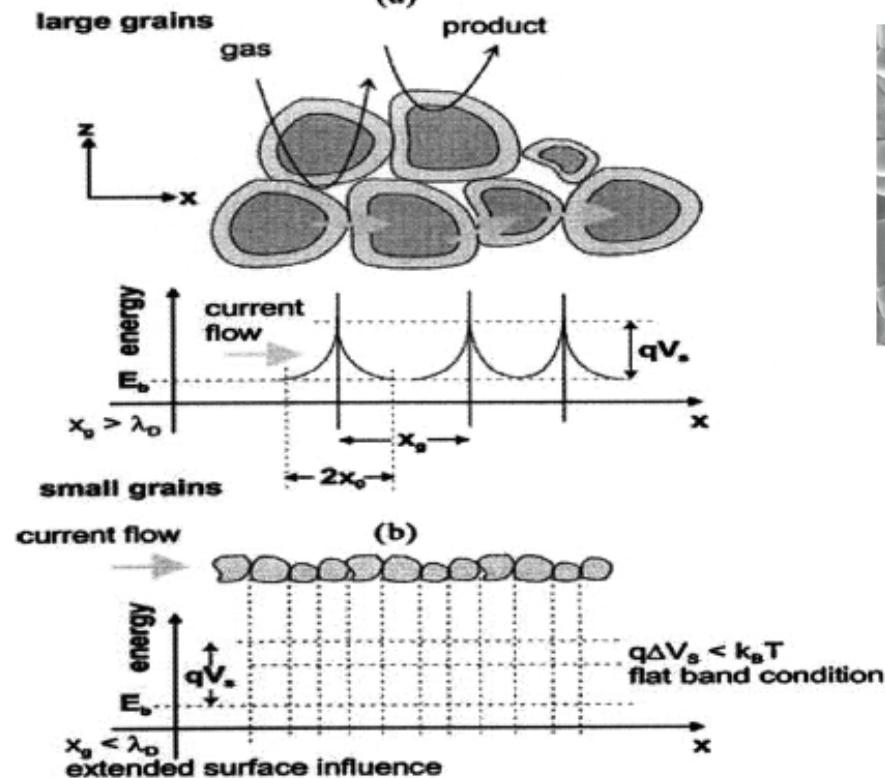
Electrochemistry knows the same Effect: Mott-Schottky

\*Weisz, P. B. (1953). Effects of electronic charge transfer between adsorbate and solid on chemisorption and catalysis. *The Journal of Chemical Physics*, 21(9), 1531-1538.

# Grain size

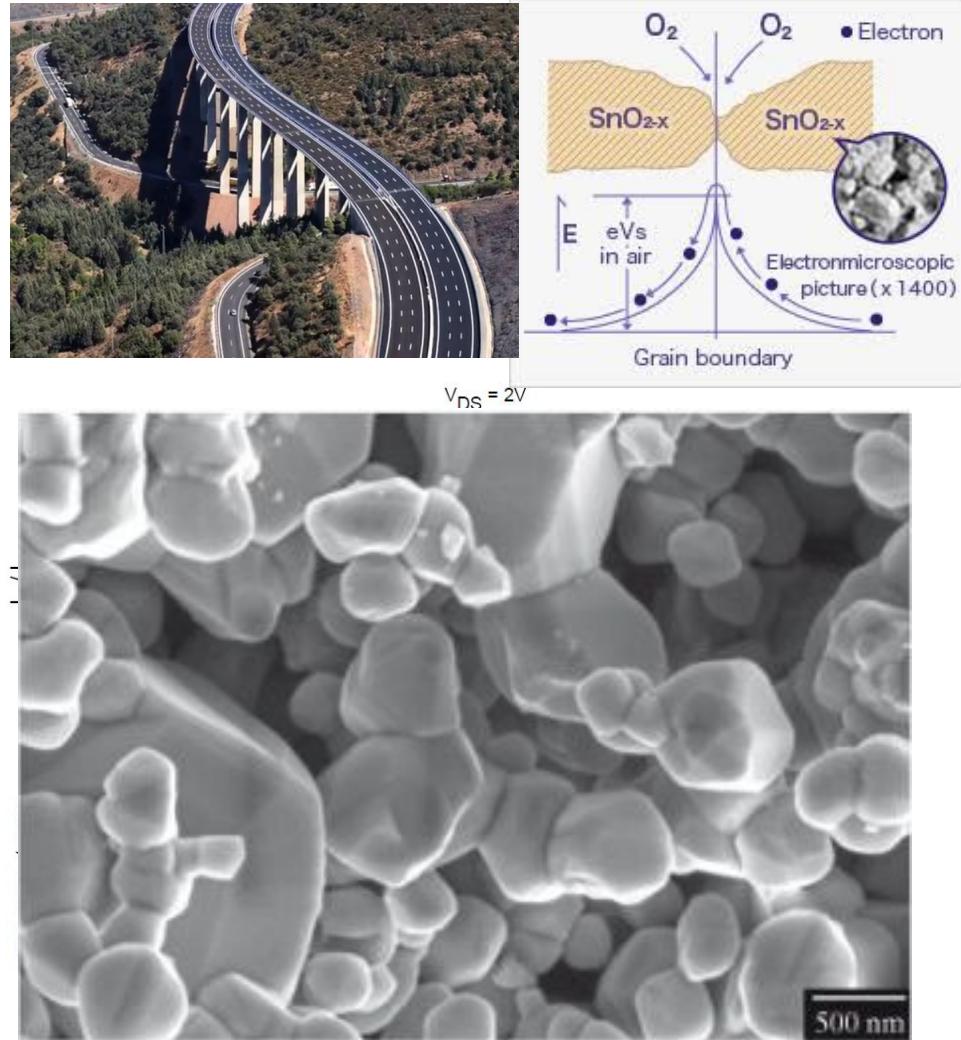
## ■ Flat band limit

- Crystallites  $d_{\text{Grain}} \sim d_{\text{Space Charge}}$
- Fermi staticis:
- Thermoionics vs tunneling



# Summary

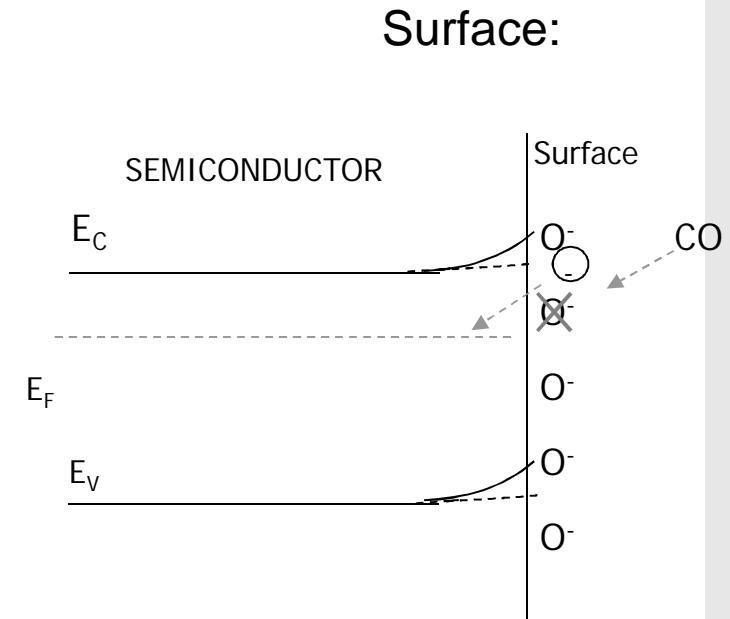
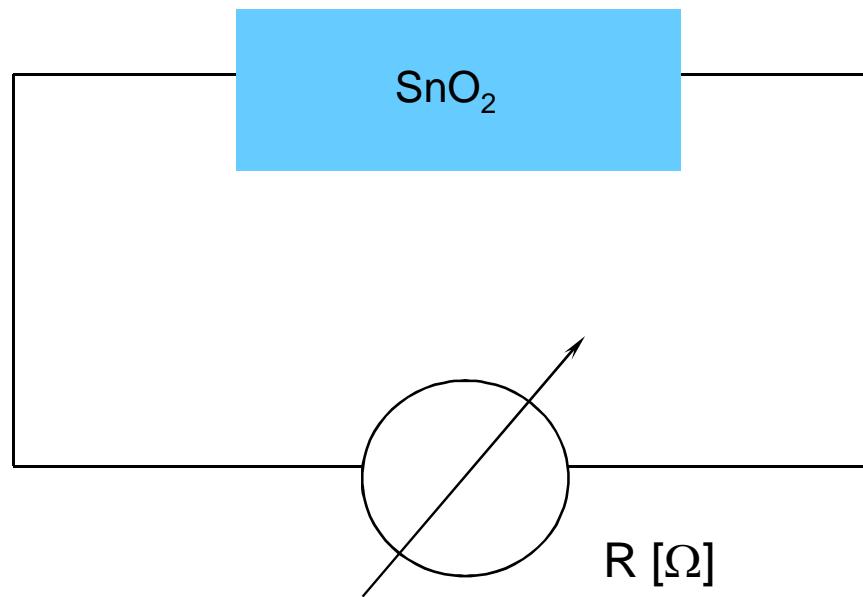
- Charge transport and transfer description
  - Flat band conditions (Fermi-Dirac statistic): Thermo-ionic
  - Wolkenstein isotherm
  - Charge transfer
- Complex description
  - Schrödinger equation
  - Shockley-Read statistics: Tunneling



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# The challenge



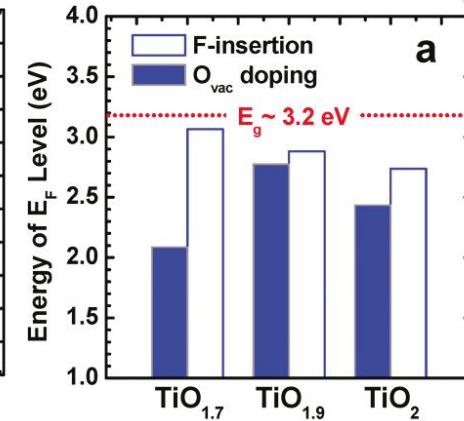
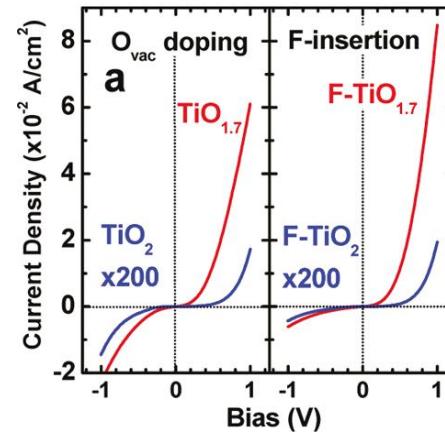
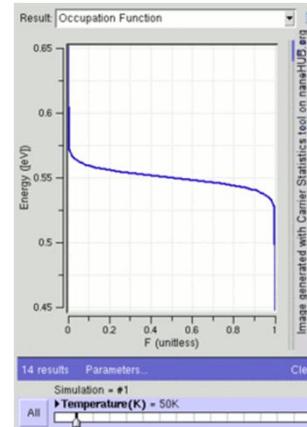
Operation temperature approx: 300°C → CMOS !

## ■ Temperature

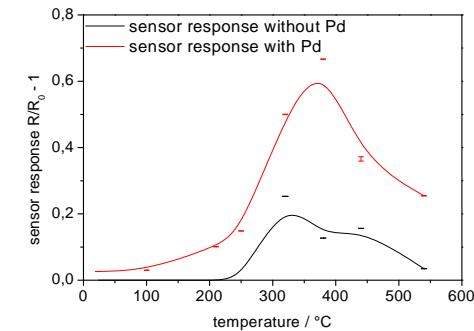
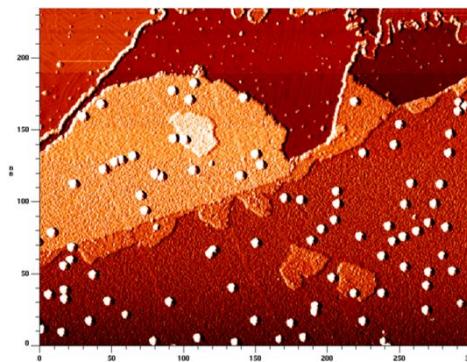
## ■ Doping change Fermi level

## ■ Surface functionalization

## ■ External electric field

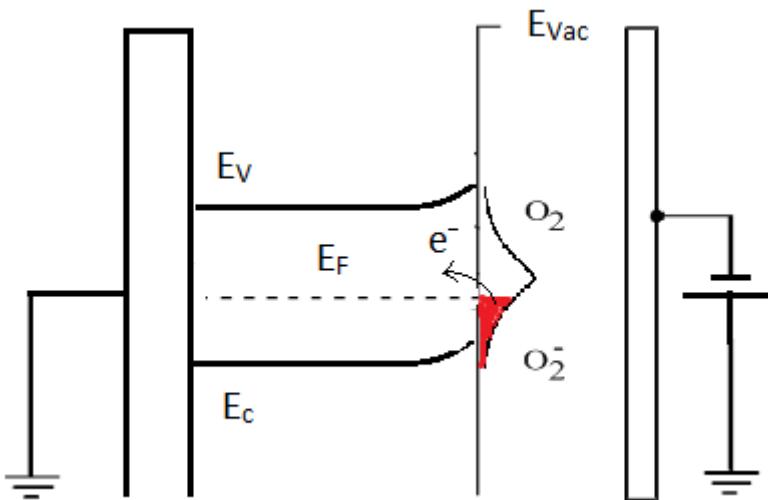


Nano Lett. 2011, 11, 751–756



## ■ EAE:

- “Surface states electrical modulation on the metal-oxide surface by means of an external electric field”



## Field-Effect devices

- JSE. **Lilienfeld**, "Method and apparatus for controlling electric currents", US Patent 1, 745, 175, **1930**.

⑩ It didn't work

⑩ Why...

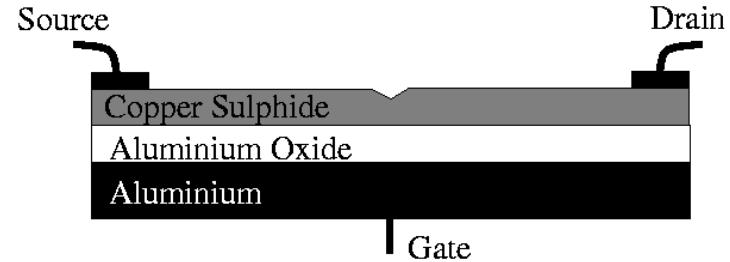
- O. **Heil** control of the layer conductivity by the surface states, **1935**.

- W. **Shockley**, GL. **Pearson**, "Modulation of Conductance of Thin Films of Semi-Conductors by Surface Charges", Phys. Rev. Vol. 74, pp. 232-233, **1948**.

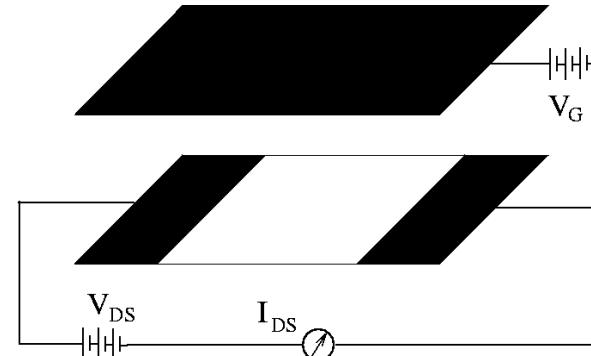
⑩ 10% were mobile.

⑩ Surface states acts like electrical tramps.

- F. **Wolkenstein**, **1958**.



Patente de Lilienfeld (1930)

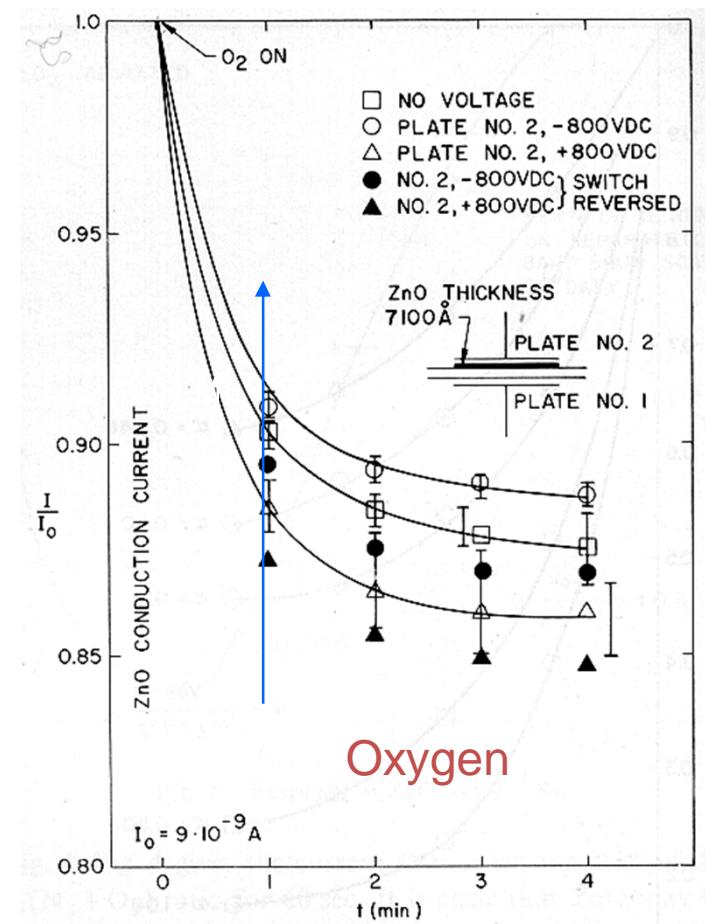
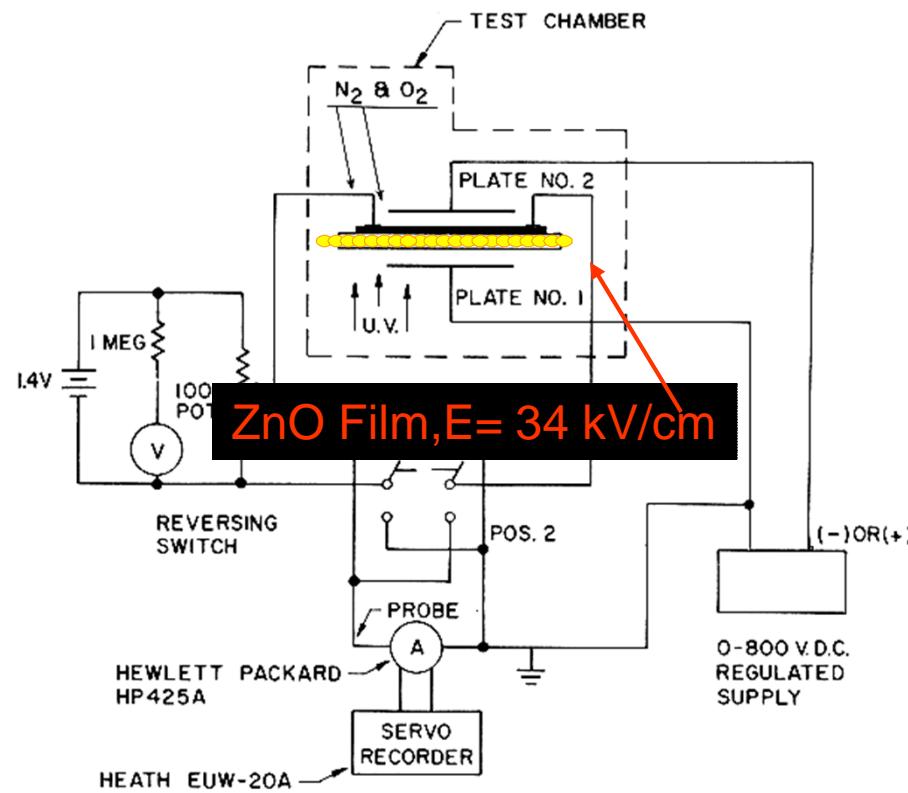


Patente de Heil (1935)

## Short history: EAE

Adsorption of Methanol in Ge: Keier and Mikheeva 1964

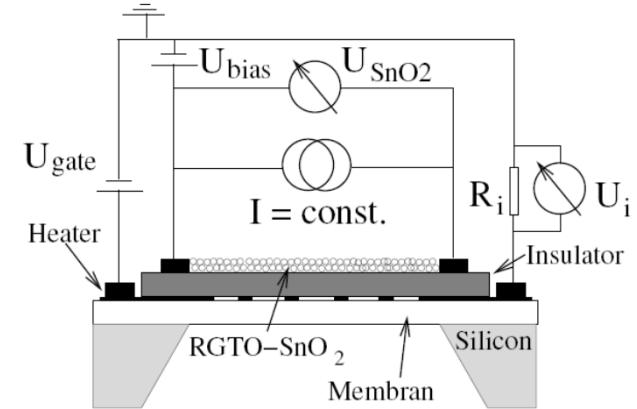
O<sub>2</sub> in ZnO: Hoenig and Lane, Surf. Sci 11, 1968



## ***With micro always high electric field***

SnO<sub>2</sub> Gate Transistor: Popova and Stoyanov

1994



Consistency to the theory: Geistlinger

1994

SnO<sub>2</sub> thin film: Hellmich and Müller

1996

Hellmich and Müller (1996)

SnO<sub>2</sub> multi-electrode Sensor:

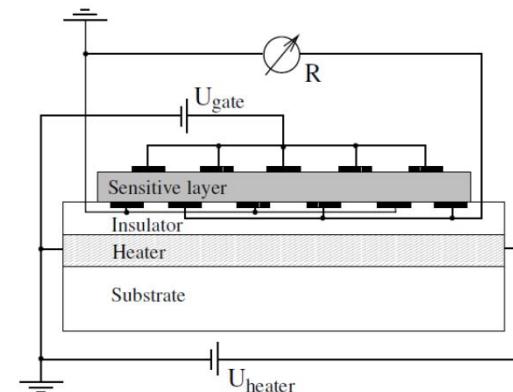
Hausner and Binder

1997

Thin film transistor gas sensor:

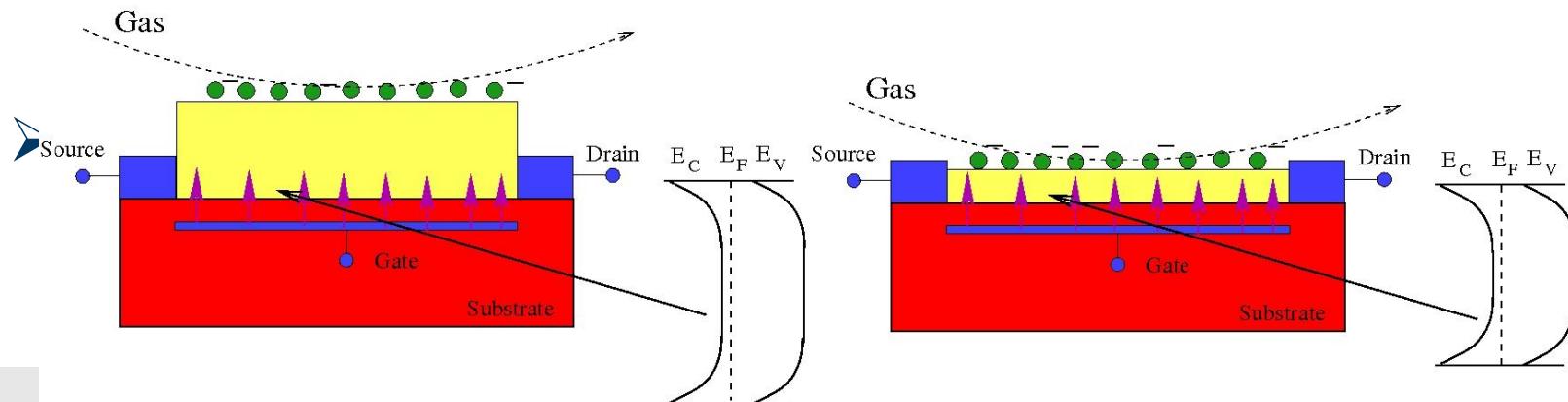
Jaegle and Wöllensteiner

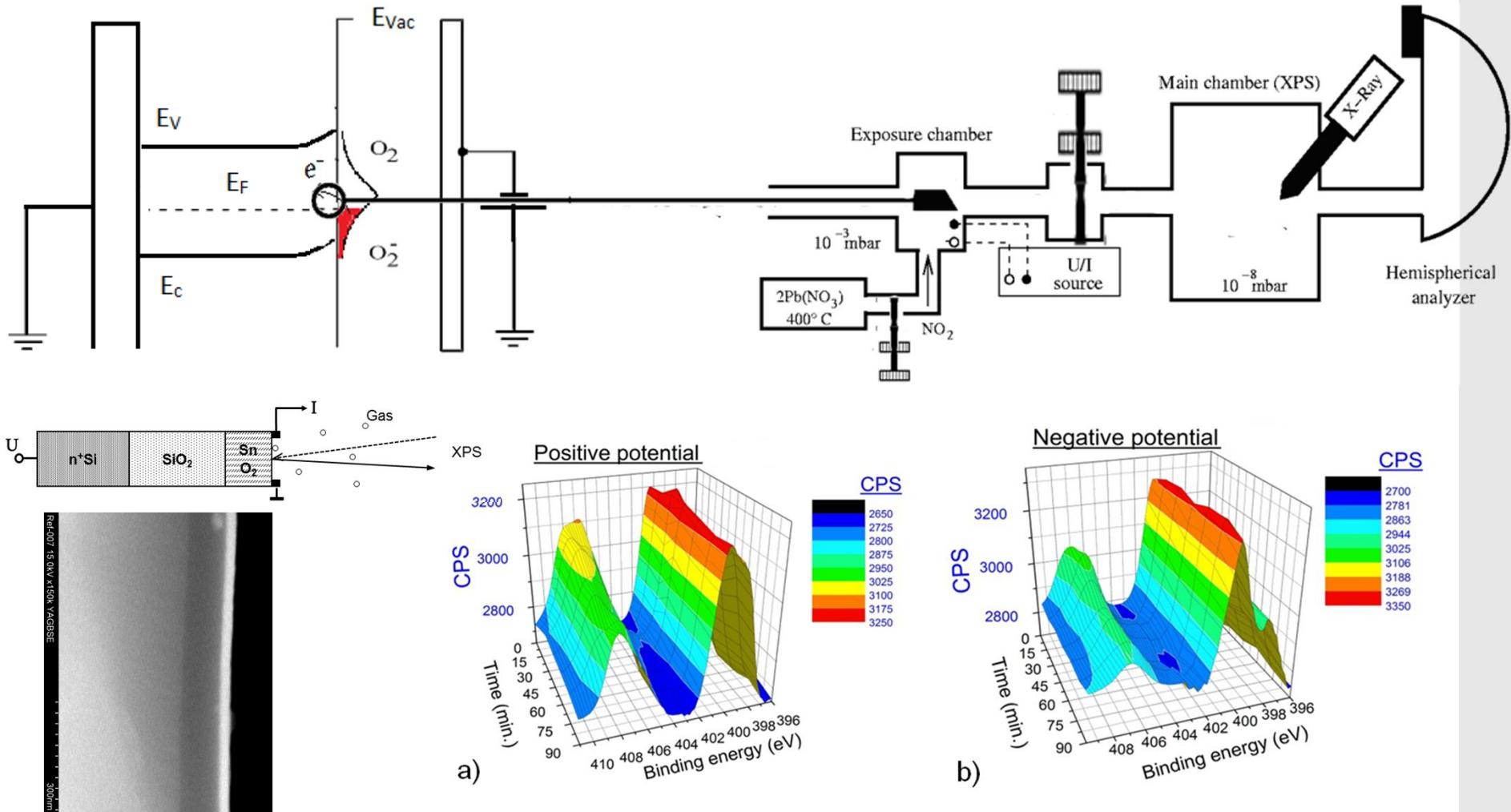
1998



Hausner and Binder (1997)

- Thickness:
  - Must be in the Debye Length range.
  - Energy levels are modulated
  - Reaction on the

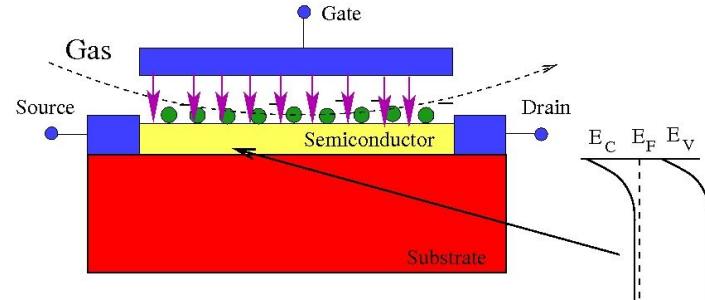




\*ChemPhysChem, 14(11), 2505-2510, 2013

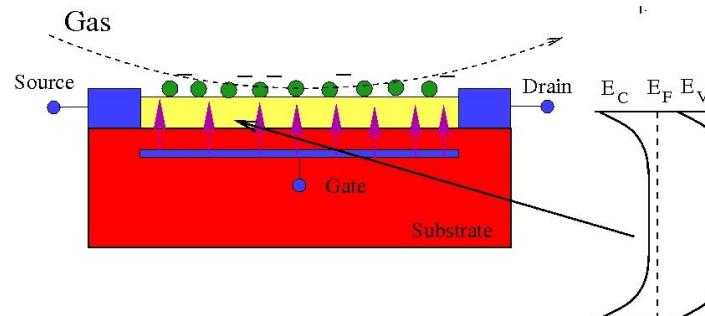
## ■ Suspended Gate

- ⑩ Effective control on the surface states
- ⑩ Very difficult to realize this geometry with standard CMOS processes.



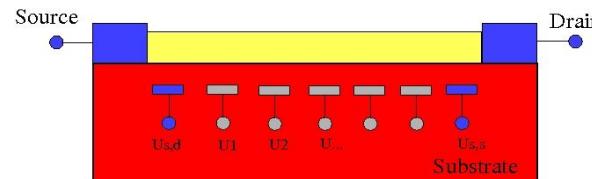
## ■ Thin-Film-Transistor Gas Sensor

- Debye-Lenght:
$$L_D = \sqrt{\frac{\epsilon_0 \cdot \epsilon_r \cdot kT}{(e^-)^2 \cdot N_D}}$$
- Effective control of the Fermi-Level
- Technology: CMOS-Standard processes

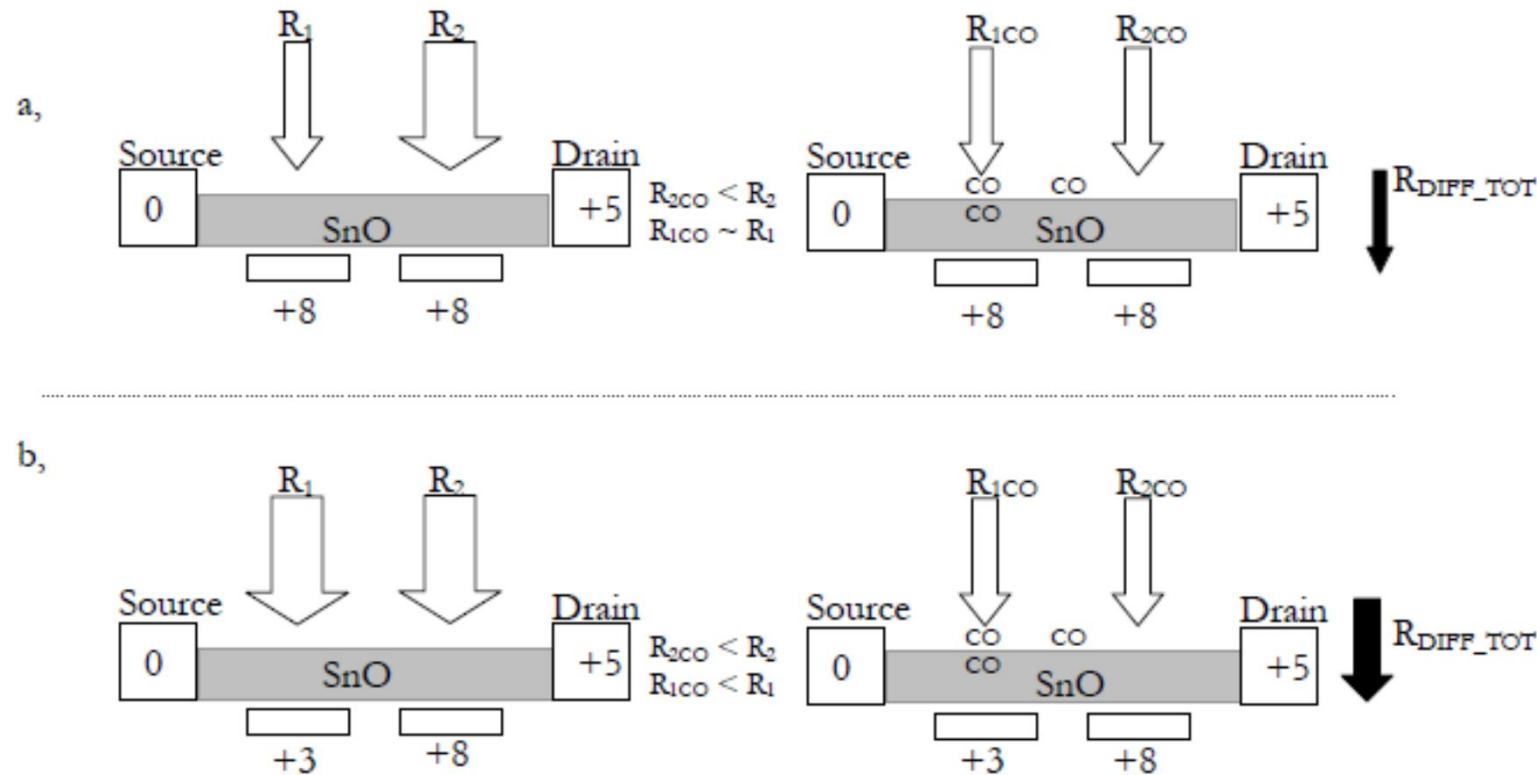


## ■ Multi-Gates

- Homogeneous surface reaction

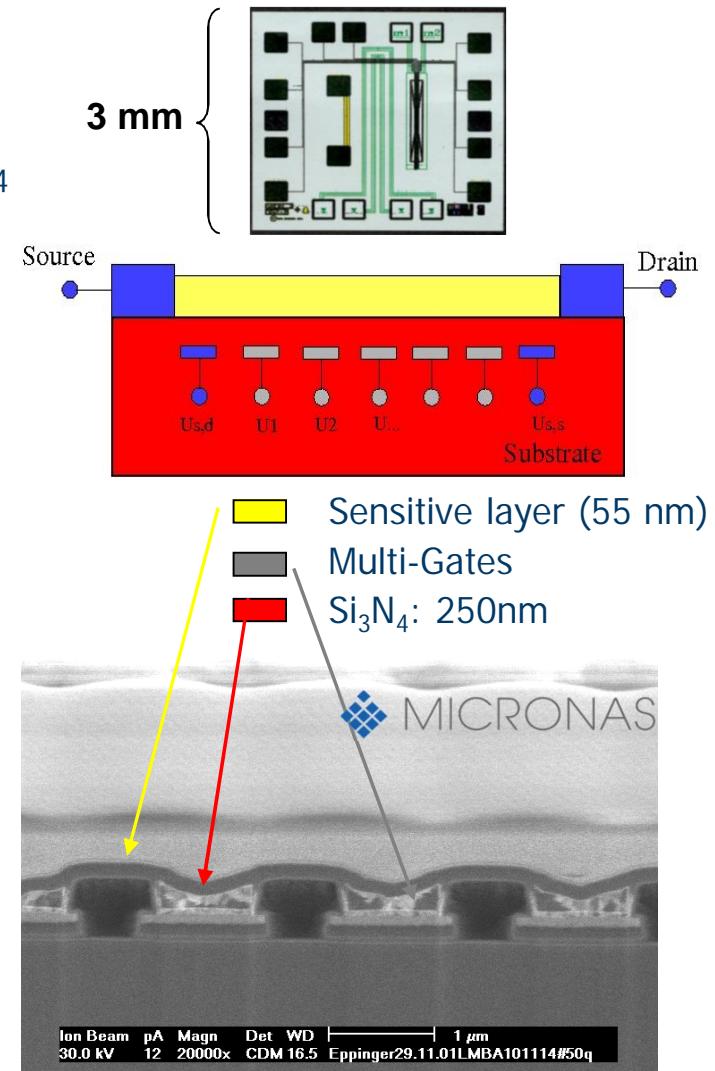


# Multigate concept



# Technology

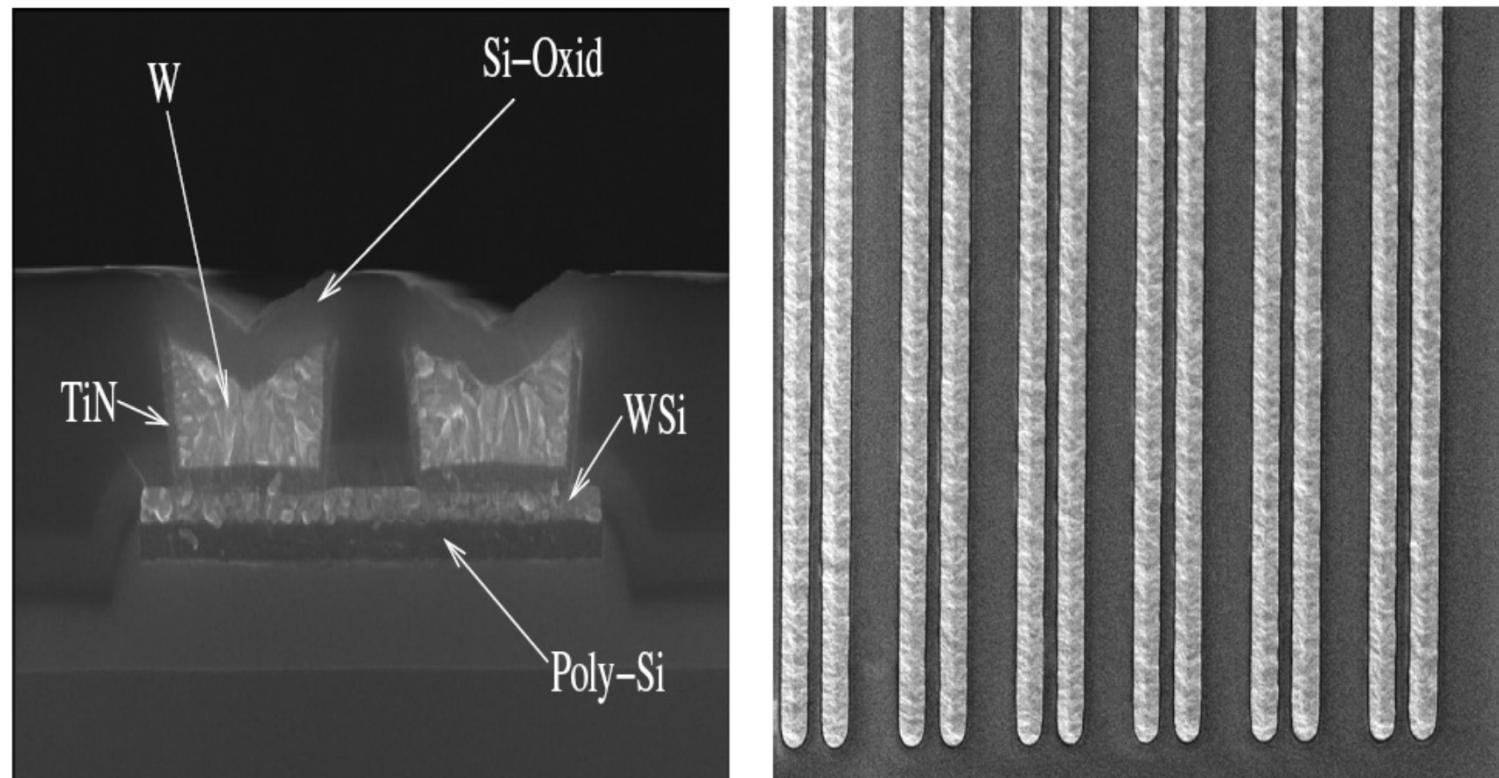
- SEM image „Multi-Gates“ Structure\*
  - Source-Drain TFT
  - Semiconductor  $\text{SnO}_2$  and insulator  $\text{Si}_3\text{N}_4$
  - Contact Pt
- Compact poly-crystalline\*\*
  - Non predominant orientation
  - Grain size from 5 to 25 nm (**Flat-band condition**)
- Sensitiven layer 55 nm



\*M. Lehmann, H. Frerichs \*\*J. Wöllensteini

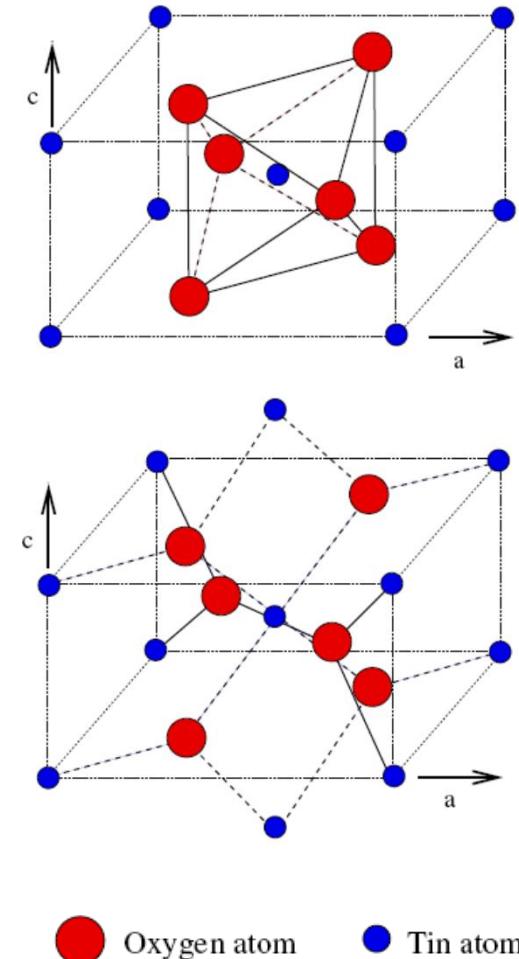
## ■ Gate electrode

- Multi-Gate-Structure, top view (Right)
- Gate electrode, Crossection (Left)



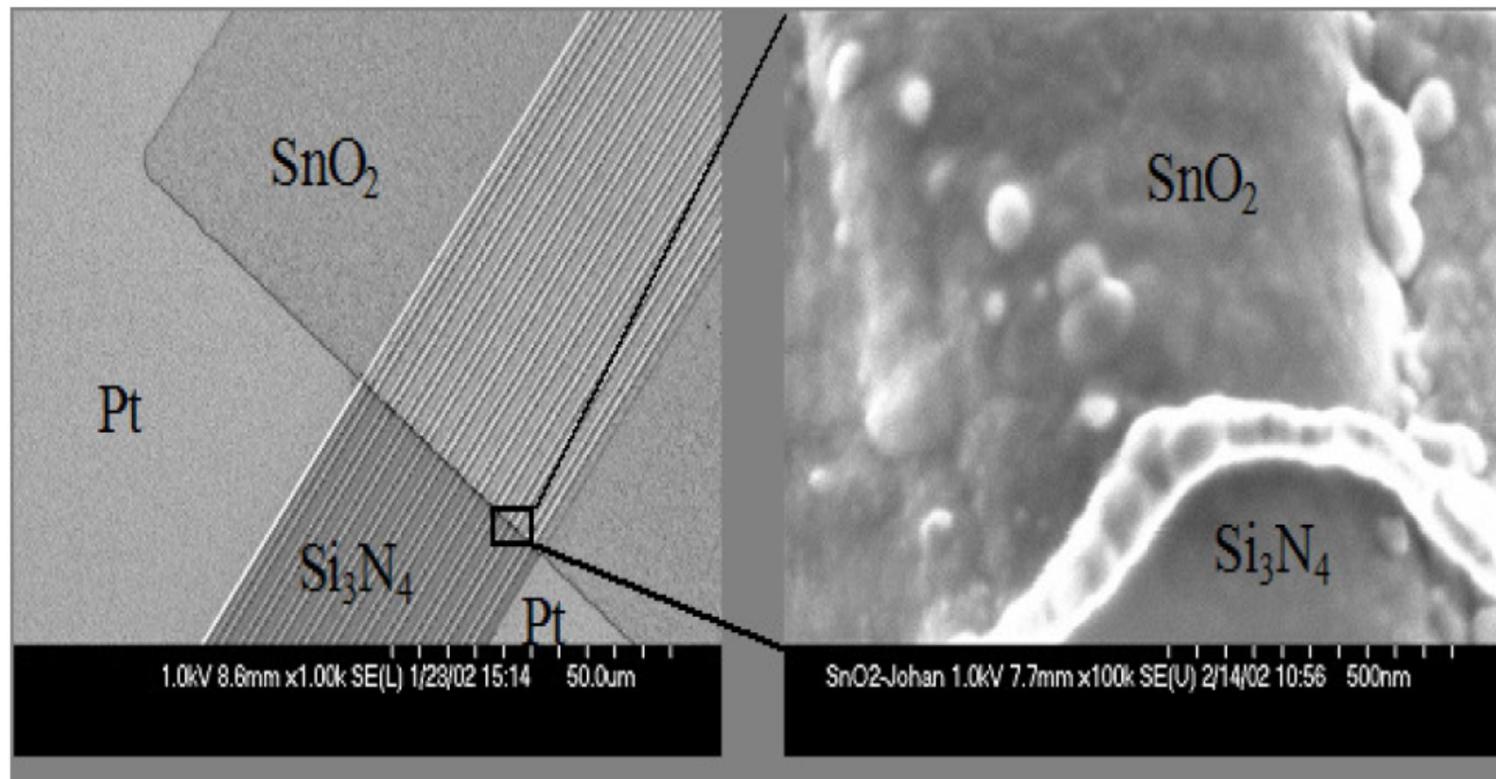
# Material SnO<sub>2</sub>

- Octahedron and rutile ( $a=b=0.47\text{nm}$ ,  $c=0.32\text{nm}$ , O-Sn-O=77°20')
- Inert to the acid-base reaction
- Heat of formation  $\Delta H=1.9E13(\text{J/mol})$
- Density at 300 K is 6,95 (g/cm<sup>3</sup>)
- Melting point at 1630°C
- Direct band gap 3,5 eV
- Small electron effective mass  $m^*=0,275 m_0$  , good conductivity
- Neutral oxygen vacancies ( $V_0$ , they play role as donators) form energy states at 0.035 eV and 0.140 eV
- It is a **n-type semiconductor** due the existence of native donators levels



## ■ SEM view of a multi-gate structure

- Source and drain area of a multiple TFT (left)
- $\text{SnO}_2$  and  $\text{Si}_3\text{N}_4$ , on right



# Electric model

## Sensor-Modell:

- Poisson + continuity equation:  
Signal behaviour

$$\Delta \cdot \psi = -\frac{\rho}{\epsilon_0 \cdot \epsilon_r}$$

$$\nabla \overrightarrow{J_{th}} + e^- \cdot \frac{\partial}{\partial t} (p - n) = 0$$

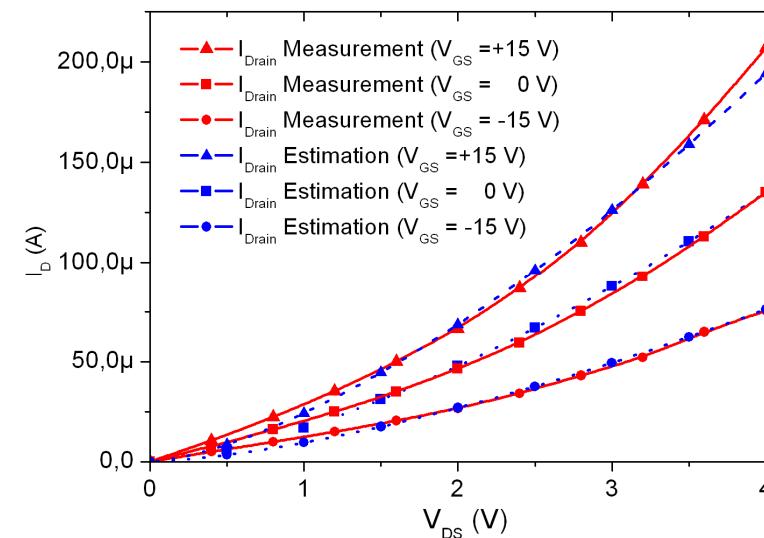
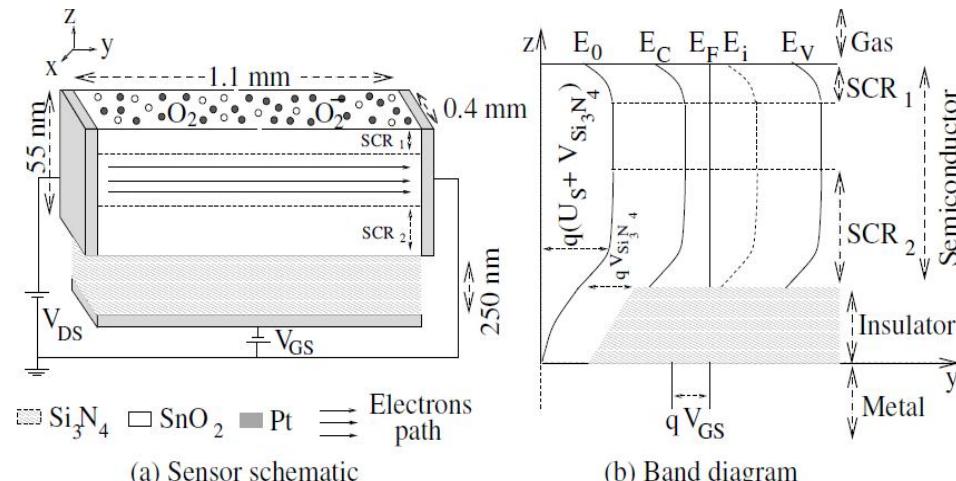
$$Q_{SS} = Q_{SCR}$$

- Wolkenstein-Modell: Surface reaction

$$\theta(p) = \frac{\beta \cdot p}{1 + \beta \cdot p}$$

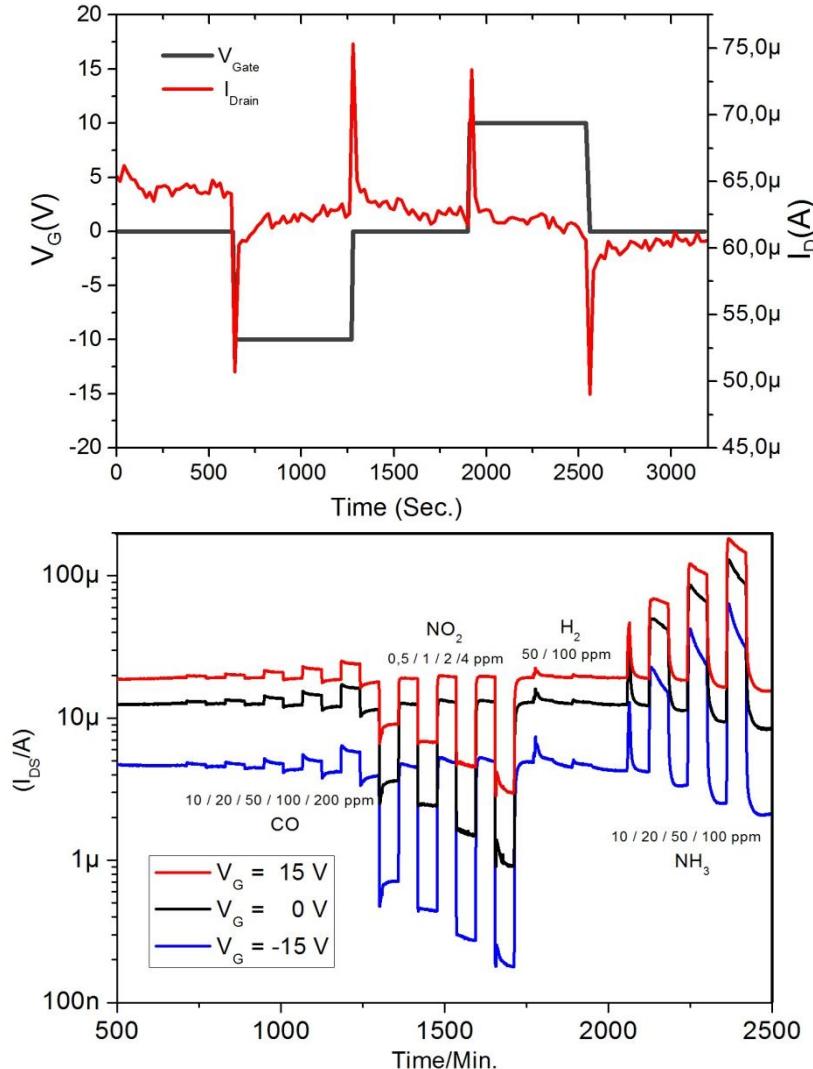
$$\beta = b \left\{ f^0 \left[ 1 + \frac{v^- f^-}{v^0 f^0} \cdot \exp \left( \frac{E_t^f - E_c}{kT} \right) \right] \right\}^{-1}$$

- Drain-current: Controlled by  $V_{DS}$
- Sensor works like MOSFET with reaction to gases



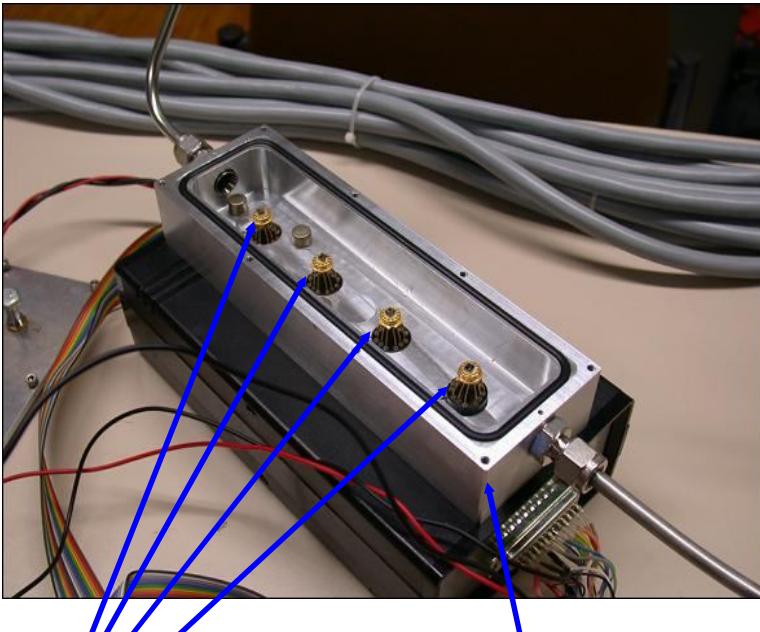
# Gas measurements

- Ideal current signal
  - Proportional to the gate potential
  - Not time depending (stable)
- Measurements with GasFET
- Solution: Pulse at the gate
- Parameter:
  - Synthetic air 80% N<sub>2</sub> and 20% O<sub>2</sub>, RH 40%, Work temperature 200°C, V<sub>DS</sub>=1 V



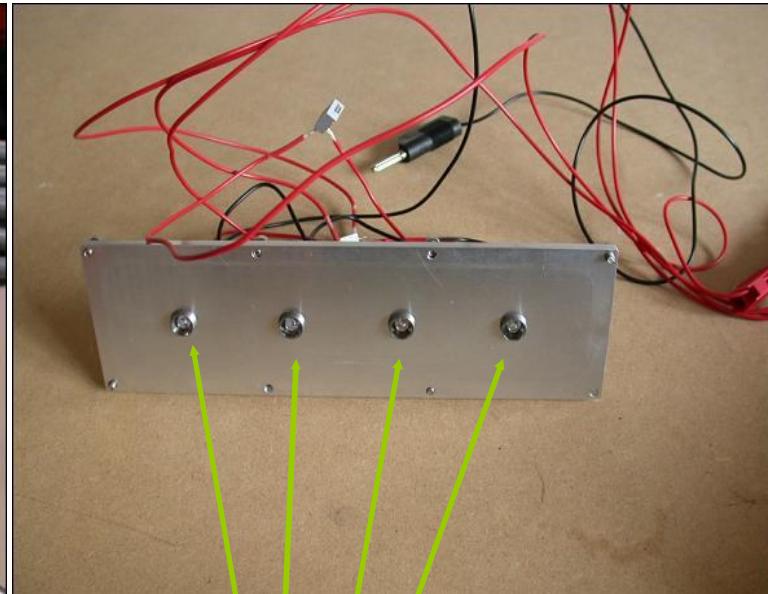
### Lichteinfluss auf gassensitive Nanoschichten

Messaufbau:



Gassensoren

Messkammer

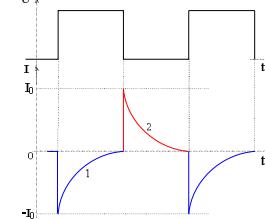
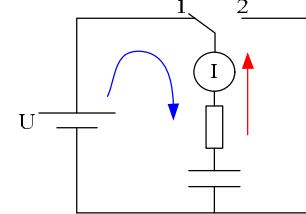


LEDs bei verschiedenen Wellenlängen  
(blau und UV)

# Signal drift

## MIS-condensator

- Low:  $C = 2.2 \text{ pf}$



## Charge carriers Generation

- Illumination (LED 340 nm)
  - $t_{\text{ON}}$ : 40-60 s
  - $t_{\text{OFF}}$ : 80-120 s

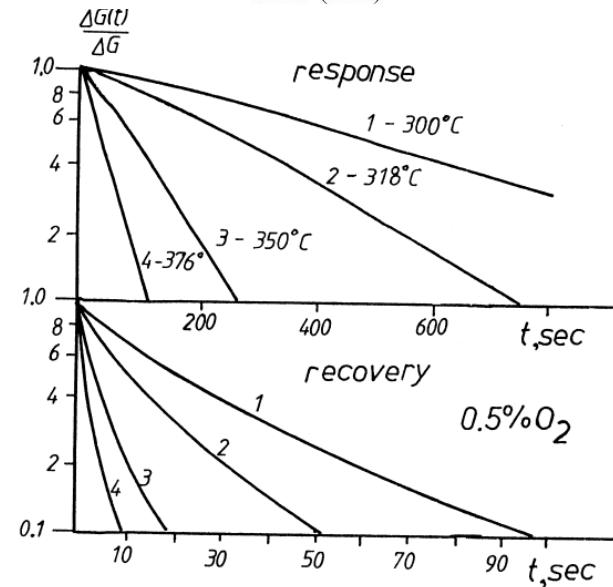
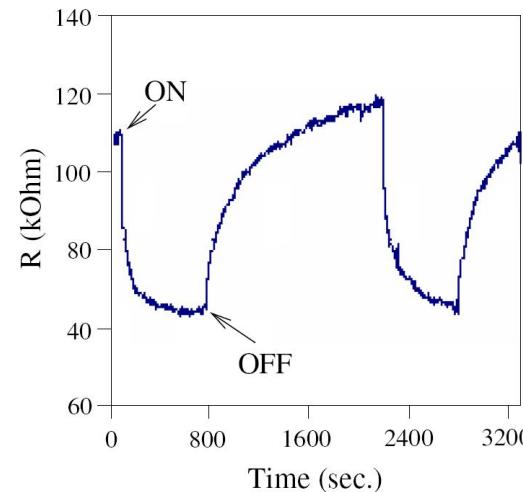
## Surface kinetic reaction\*

- Simulation parameters:  $10^{10}$  times smaller than real parameters

## Chemical reaction

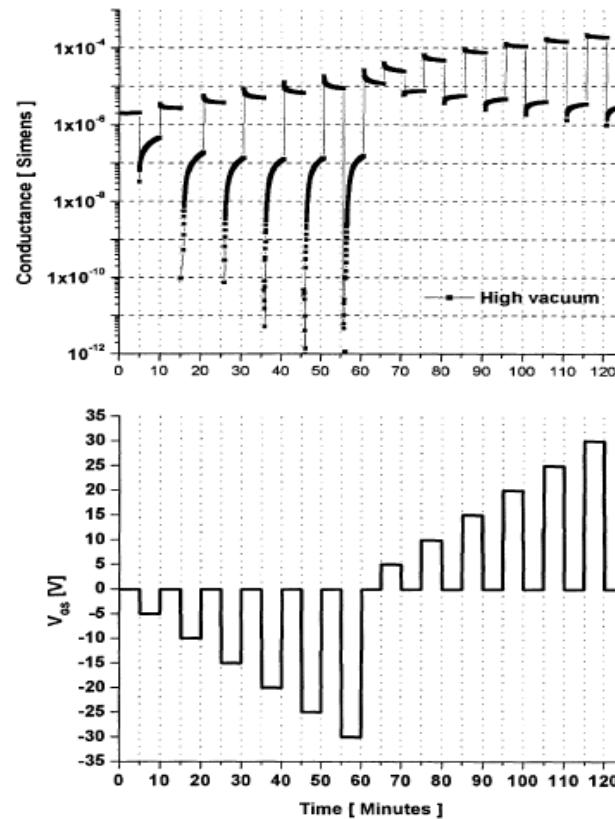
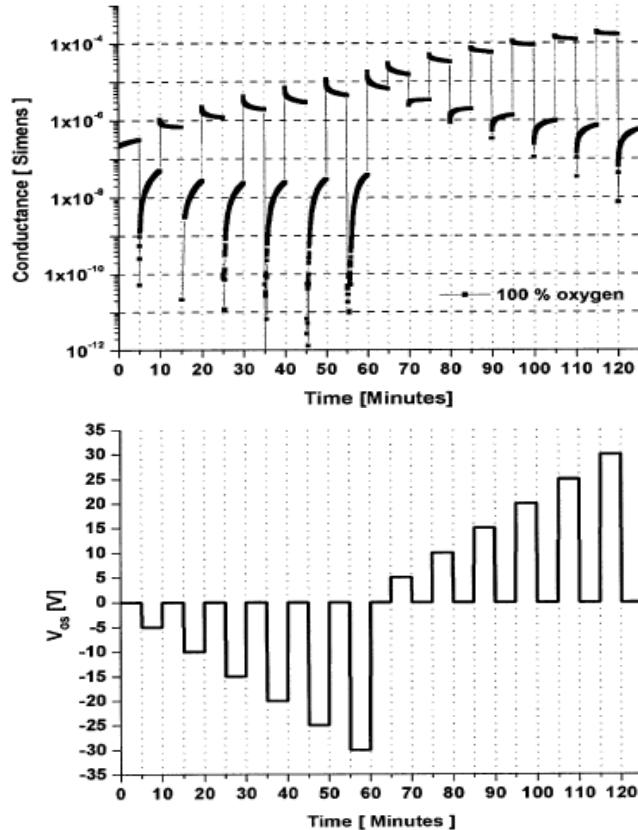
- Vacuum measurements, same signal behaviour

\*V. Brynzari et al. Sensors and Actuators B, 61, 143-153, 1999



# High vacuum

H.L. Tuller, Massachusetts Institute of Technology\*



„We're all in the same boat“

## Doping diffusion by electric field

- Fick diffusion

$$\frac{\partial N}{\partial t} = D_{vac} \Delta \cdot N$$

- Fokker-Planck-Equation

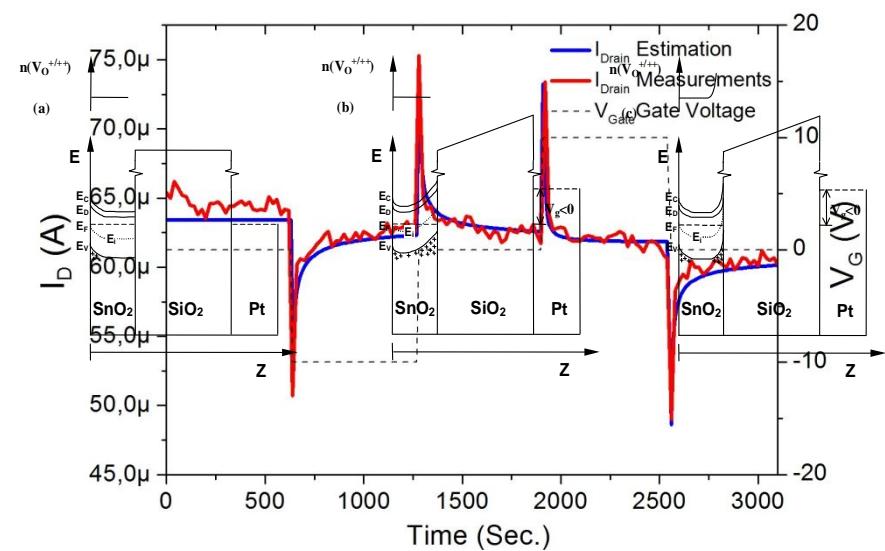
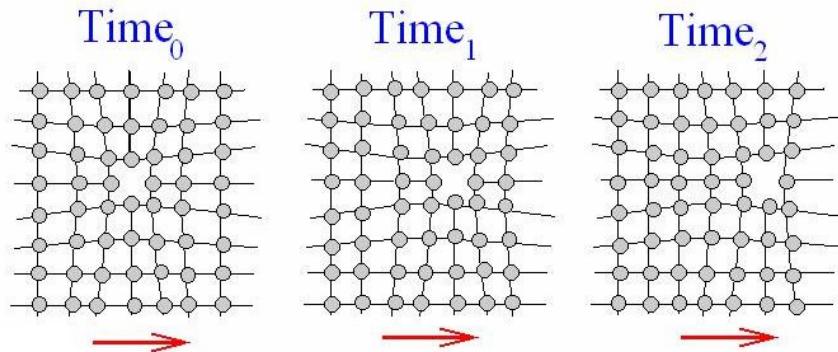
$$\frac{\partial N}{\partial t} = \frac{e^-}{m_{vac}^{eff}} \nabla \cdot (E \cdot N) + D_{vac} \Delta \cdot N$$

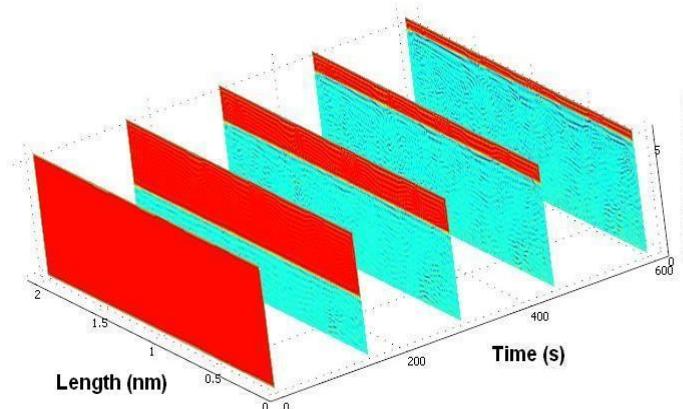
- Drift-Diffusions-Equation

$$\frac{\partial n}{\partial t} = -\mu_n \nabla \cdot (E \cdot n) + D_{e^-} \Delta \cdot n$$

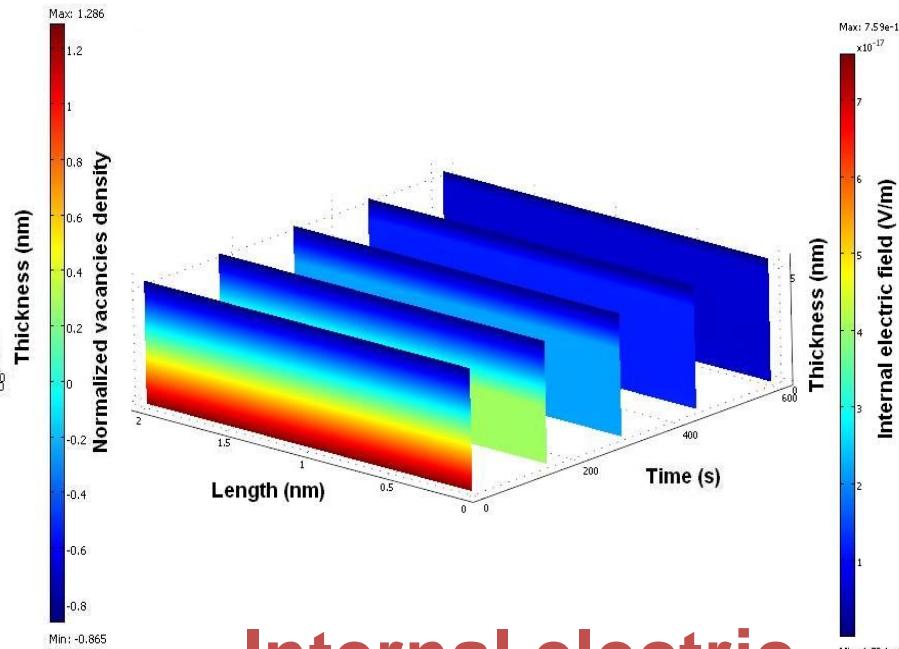
- Coupling of FP+DD by Poisson-equation

$$\Delta \cdot \psi = -\frac{e^-}{\epsilon_0 \cdot \epsilon_r} (n - N)$$





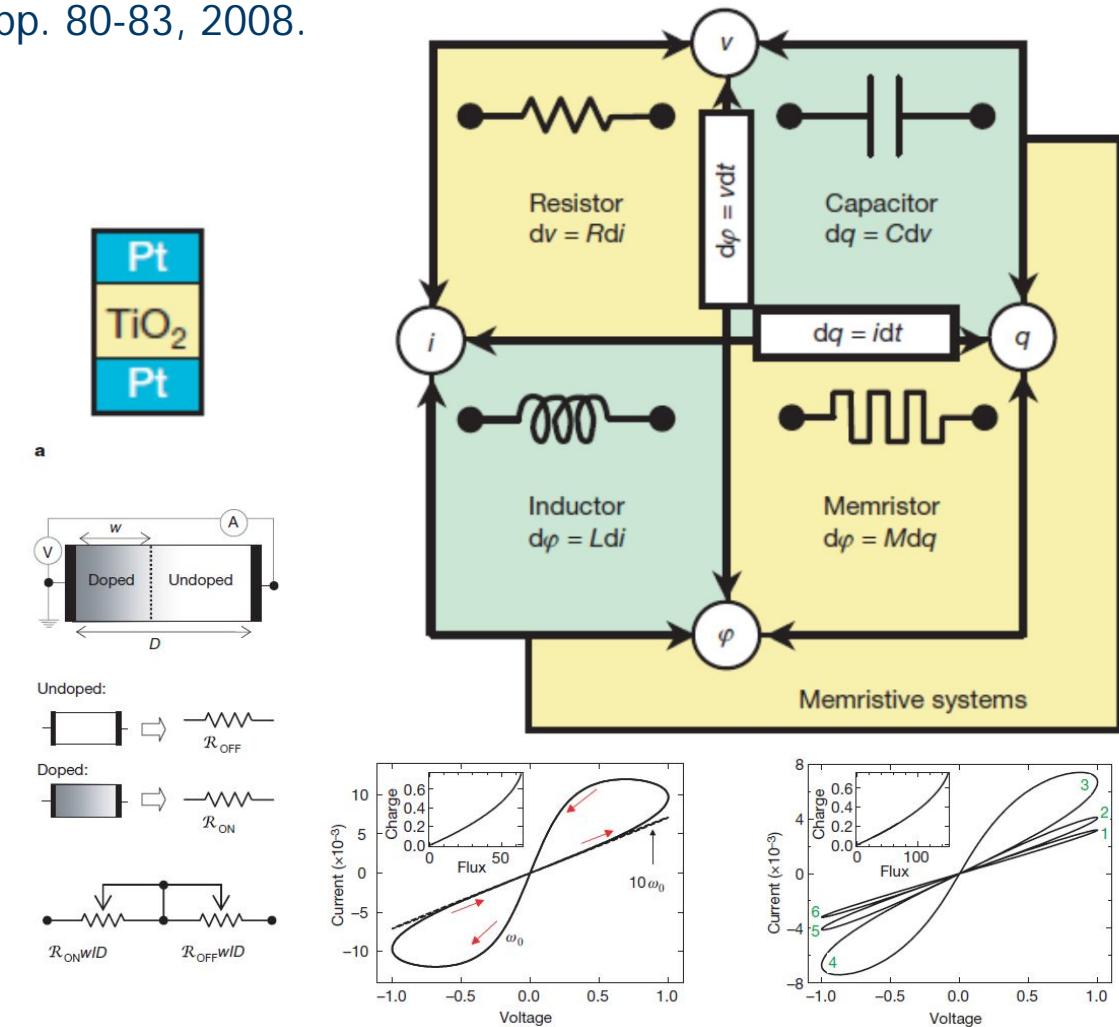
**Normalized  
vancies  
distribution.**



**Internal electric  
field  
distribution.**

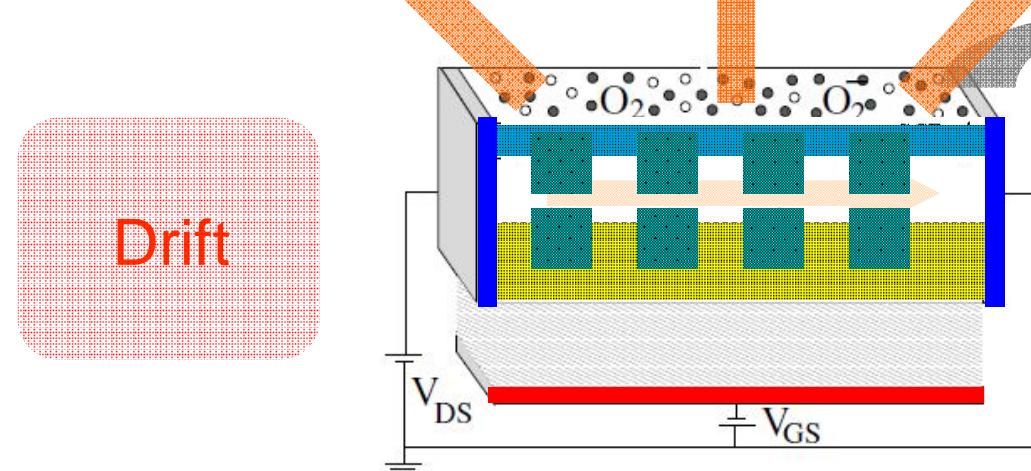
# Memristor

D.B. Strukov, Nature, Vol. 453, pp. 80-83, 2008.



# Summary

EAE Adsorption      Co-Adsorption      Gas-adsorption



Doping

Continuity equation

Poisson-equation

Standard Halbleiter-Theorie

Electric Potential

	$N_D$	$I_{DS}$	$SCR_1$	$SCR_2$
Standard	✗	✓	✗	✓
Gasad.	✗	✓	✓	✗
Co-Ads.	✗	✓	✓	✗
EAE	✗	✓	✓	✓
Drift	✓	✓	✓	✓

# Outline

- Motivation
- Classic semiconductor theory
  - Ideal „Bulk“
- Semiconductor theory at the bulk edge
  - Semiconductor surface-gas interactions
- One example of device operation
  - Electro-adsorptive effect
  - Thin film transistor
  - Drift of vacancies (dopants)
- Remarks

# Conclusions

- Charge transport in semiconductors can be accurately described
  - Semiconductor theory
- IMPORTANT: Not only the material properties are important in the charge transport also geometry!
- Gas-solid interaction influences the transport of charge in solid due to a charge transfer process.
  - Atomistic interaction at the surface modifies the conductivity of solid
  - It can be accurately described by an extended semiconductor theory
    - ⌘ Wolkenstein theory
    - ⌘ Charge transfer (band bending)
  - IMPORTANT! Check if this model is applicable directly to your system
    - ⌘ Grain size effects (Debye lenght)
- Note: In catalytic processes the surface change continuosly:
  - Dynamic phase transitions



# Greetings



# Thank you!



JOHANNES  
**GUTENBERG**  
UNIVERSITÄT  
MAINZ



 MICRONAS



  
Fraunhofer Institut  
Physikalische  
Messtechnik



  
Bundesministerium  
für Bildung  
und Forschung

Alexander von Humboldt  
Stiftung/Foundation

