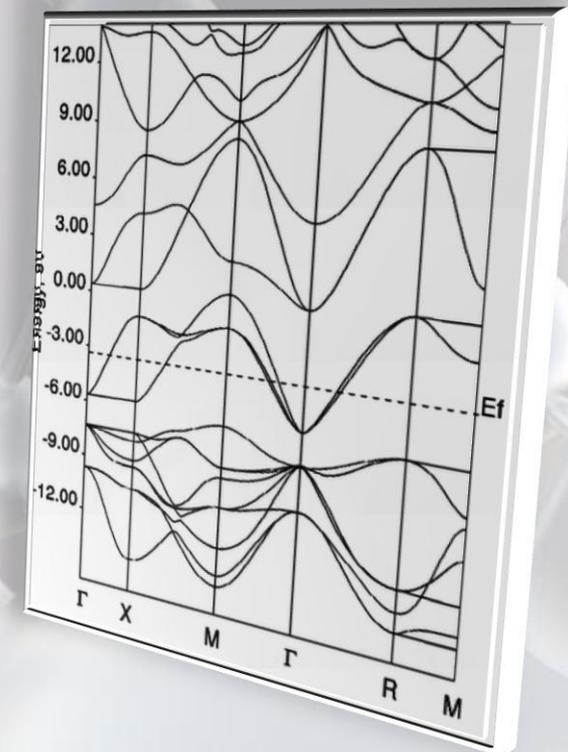
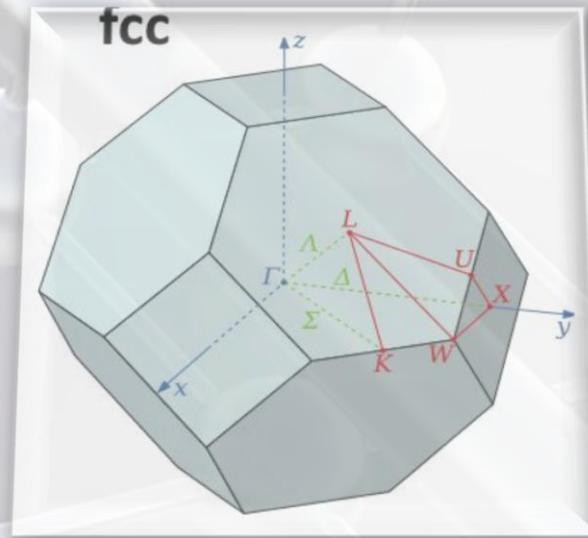
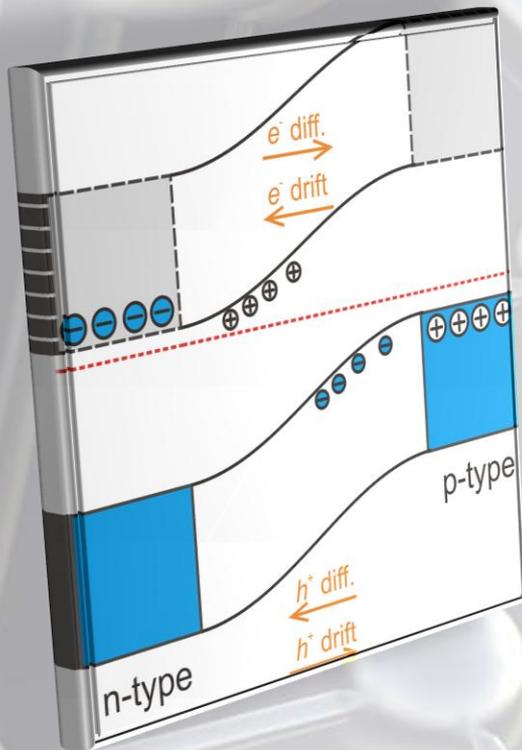


# Electronic Structure of Solids

FHI-AC Lecture Series

27<sup>th</sup> October 2017



# Outline

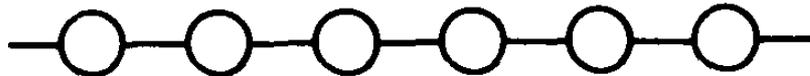
- **Part 1: Basics of Bands**
  - How bands are formed from molecular orbitals
  - The meaning of k-space
  - Density of states
  - Electron filling
  - The Fermi level
  - Metals vs semiconductors
  - Conduction
  - Doping
- **Part 2 : The Fermi level, Work Function and Band Bending**
  - Significance of Fermi level
  - Understanding work function
  - Interfaces and band bending
  - Relevance to catalysts

# Part 1: Basics of Band Structure

# 1. Chemist's View of Electronic structure in Solids

- This material is taken from:  
“Solids and Surfaces: A Chemist's View of Bonding in Extended Structures”  
By Ronald Hoffmann

# 1.1 Orbitals and bands in 1-dimension

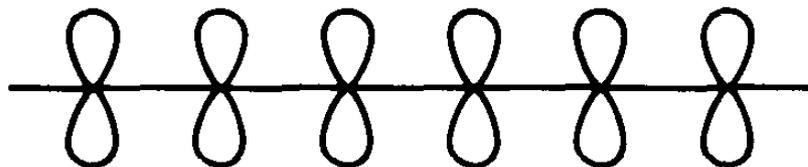


1

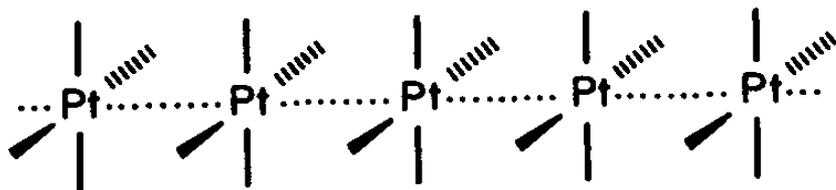
Equally-spaced H atoms



$\pi$ -system of delocalized polyene



2

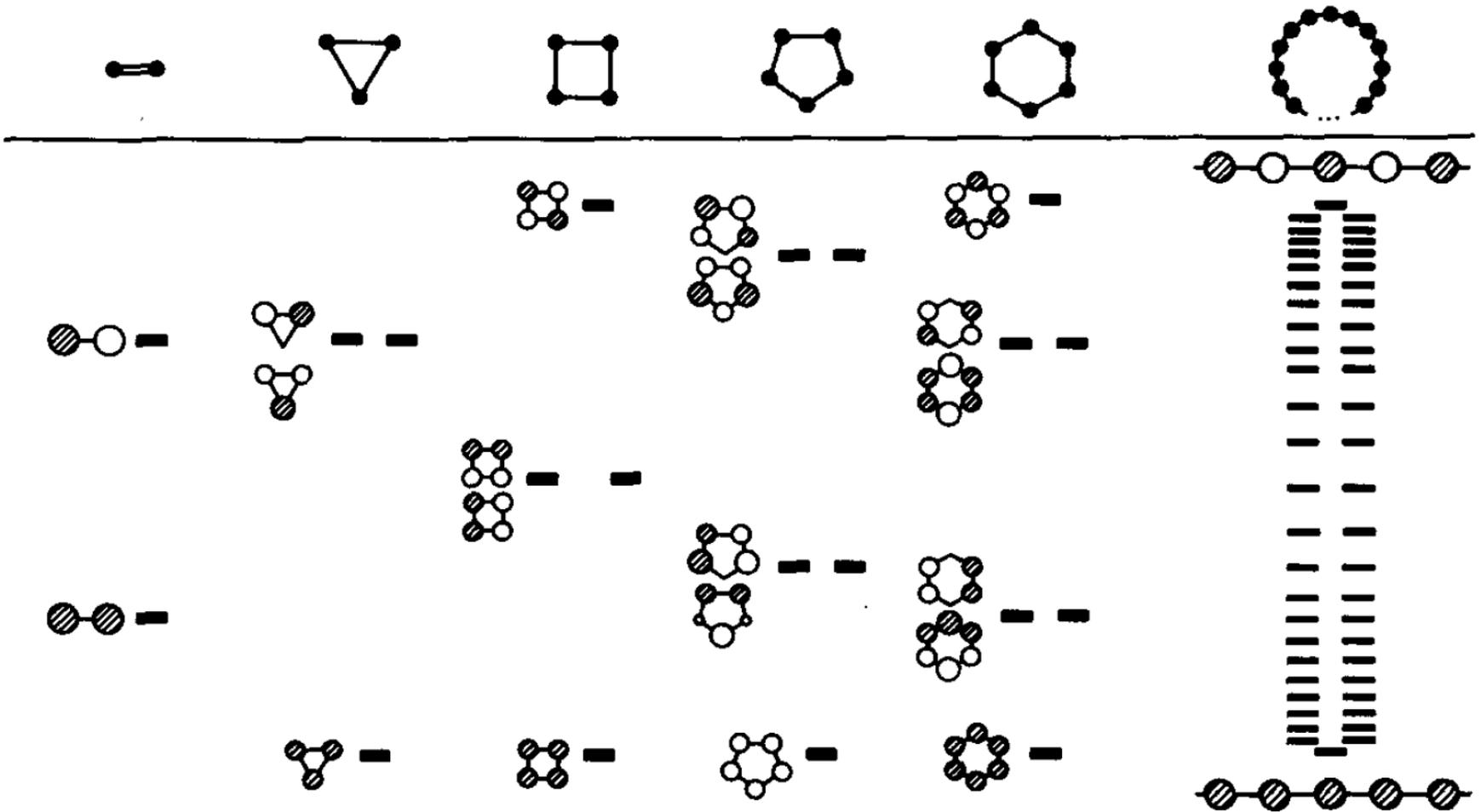


Stack of square planar Pt(II) complexes  
(e.g.  $\text{PtH}_4^{2-}$ )

3

# 1.1 Infinite chain of H atoms

cyclic boundary condition

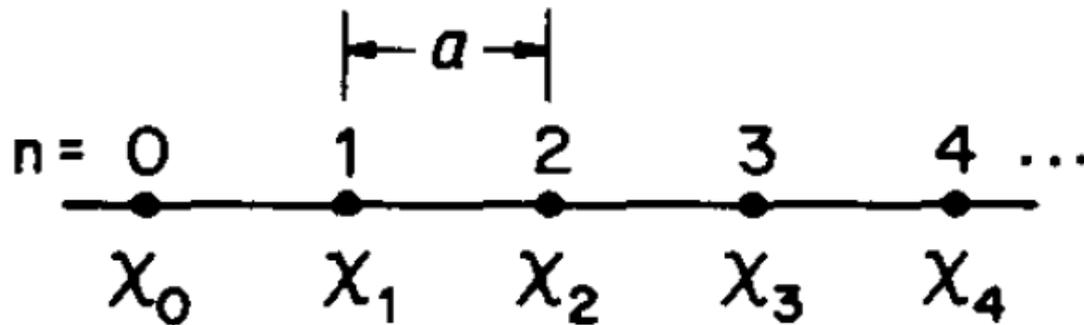


# 1.2 Bloch functions and band structures in k-space

Better way to write ... translational symmetry

use an index  $n = 1, 2, 3, 4, \dots$

Then there is a basis function (an H 1s orbital),  $\chi_0, \chi_1, \chi_2, \chi_3, \dots, \chi_n$



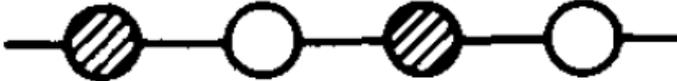
$$\psi_{\mathbf{k}} = \sum_n e^{i\mathbf{k}n a} \chi_n$$

# 1.2 Bloch functions and band structures in k-space

Quick test...

does it work for  $k = 0$  and  $k = \pi/a$  ?

$$\begin{aligned} k = 0 \quad \psi_0 &= \sum_n e^{0} \chi_n = \sum_n \chi_n \\ &= \chi_0 + \chi_1 + \chi_2 + \chi_3 + \dots \end{aligned}$$

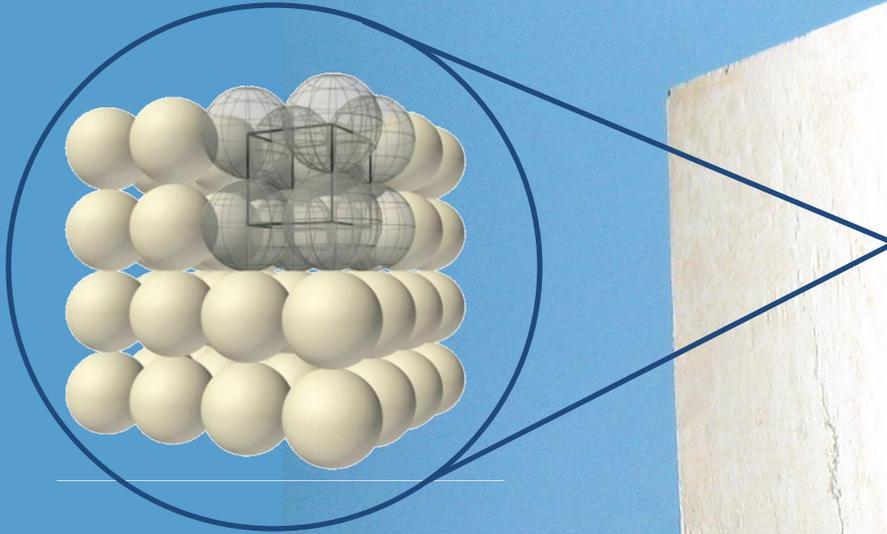

$$\begin{aligned} k = \frac{\pi}{a} \quad \psi_{\frac{\pi}{a}} &= \sum_n e^{\pi i n} \chi_n = \sum_n (-1)^n \chi_n \\ &= \chi_0 - \chi_1 + \chi_2 - \chi_3 + \dots \end{aligned}$$


Higher the magnitude of  $k$ , the more number of nodes.

$k$  has a range. Values outside this range are not unique, and reproduce values of lower  $k$   
 $|k| \leq \pi/a$  (This is called the first Brillouin zone, i.e. the range of unique  $k$ )

# 1.2 k-space

How many values of  $k$  are there?



As many as the number of translations in the crystal

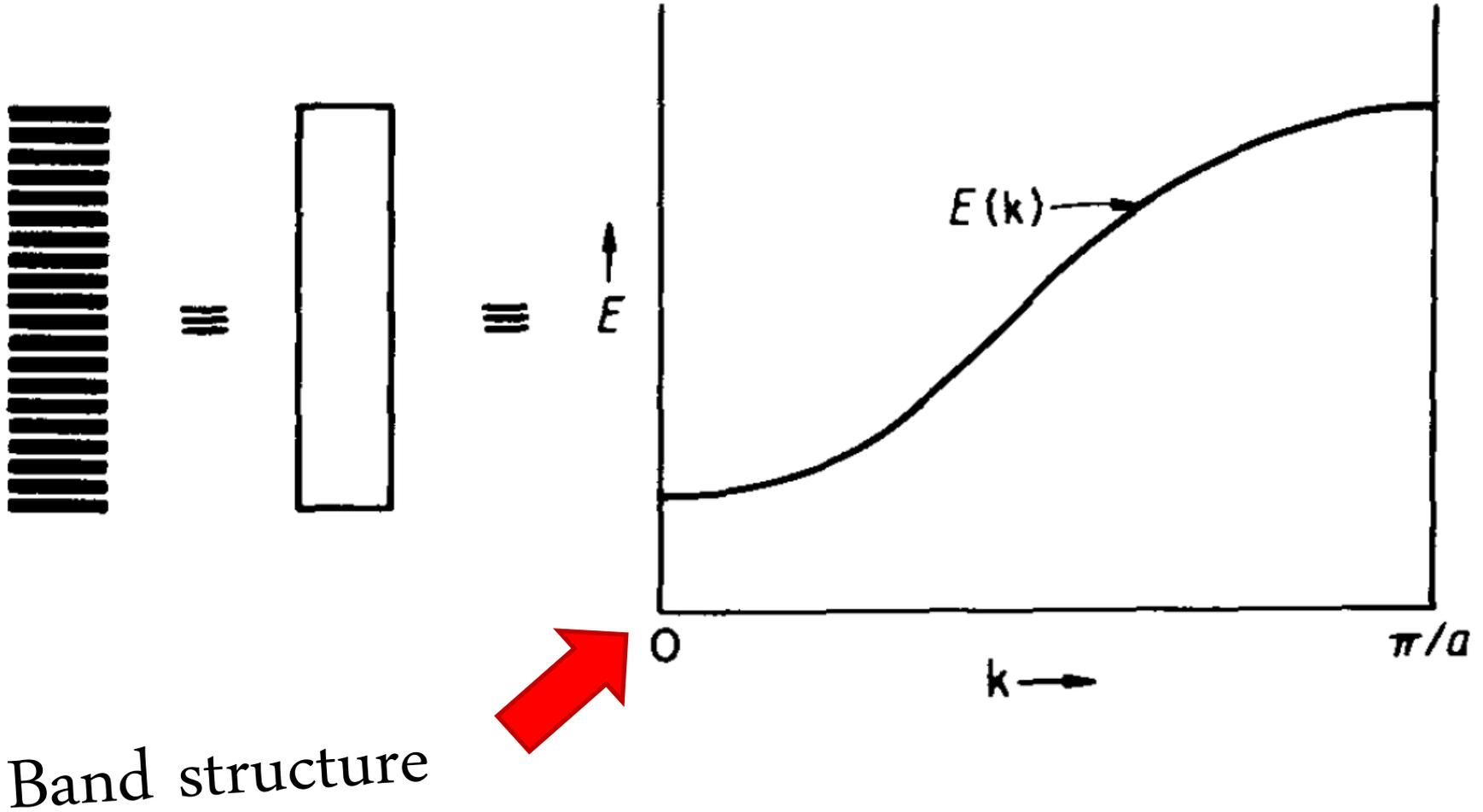
(i.e. as many as the number of microscopic unit cells in the macroscopic crystal)

$\sim 10^{23}$

# 1.2 k-space

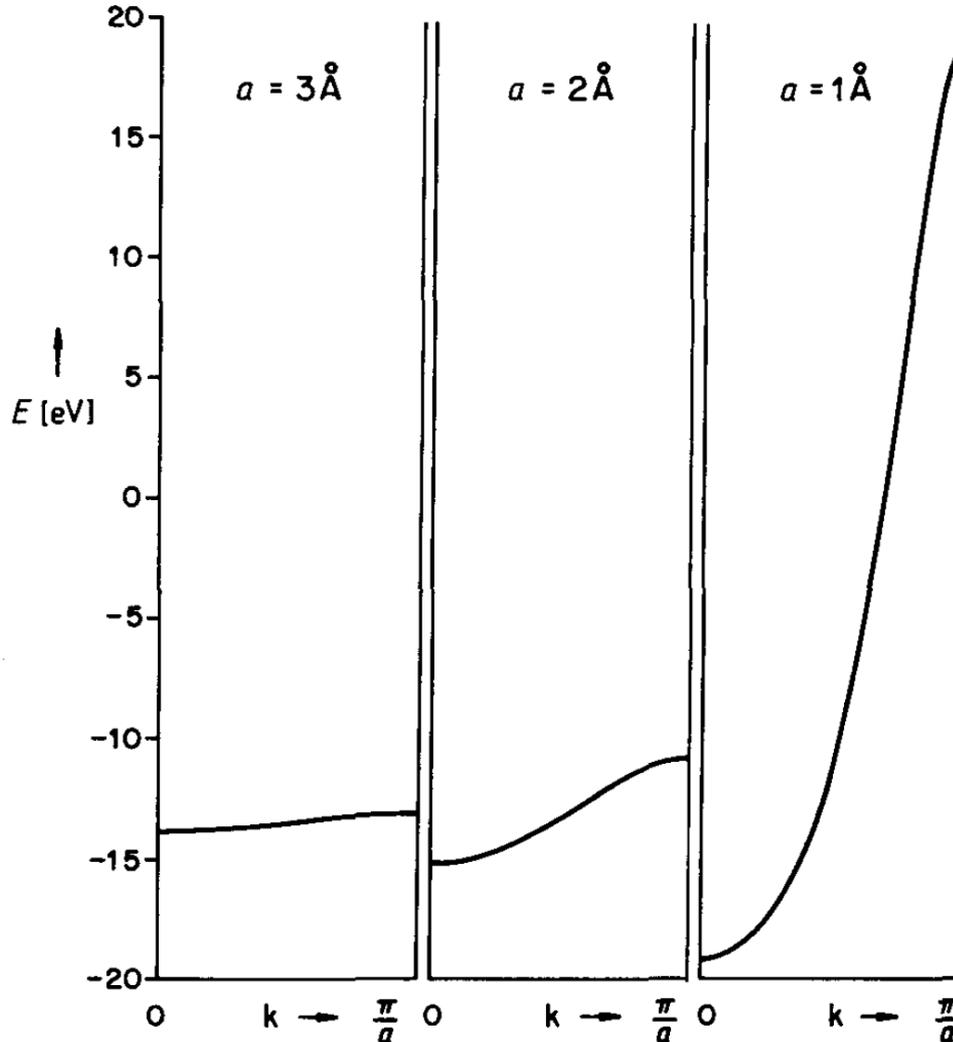
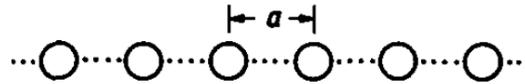
- k-space is also known as reciprocal space or momentum spaces
- Relationship between  $k = 2\pi/\lambda$  and momentum is from de Broglie's  $\lambda = h/p$
- i.e. k is proportional to p (momentum of the electron wavefunction)
- k is not only a symmetry label and a node counter, it is also a wave vector (i.e. it measures momentum)

# 1.2 k-space



# 1.2 Bandwidth

a.k.a. dispersion



The larger the atomic spacing, the closer the band resembles a free H atom.

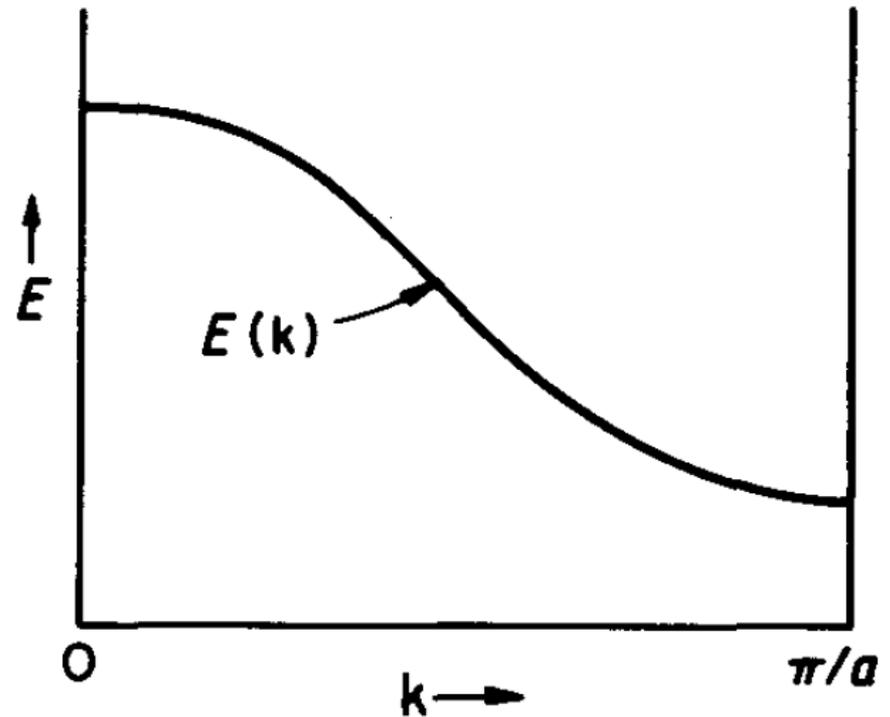
# 1.3 p-orbitals in solids

When the basis set (i.e. orbitals) have inversion symmetry (e.g. p-orbitals) then bands run in the opposite direction (i.e.  $k = 0$  has higher energy than  $k = \pi/a$ )

$$\psi_0 = \chi_0 + \chi_1 + \chi_2 + \chi_3 + \dots$$

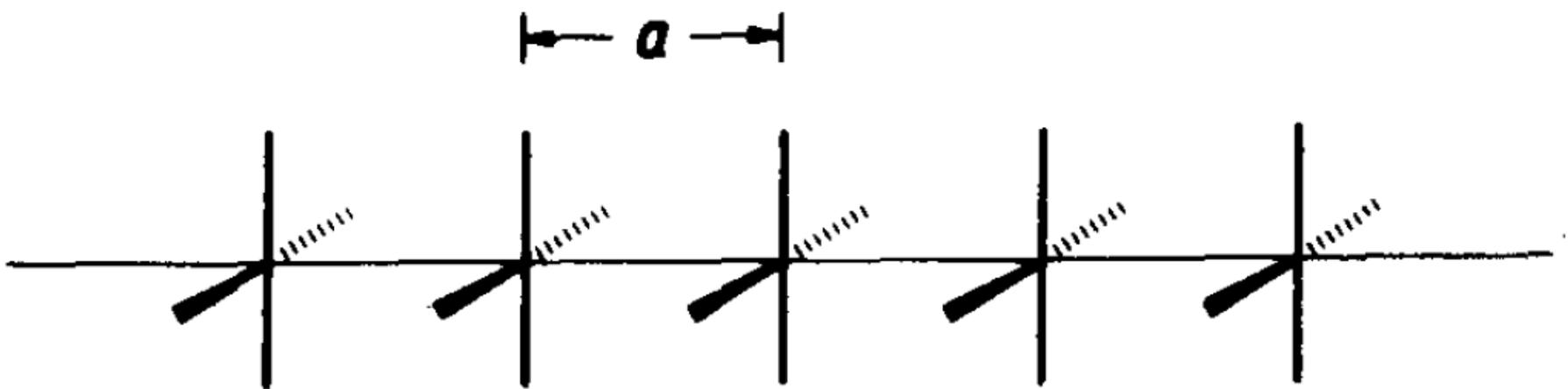


$$\psi_{\frac{\pi}{a}} = \chi_0 - \chi_1 + \chi_2 - \chi_3 + \dots$$

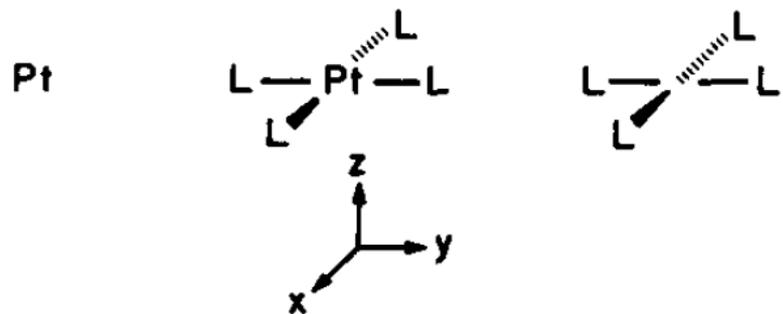
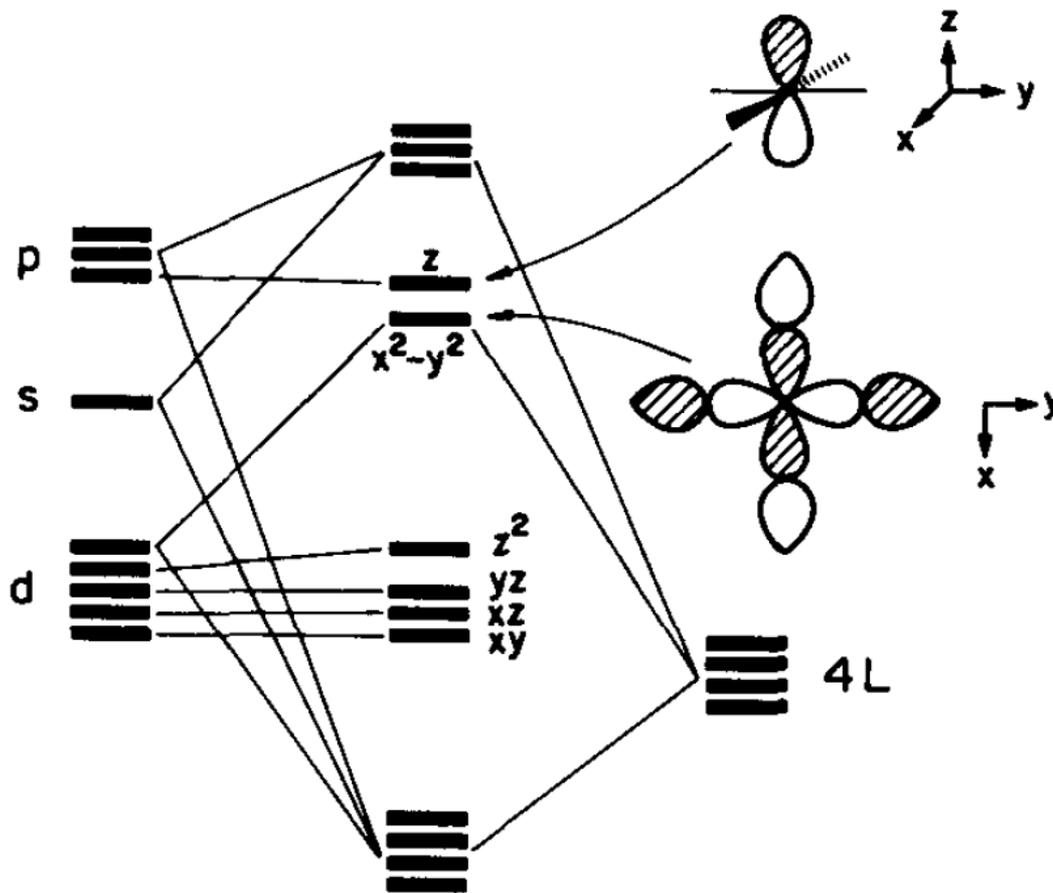


# 1.3 Stacks in solids

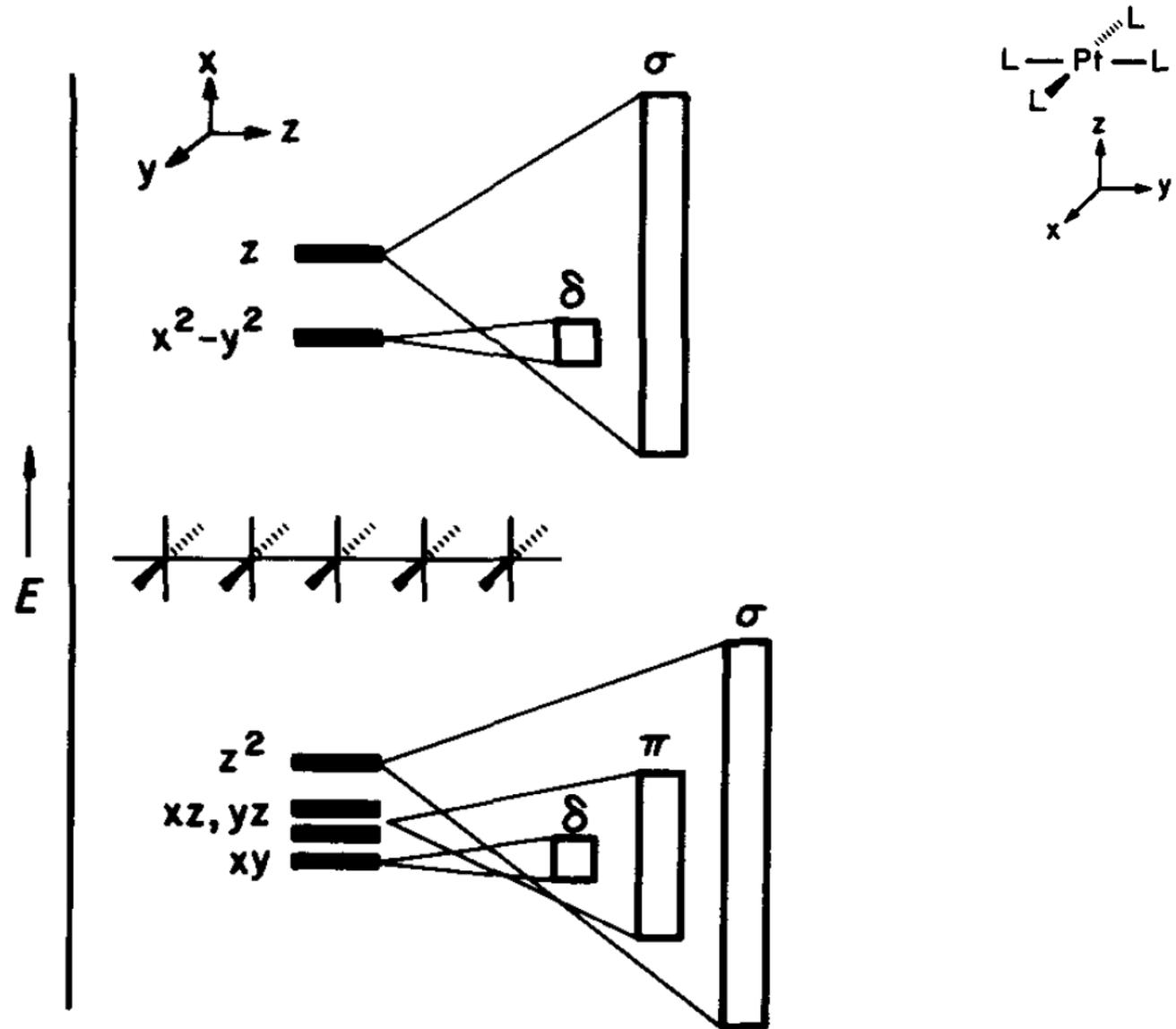
Getting more complicated  
with  $d^8$   $\text{PtL}_4$  complexes



# 1.3 Frontier orbitals of monomer



# 1.3 Bands from frontier orbitals



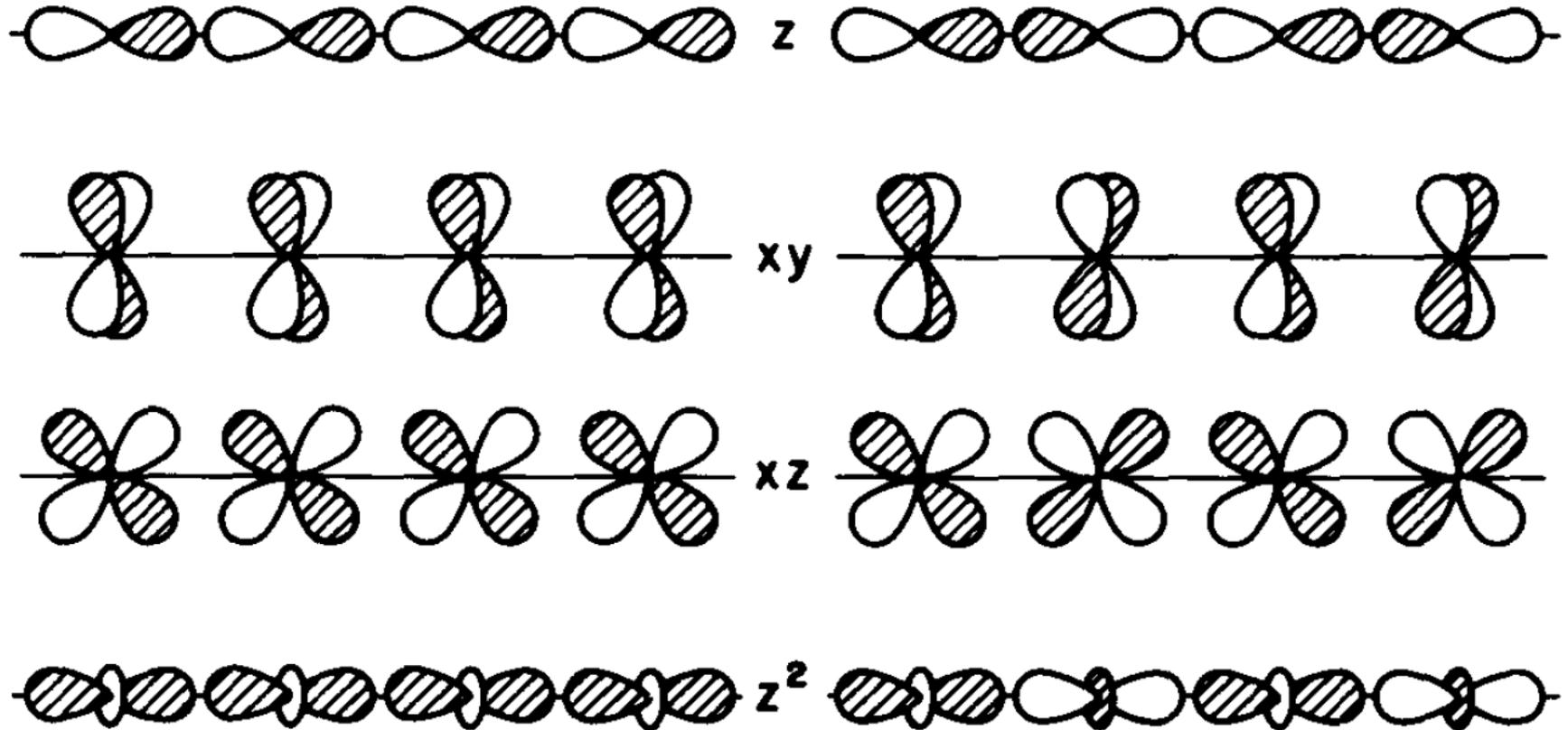
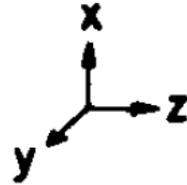
# 1.3 Orbitals at $k = 0$ and $k = \pi/a$

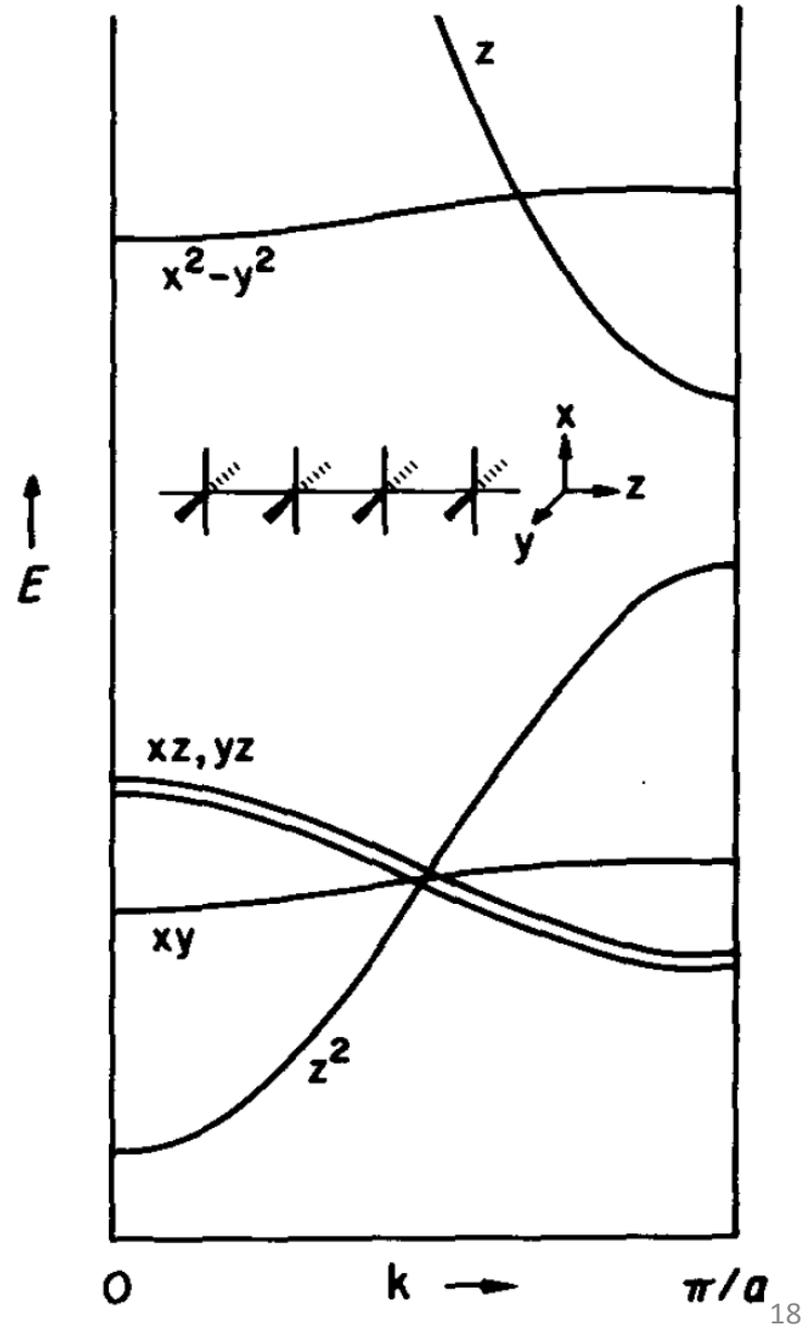
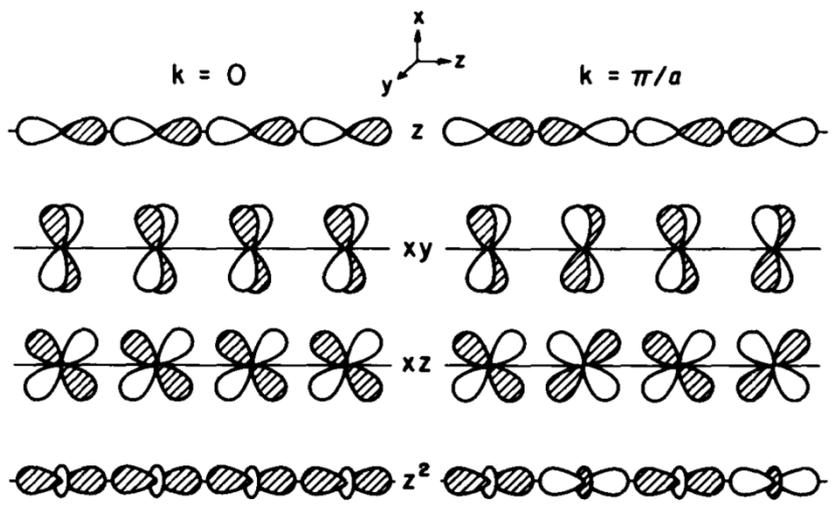
Zone center

Zone edge

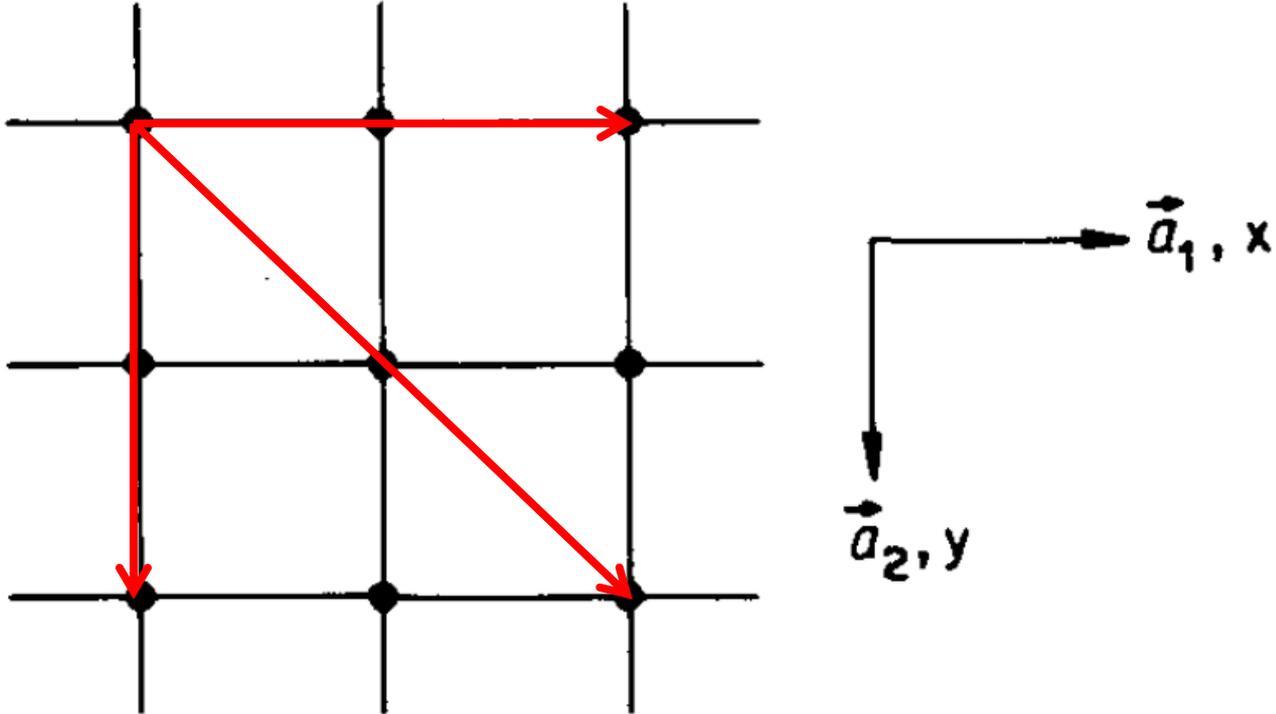
$k = 0$

$k = \pi/a$





# 1.4 Bands in 2 Dimensions



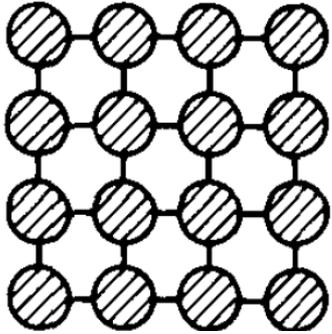
We can factor the Schrödinger equation of the crystal into wavefunctions along the x- and y-axis, each being identical to the 1D representations.

Then we have two kinds of wave vectors,  $k_x$  and  $k_y$

$$0 \leq |k_x|, |k_y| \leq \pi/a \quad (a = |\mathbf{a}_1| = |\mathbf{a}_2|)$$

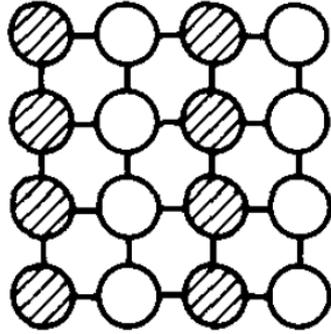
# 1.4 Solutions in 2D for s-orbitals

$$k_x = 0, k_y = 0$$



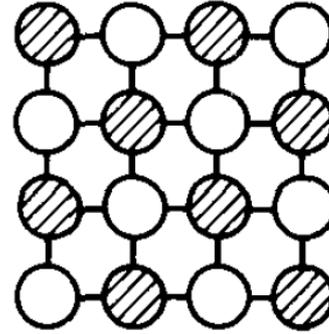
$\Gamma$

$$k_x = \pi/a, k_y = 0$$



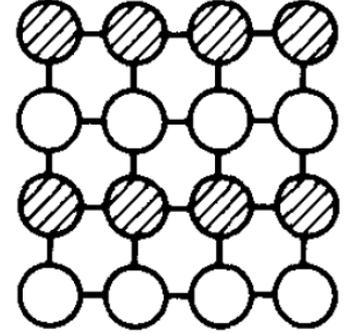
X

$$k_x, k_y = \pi/a$$

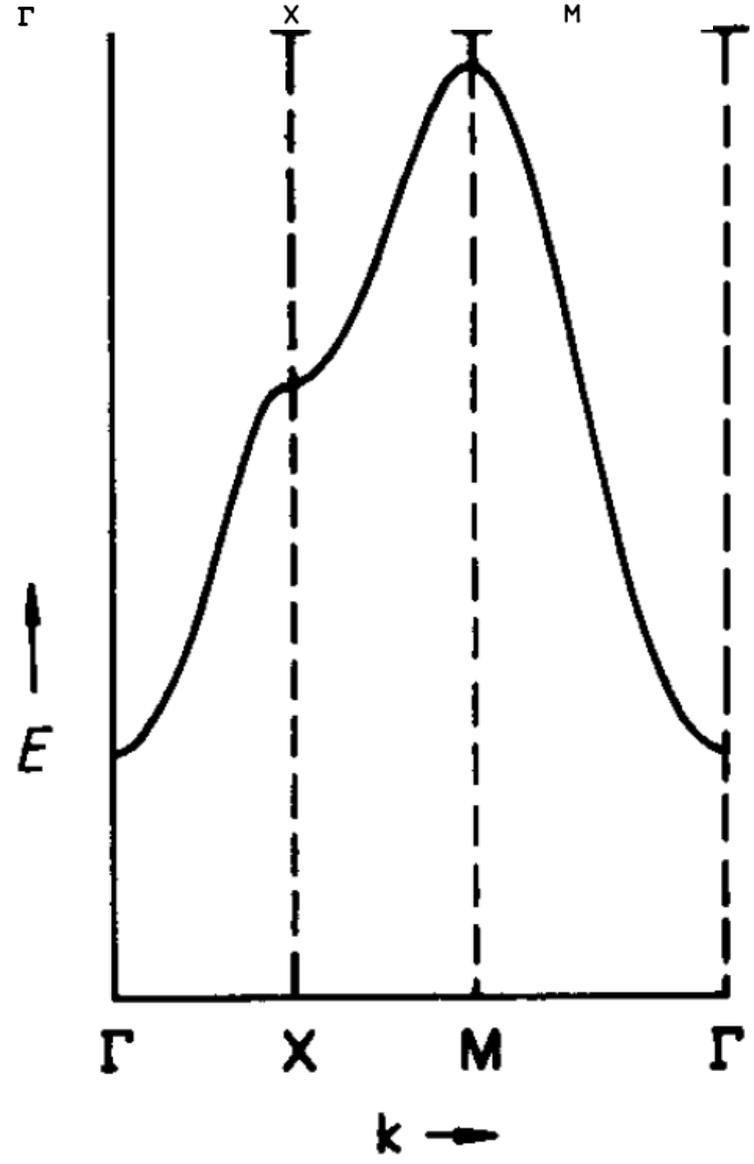
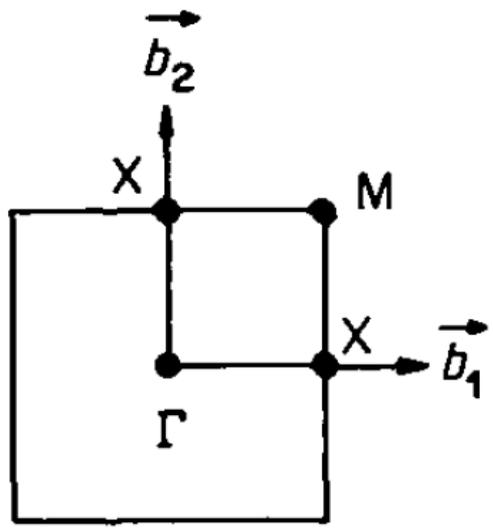
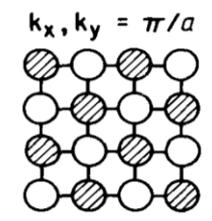
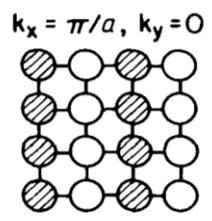
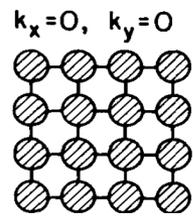


M

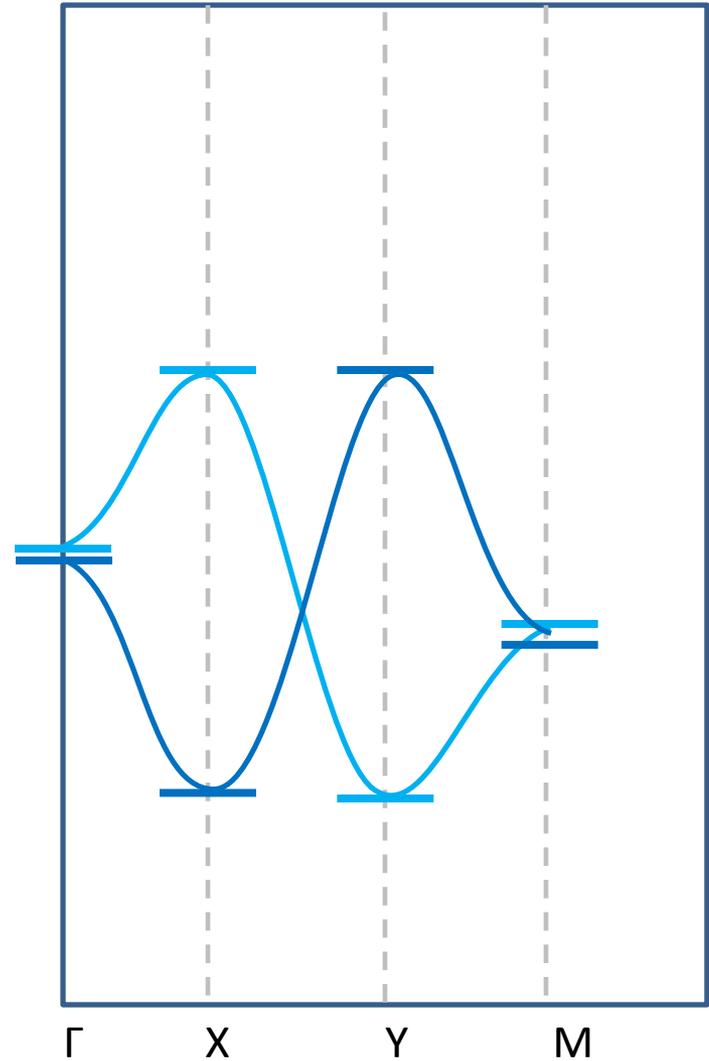
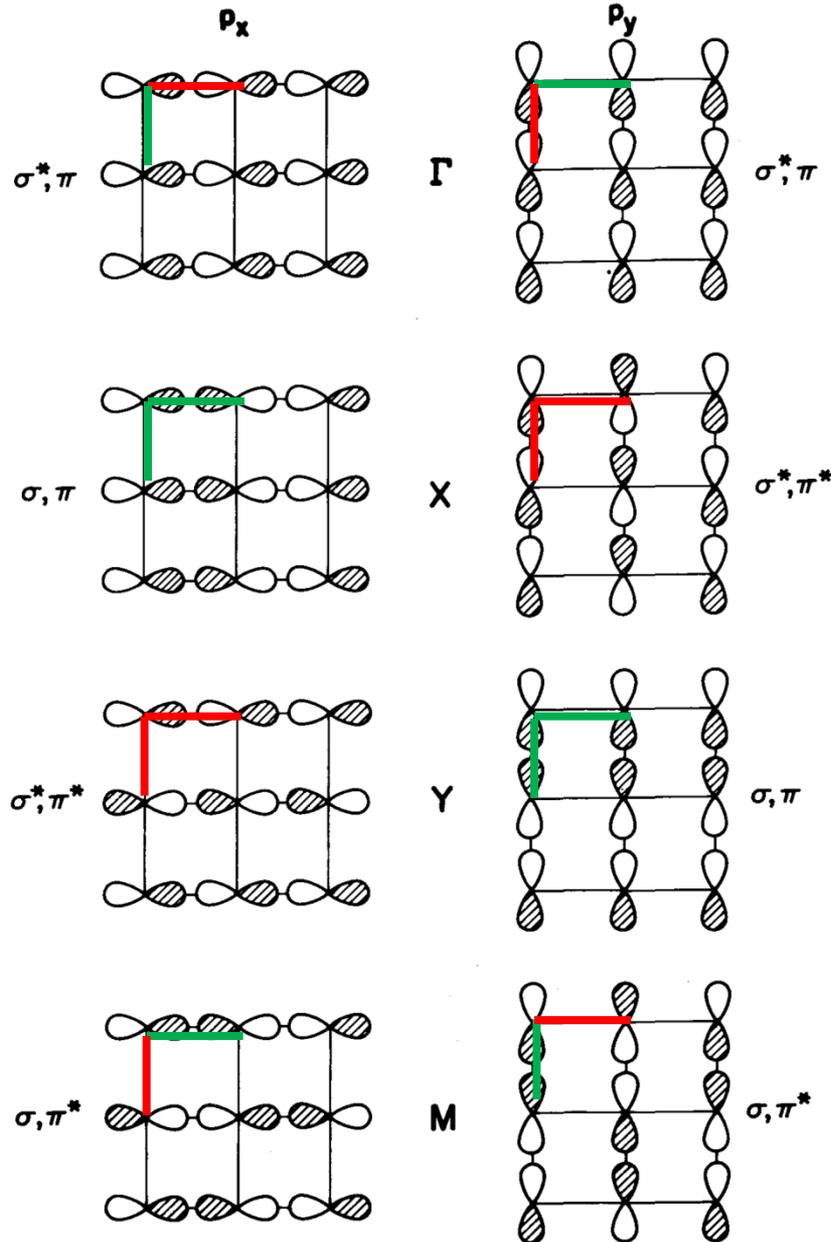
$$k_x = 0, k_y = \pi/a$$

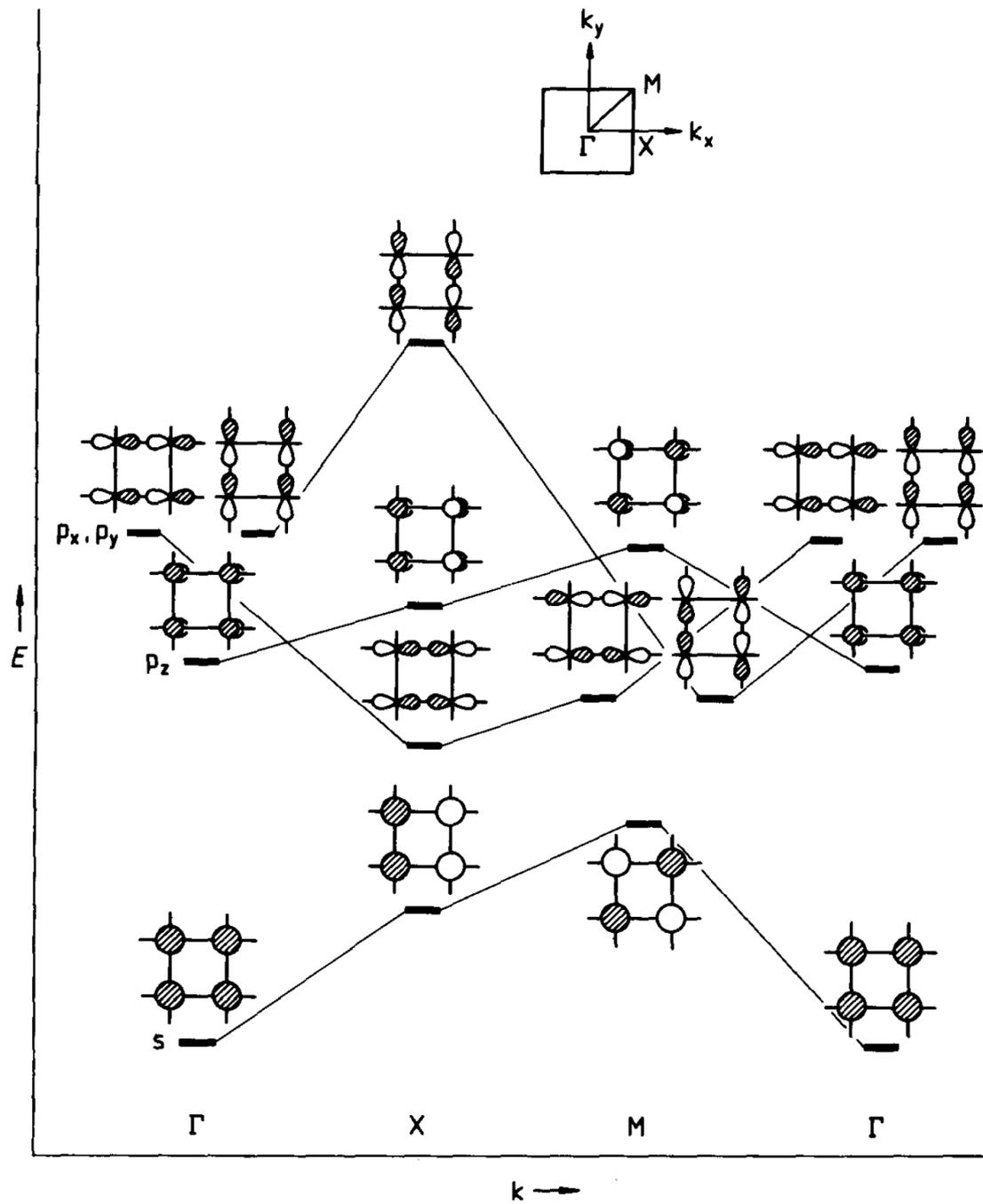


X

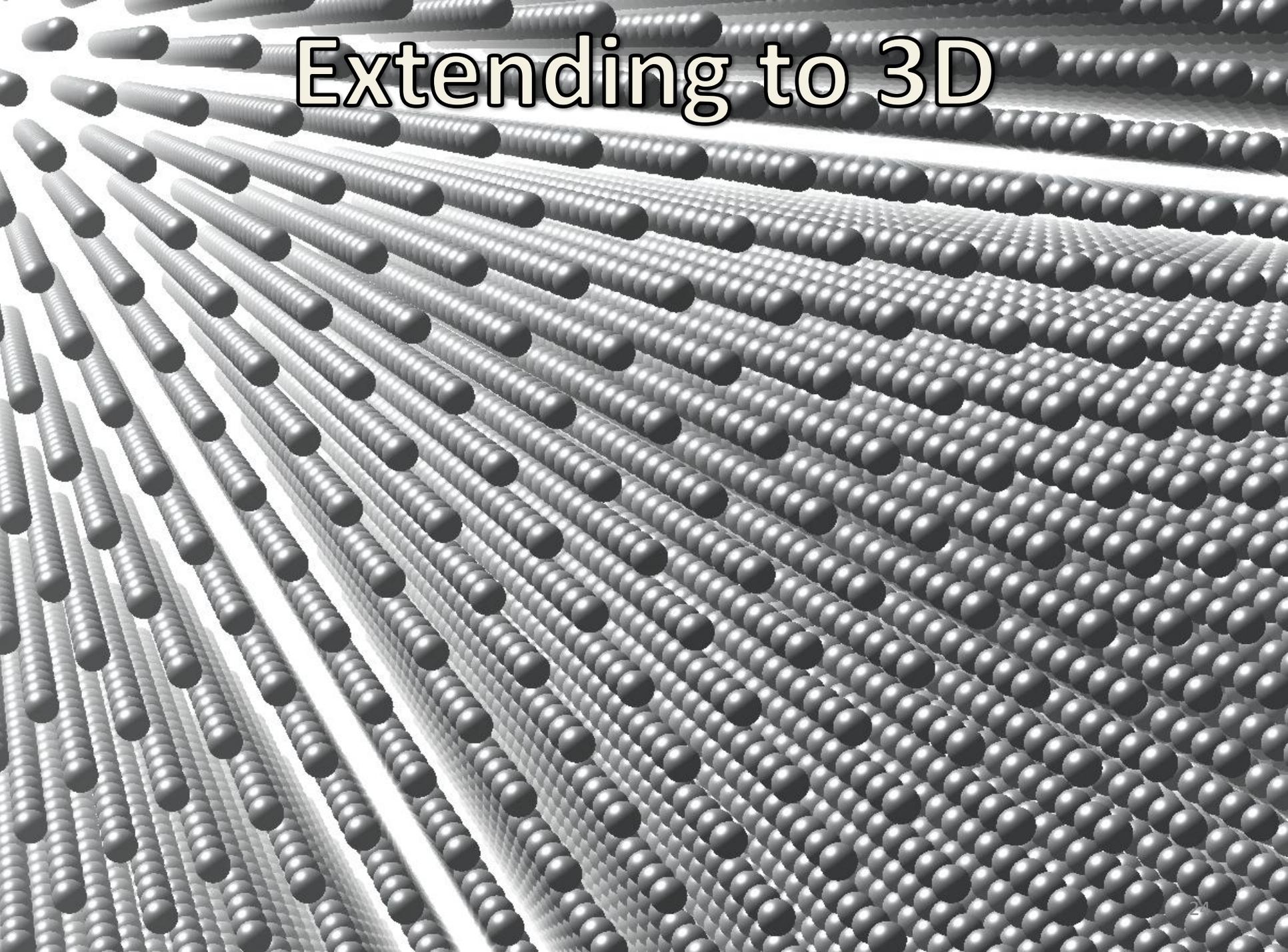


# 2D Solutions for p-orbitals



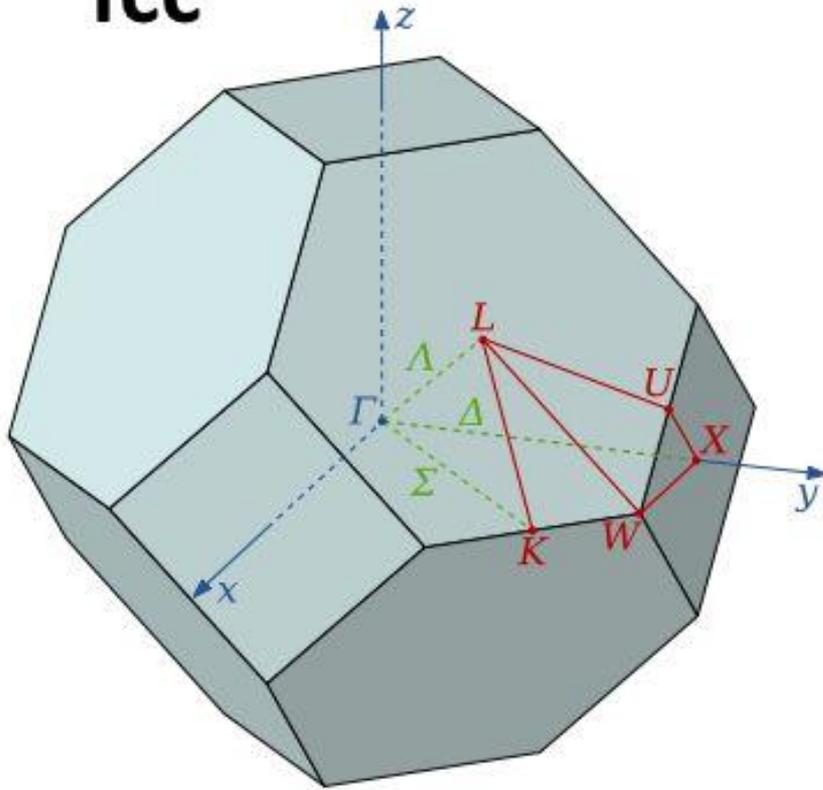


# Extending to 3D



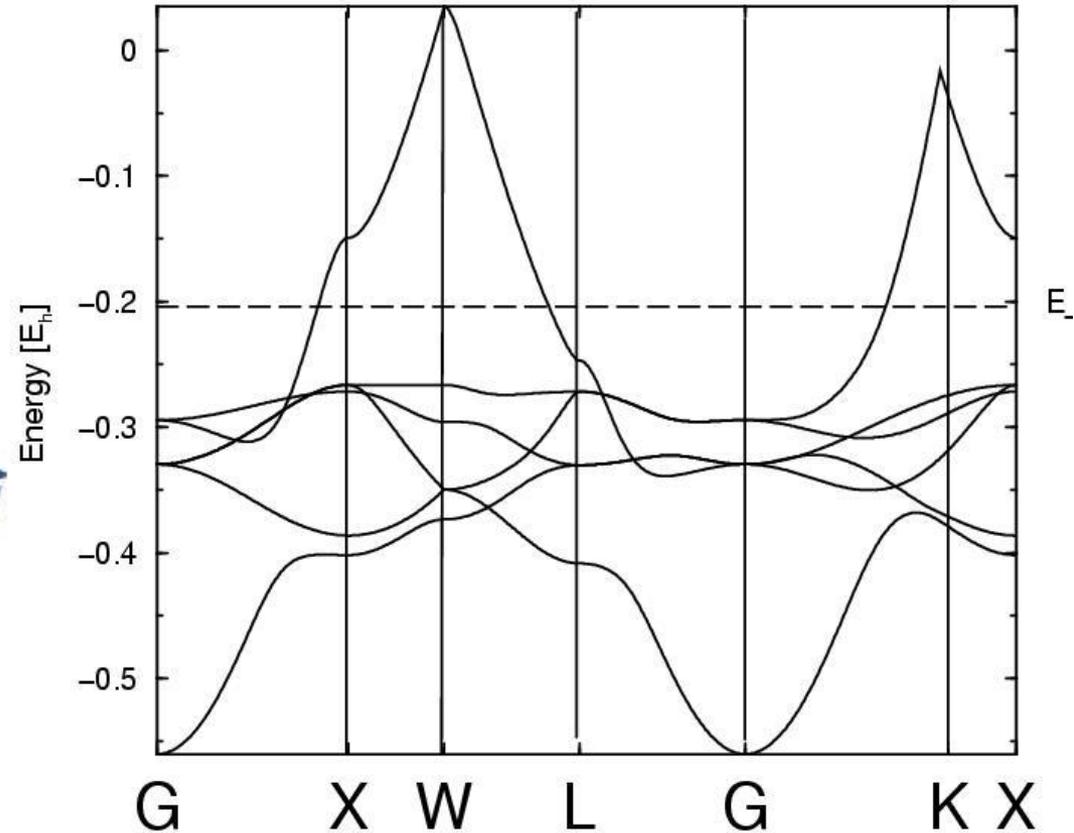
# 1.5 Brillouin Zone in 3D

**fcc**



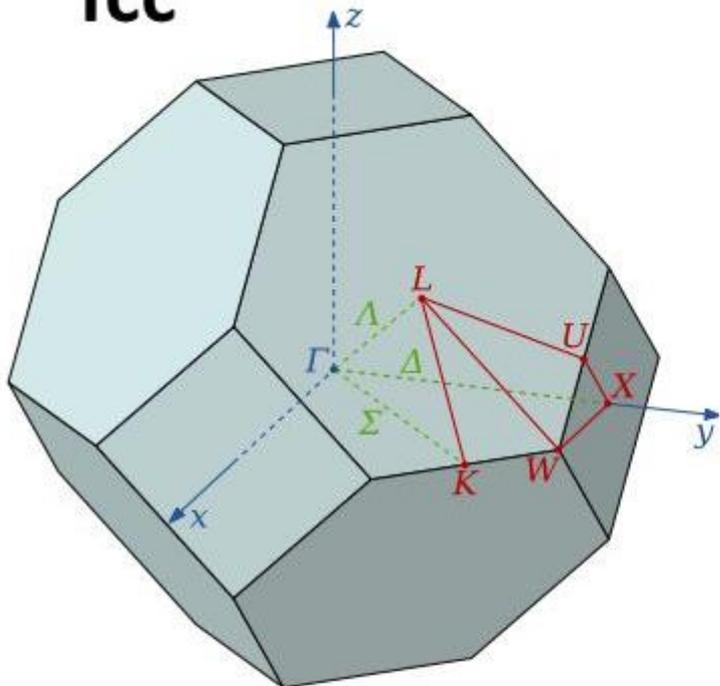
**Cu**

LDA band structure

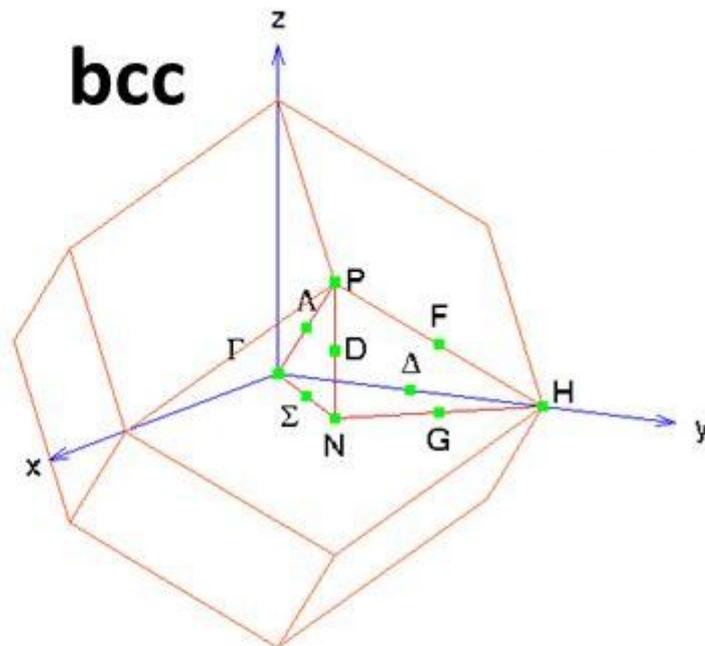


# 1.5 Brillouin Zone of other lattices

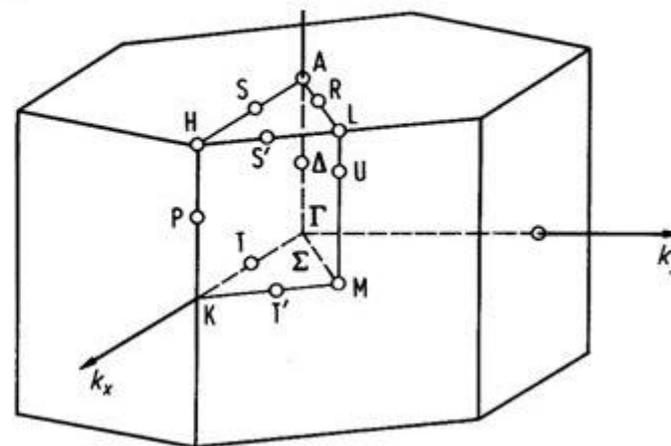
fcc



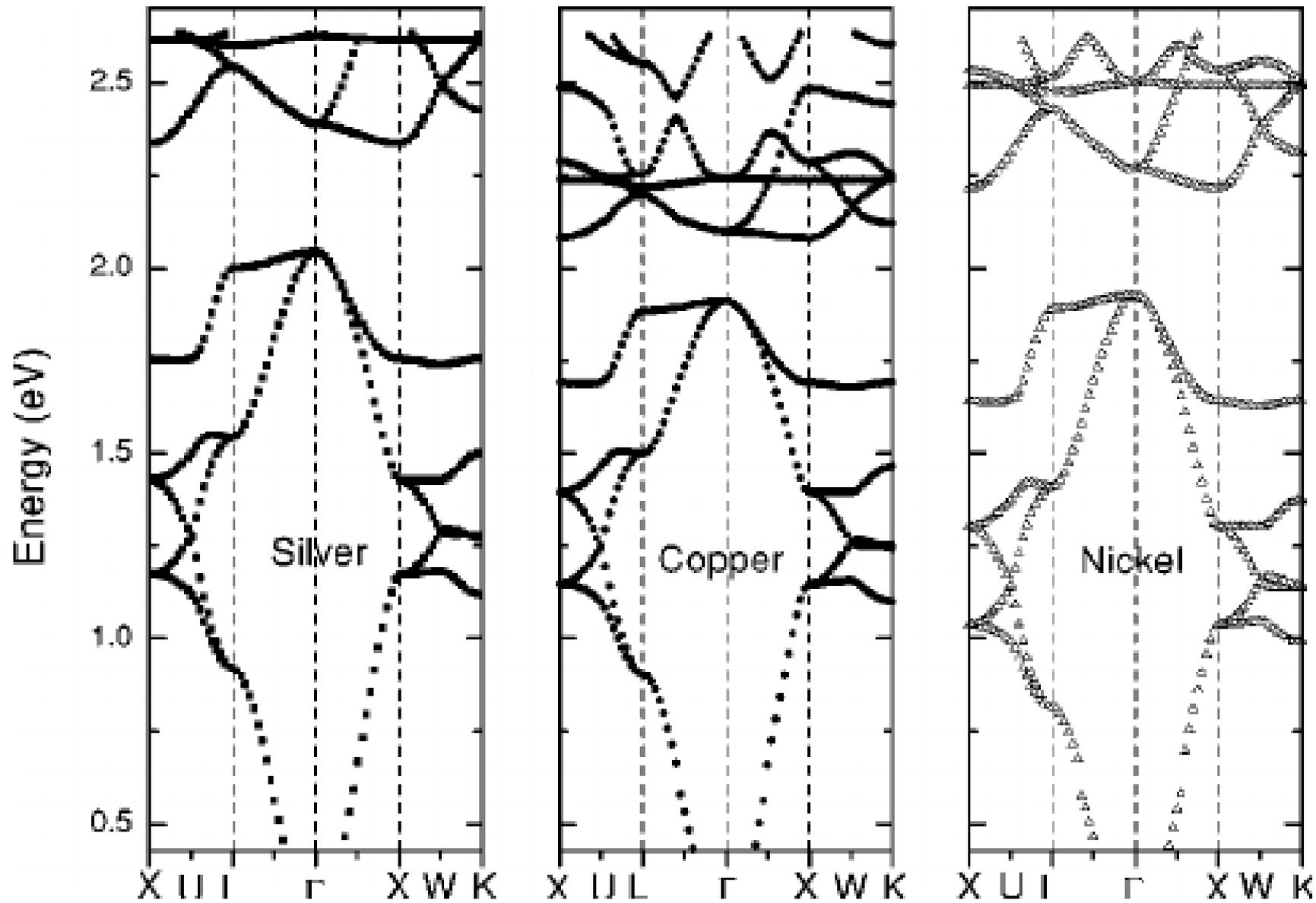
bcc

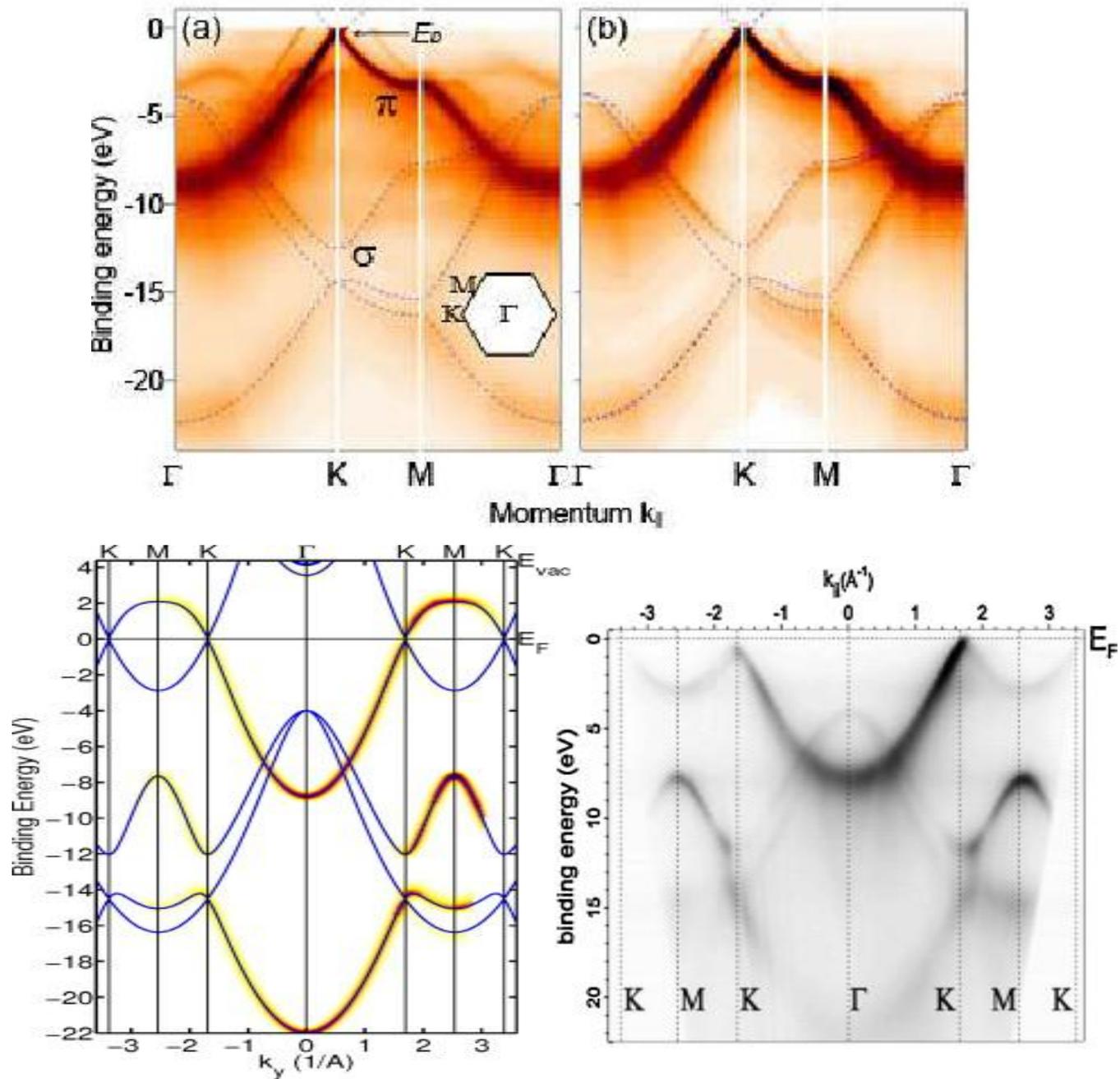


hcp

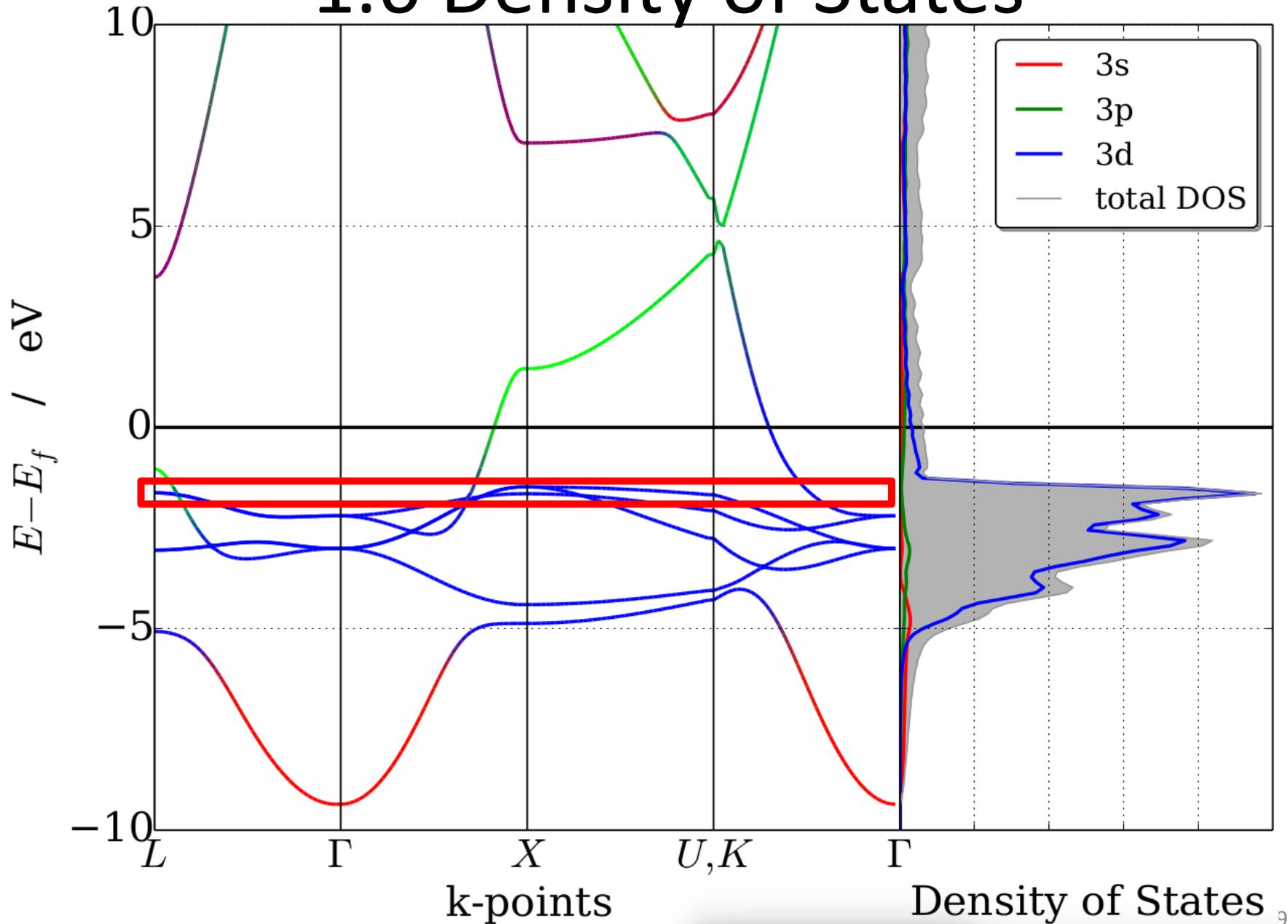


•The BZ reflects lattice symmetry

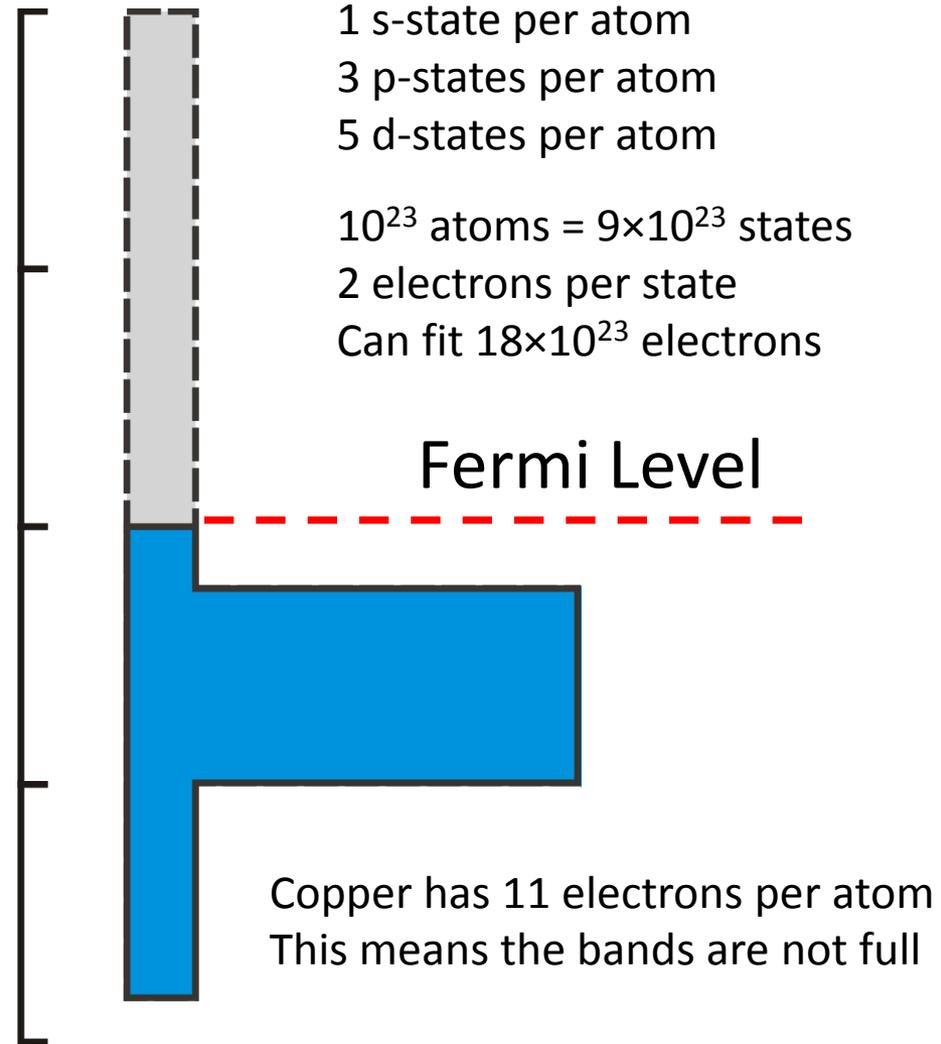
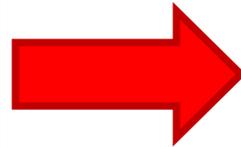
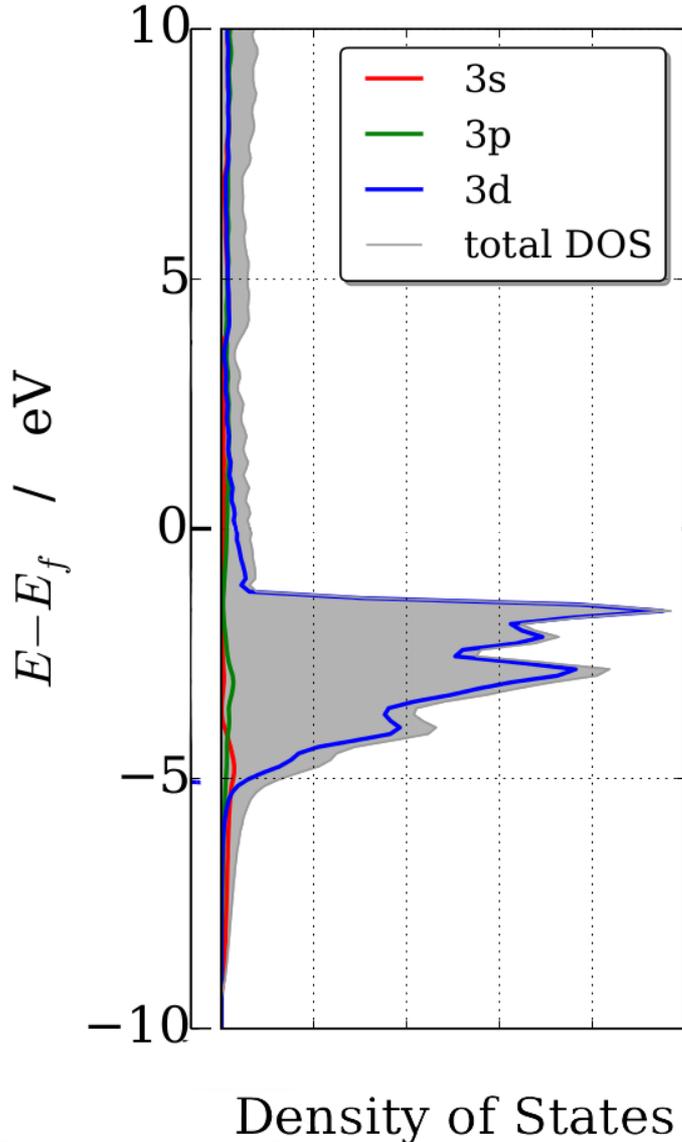




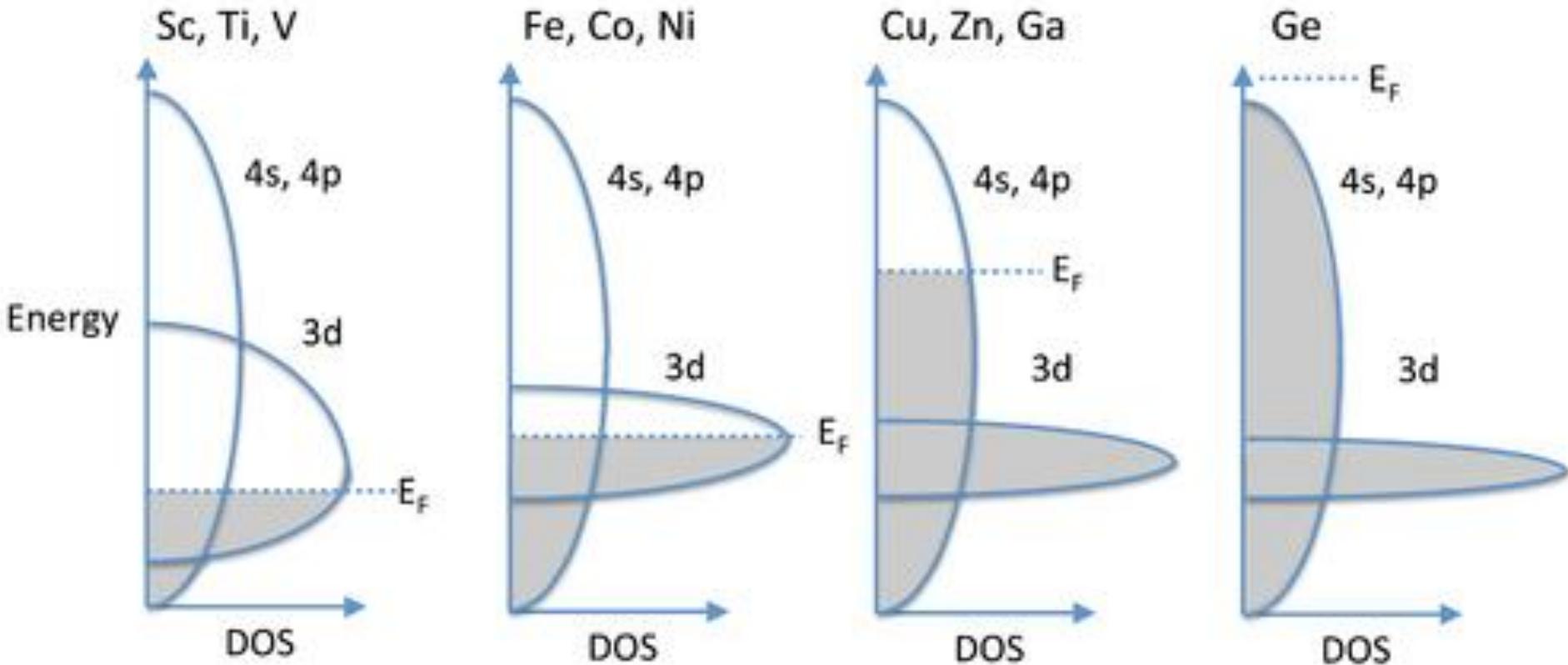
# 1.6 Density of States



# 1.6 Filling with electrons

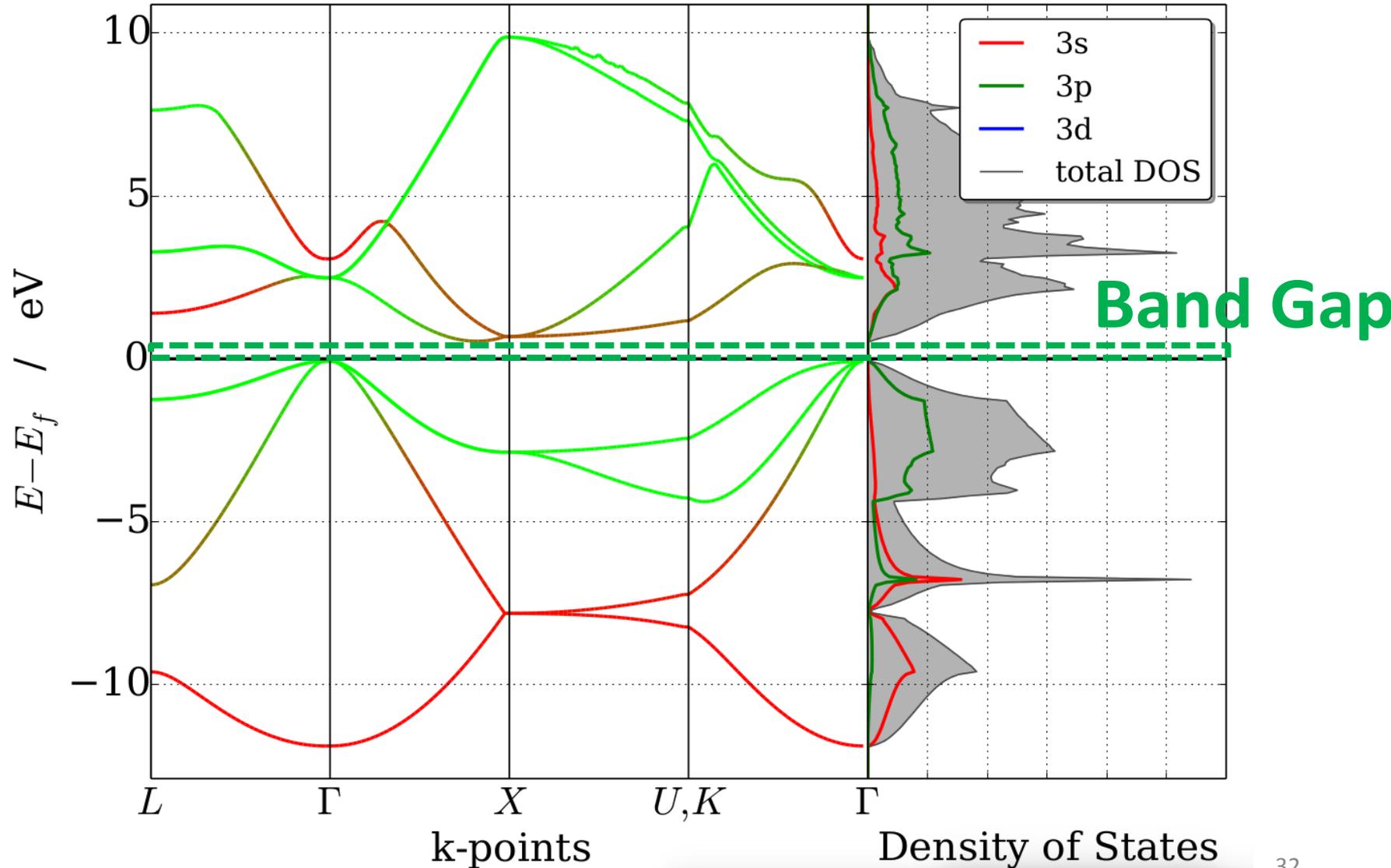


# 1.6 Band filling of the transition metals

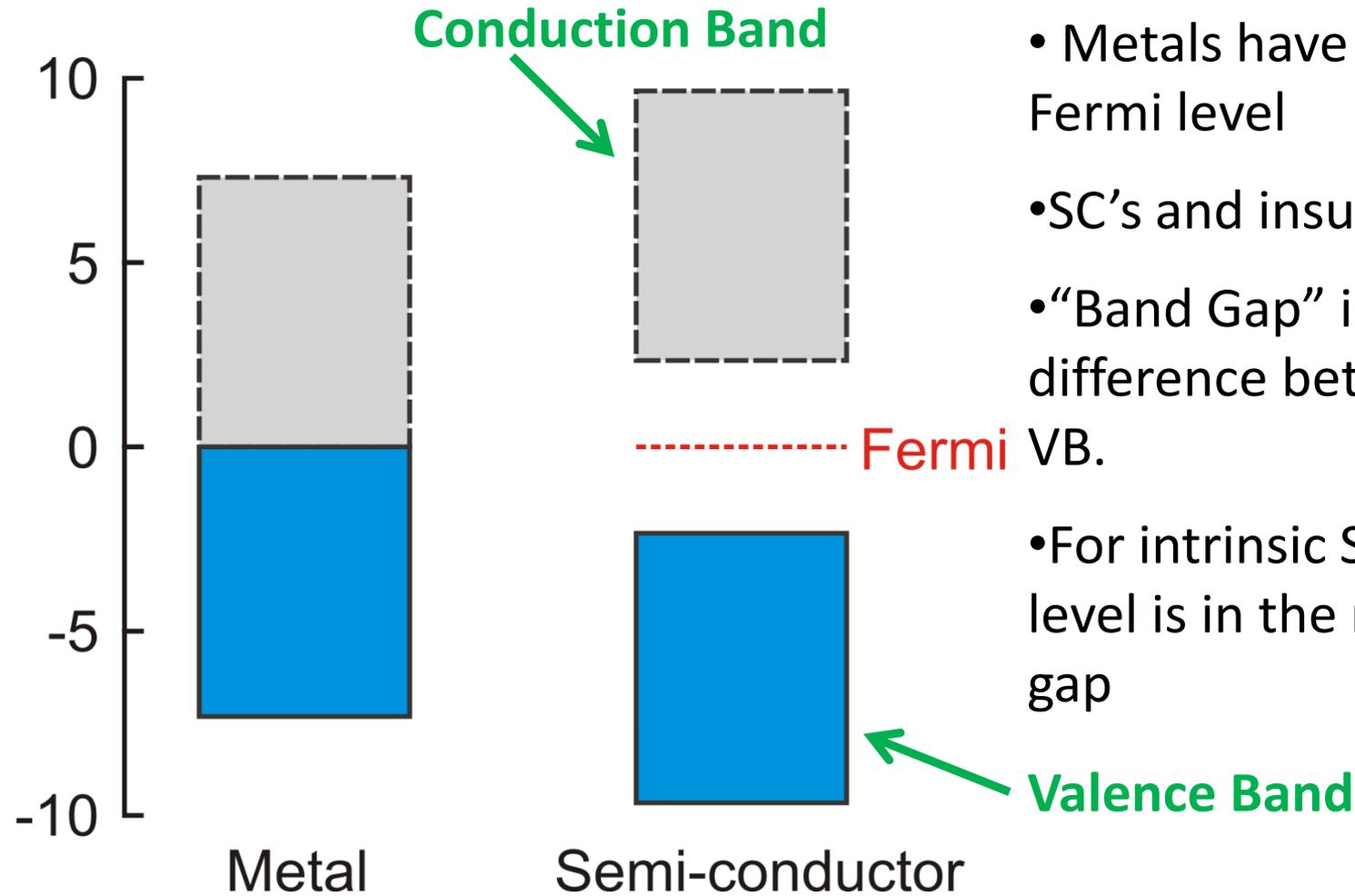


# 1.7 Semiconductors and band gaps

Bands diagram of silicon

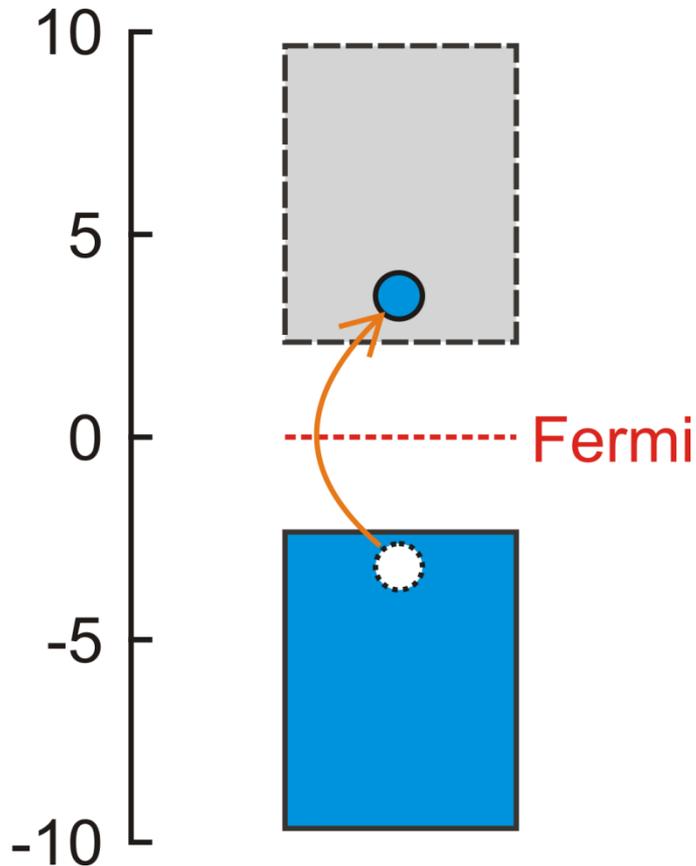


# 1.7 Metals vs Semiconductors



- Metals have states at Fermi level
- SC's and insulators do not
- "Band Gap" is the difference between CB and VB.
- For intrinsic SC's the Fermi level is in the middle of the gap

# 1.7 Conduction

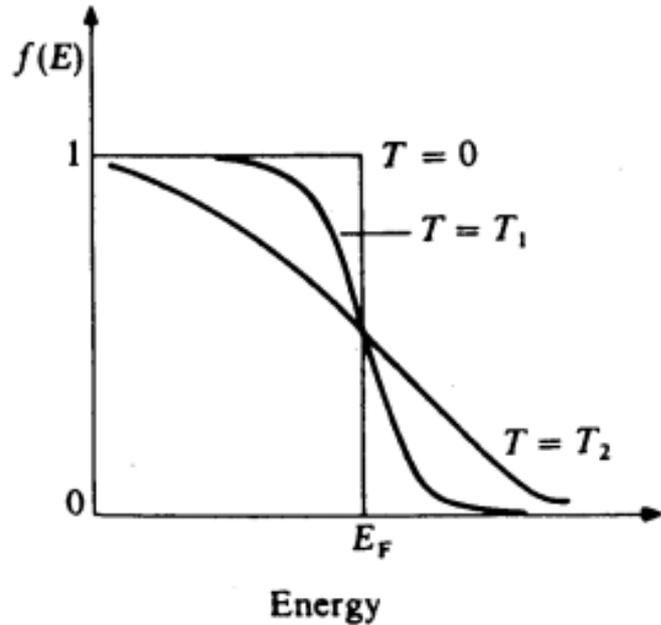


- Elect. Conduction occurs via excitation of  $e^-$  into unoccup. States
- This generates electrons in the CB and holes (i.e. electron vacancies) in the VB.
- Both lead to conductivity
- Average thermal energy ( $kT$ ) at room temp. is 0.013 eV
- Only the electrons within a few  $kT$  of the Fermi level can contribute to conduction

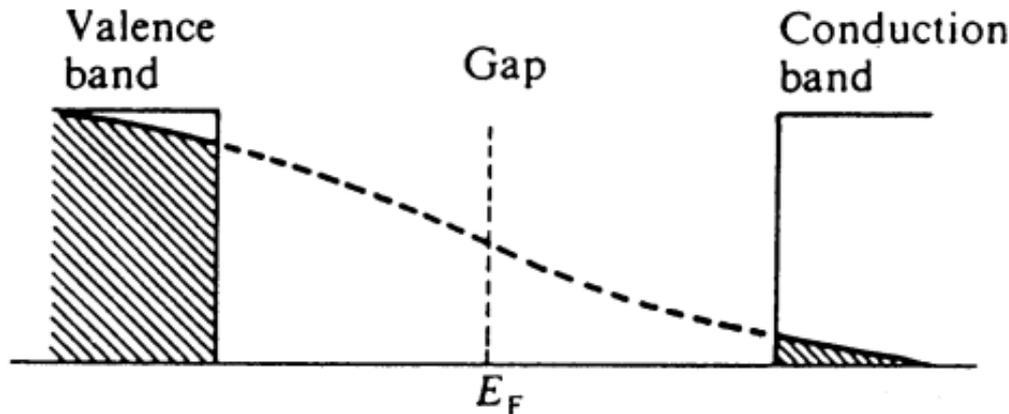
# 1.7 Fermi-Dirac Function

$$f_{\text{Fermi-Dirac}}(E) = \frac{1}{e^{(E - E_F)/kT} + 1}$$

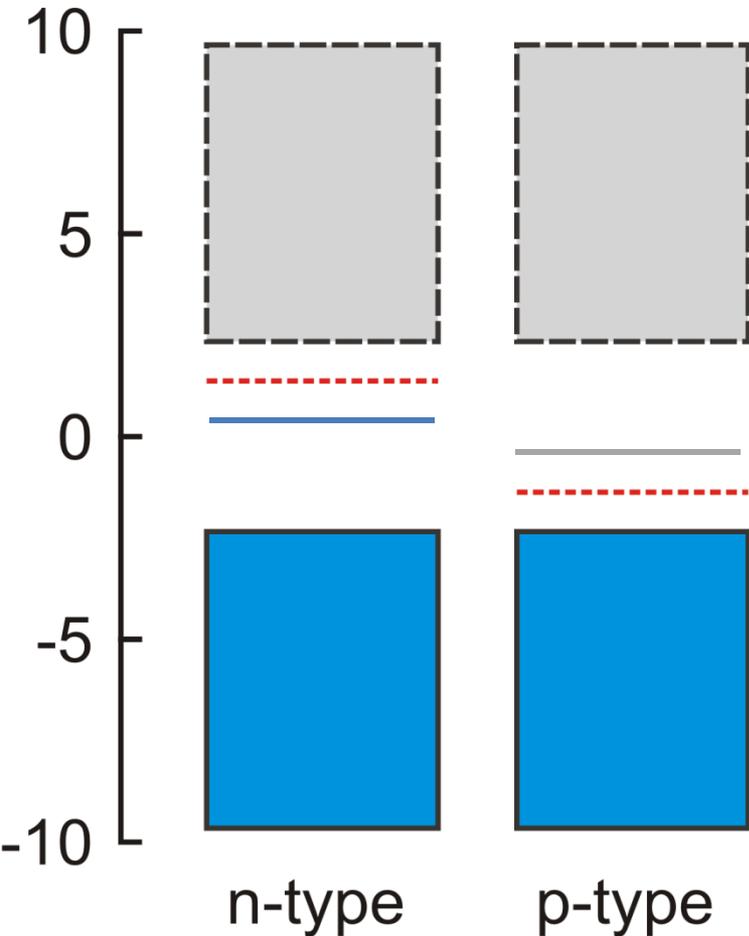
# 1.7 Carrier Concentrations



- At 0 K, the electron population distribution is a step function
- As  $T$  increases, CB states become filled and VB states become unoccupied.
- Fermi function is used to calculate this.
- Fermi function can be used to calculate charge carrier concentrations in semiconductors.



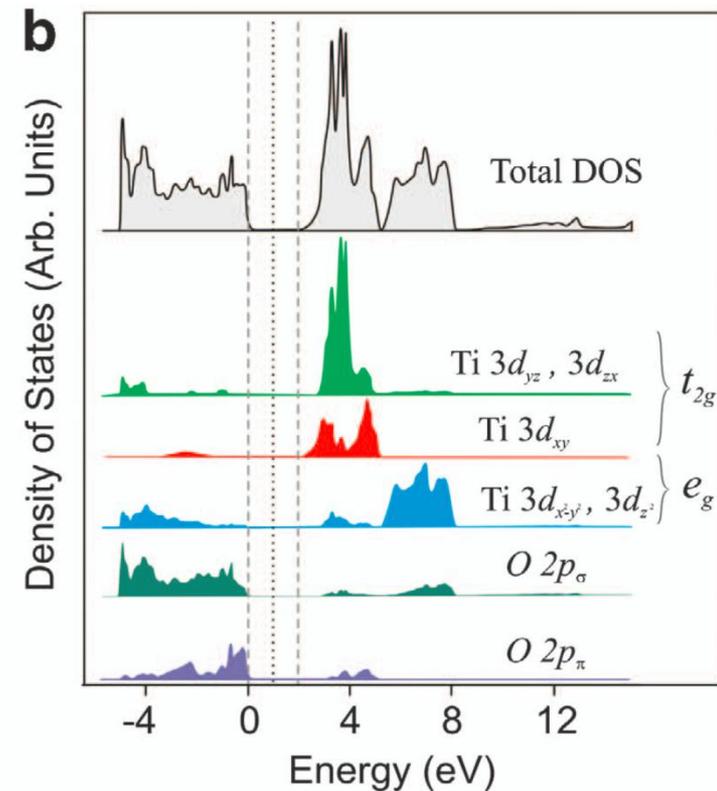
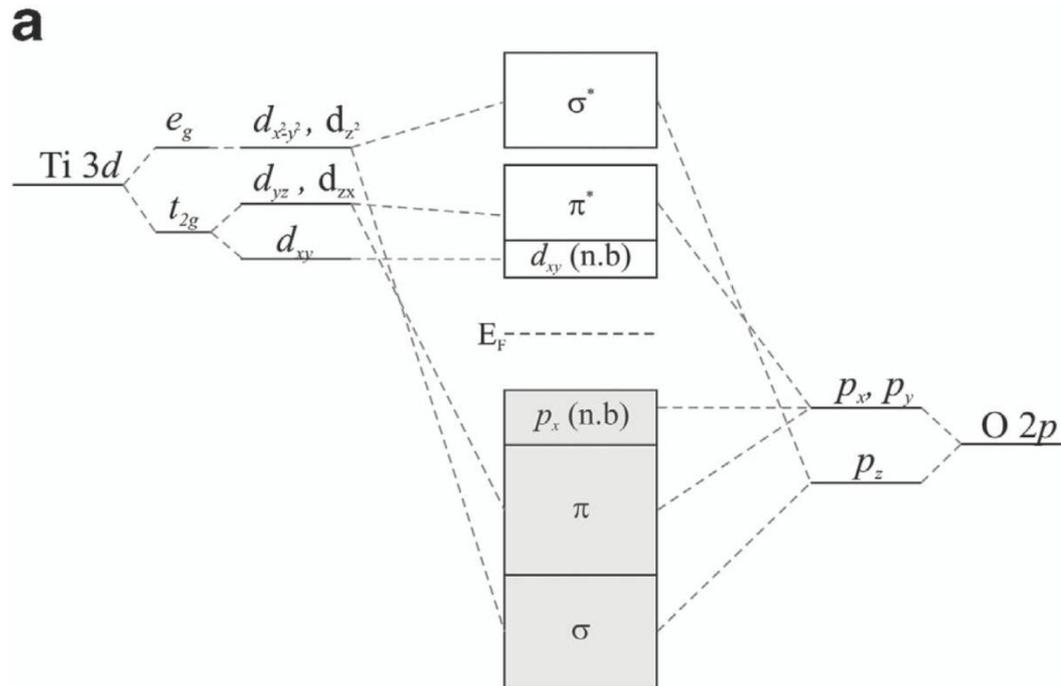
# 1.8 Doping Semiconductors



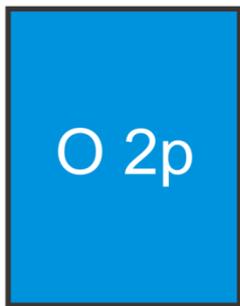
- Impurities can be added to SC's to increase conductivity
- N-type dopants add filled states close to the CB
- P-type dopants add empty states close to the VB
- N-type shifts Fermi toward CB
- P-type shifts Fermi toward VB
- In Si, P is an n-type dopant, B is a P-type dopant.

# 1.8 Doping with defects

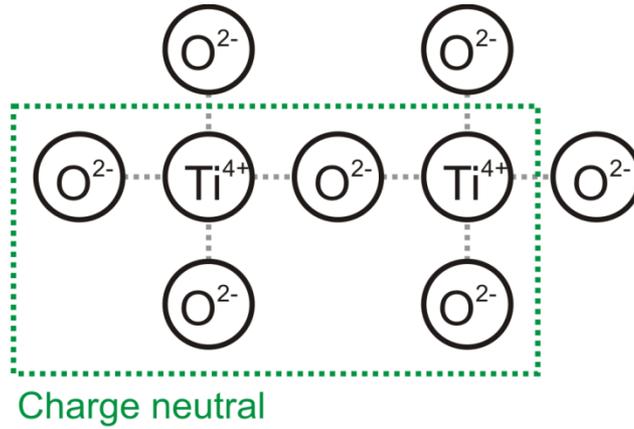
- Consider  $d^0$  oxides
- d-orbitals are empty, O 2p orbitals are full
- Examples:  $\text{TiO}_2$ ,  $\text{MoO}_3$ ,  $\text{V}_2\text{O}_5$ ,  $\text{WO}_3$



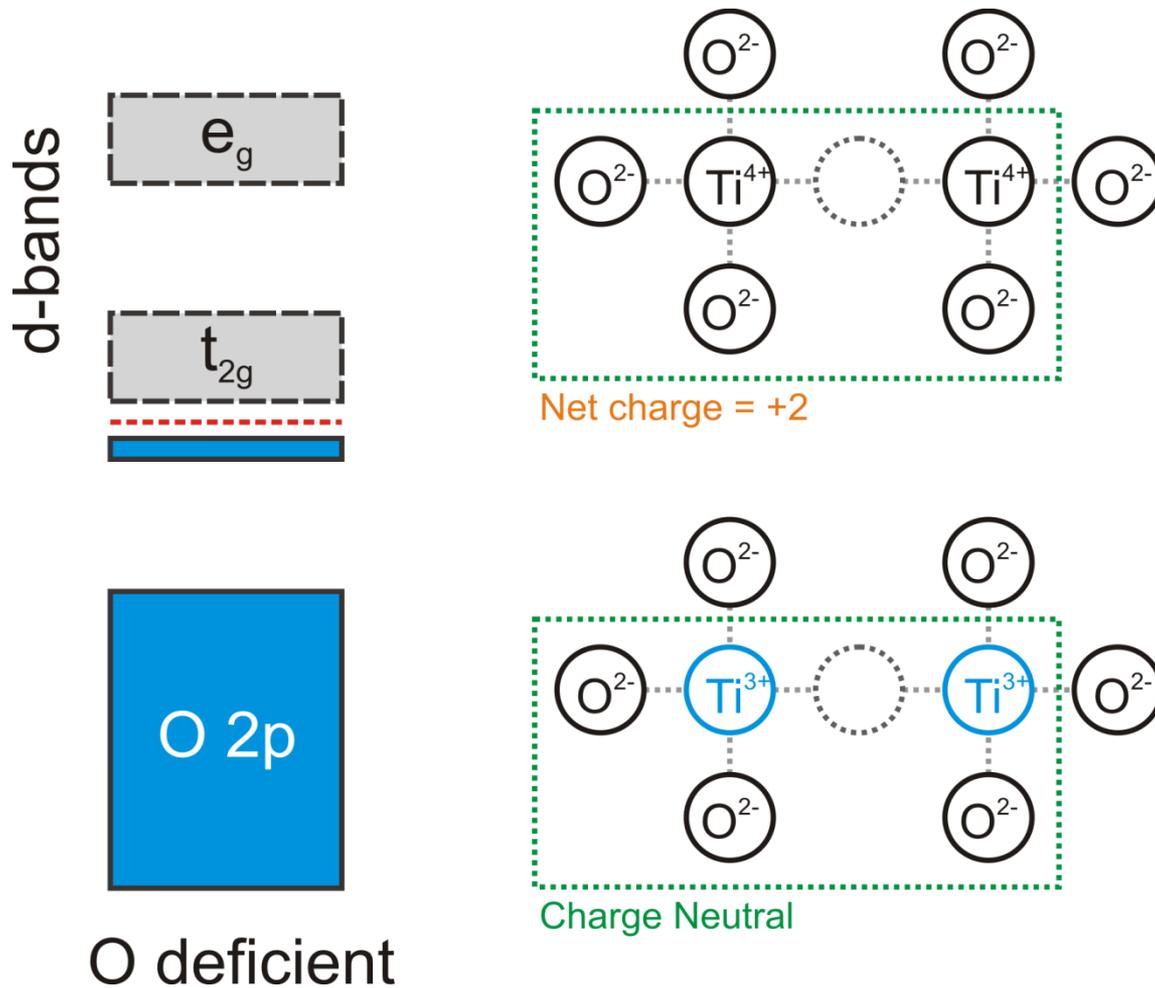
d-bands



Defect free



- Perfect crystal has Fermi at mid-gap



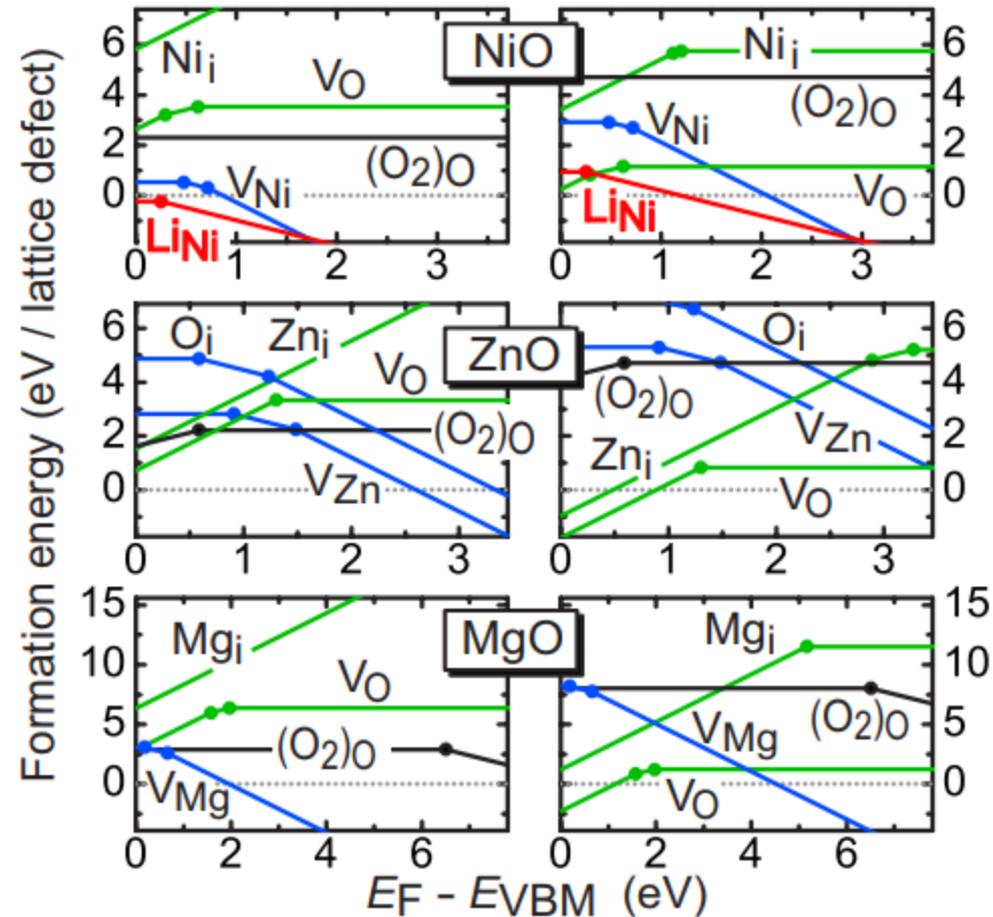
- O vacancy requires two  $Ti^{4+}$  to become  $Ti^{3+}$
- Therefore, Ti 3d band not empty
- Electrons on  $Ti^{3+}$  sites are n-type dopants

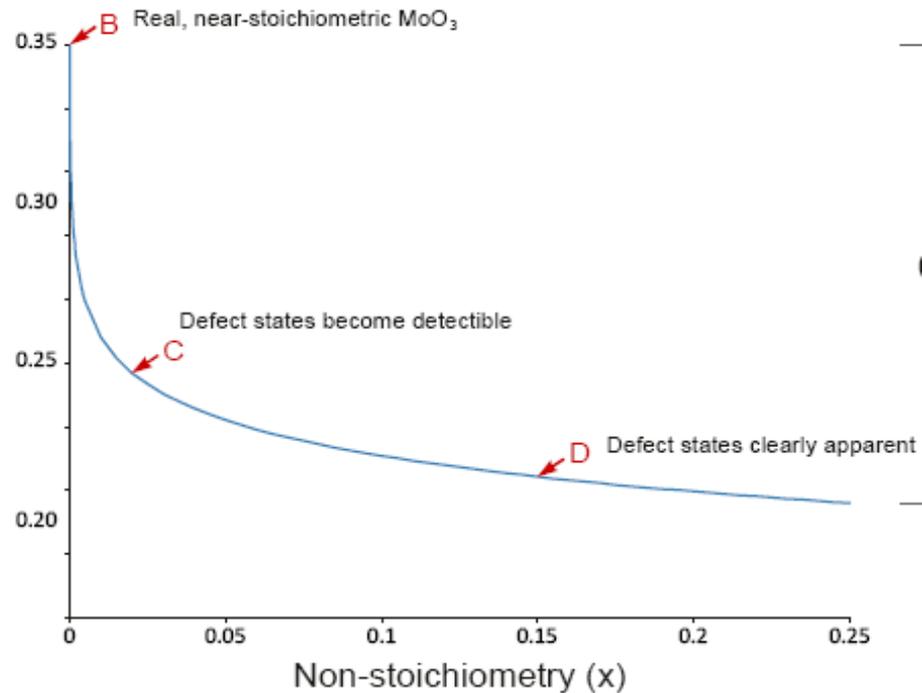
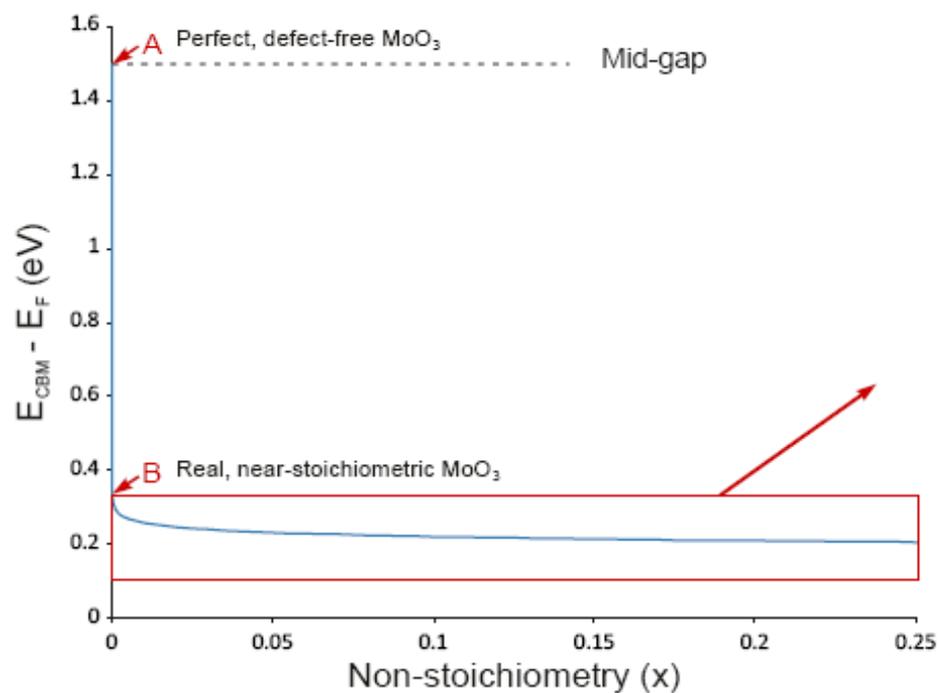
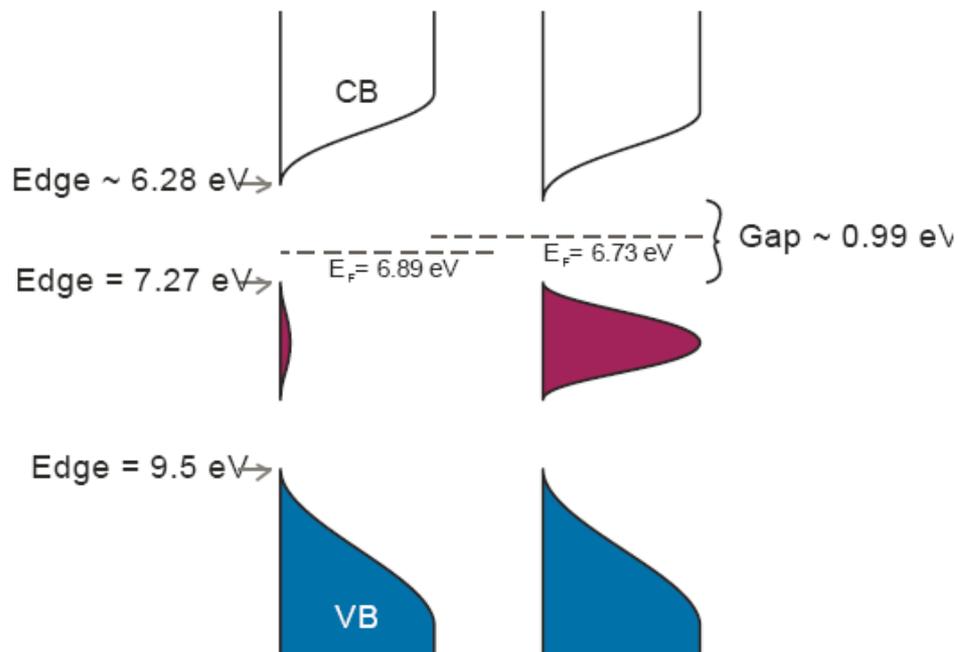
# Doping with defects

- Due to entropy, there is always some finite concentration of defects.
- Therefore, oxides are always intrinsically “doped”
- Type of doping depends on thermodynamic stability of defects
- O vacancies are n-type dopants
- Metal vacancies are p-type dopants
- O interstitials are p-type dopants
- Metal interstitials are n-type dopants
- Doping degree changes with temperature and O-partial pressure

# Doping with defects

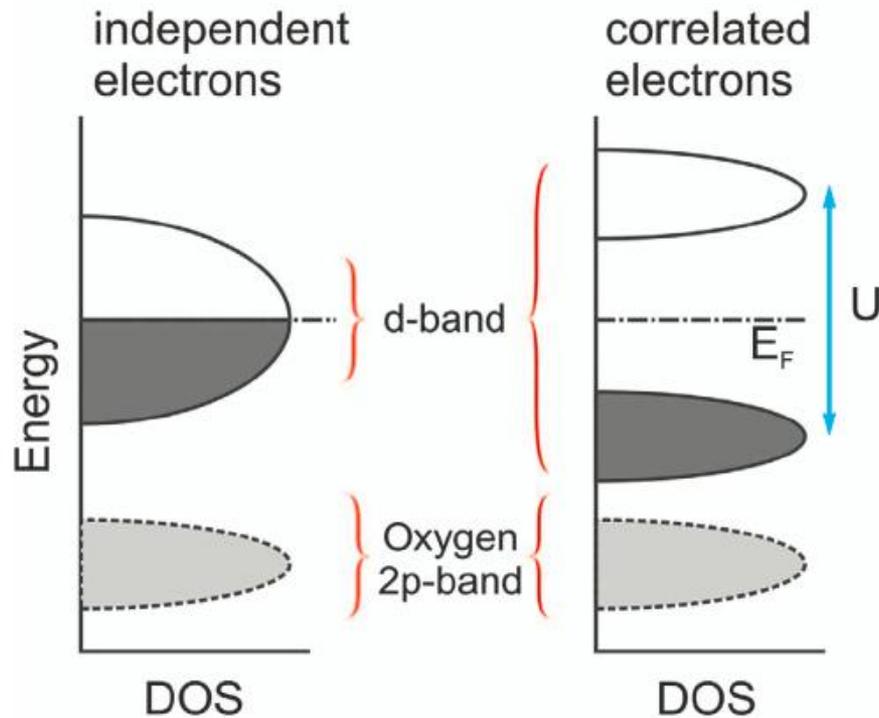
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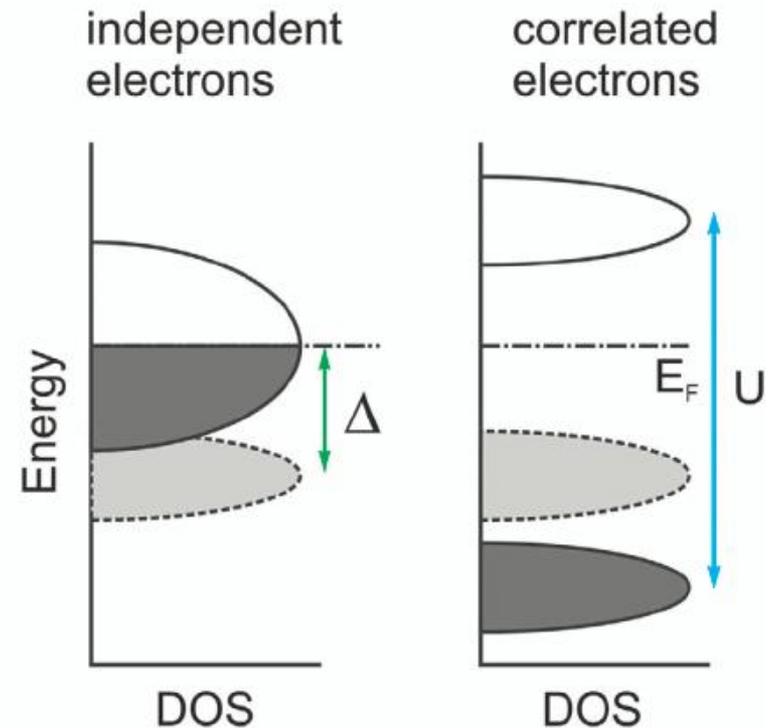


# Non-classical Insulators

**a** Mott-Hubbard Insulator

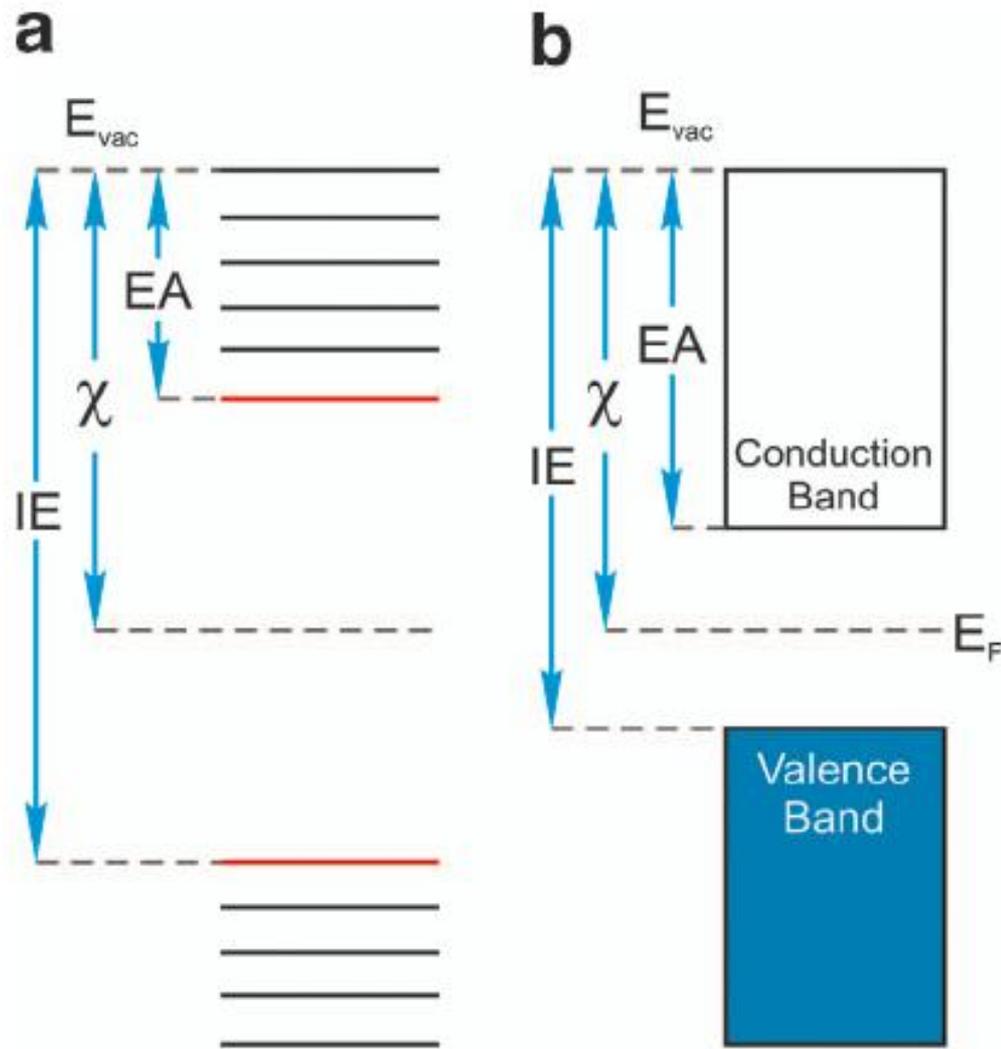


**b** Charge-Transfer Insulator



# Part 2: The Fermi Level, Work Function and Band Bending

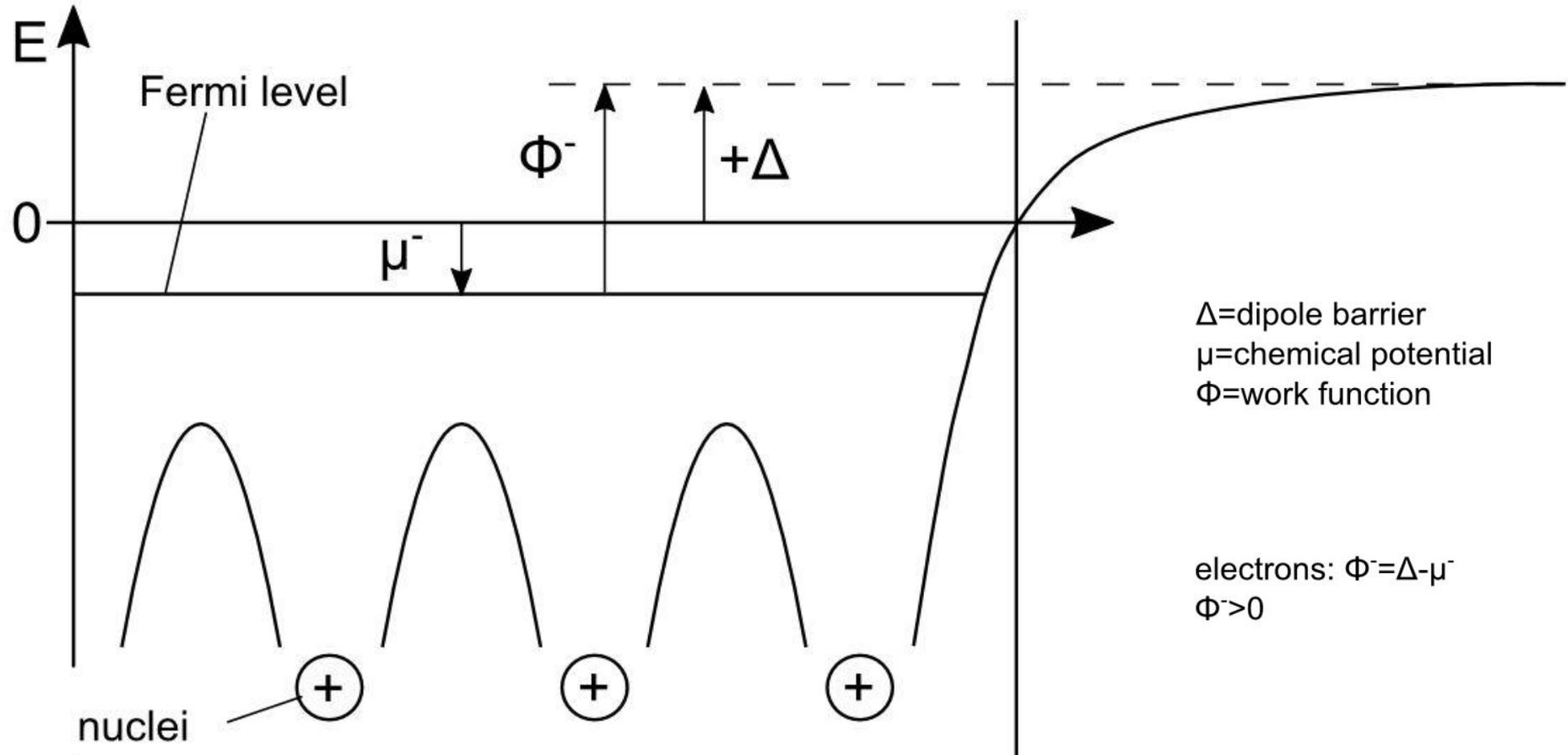
## 2.1 Fermi level is related to electronegativity



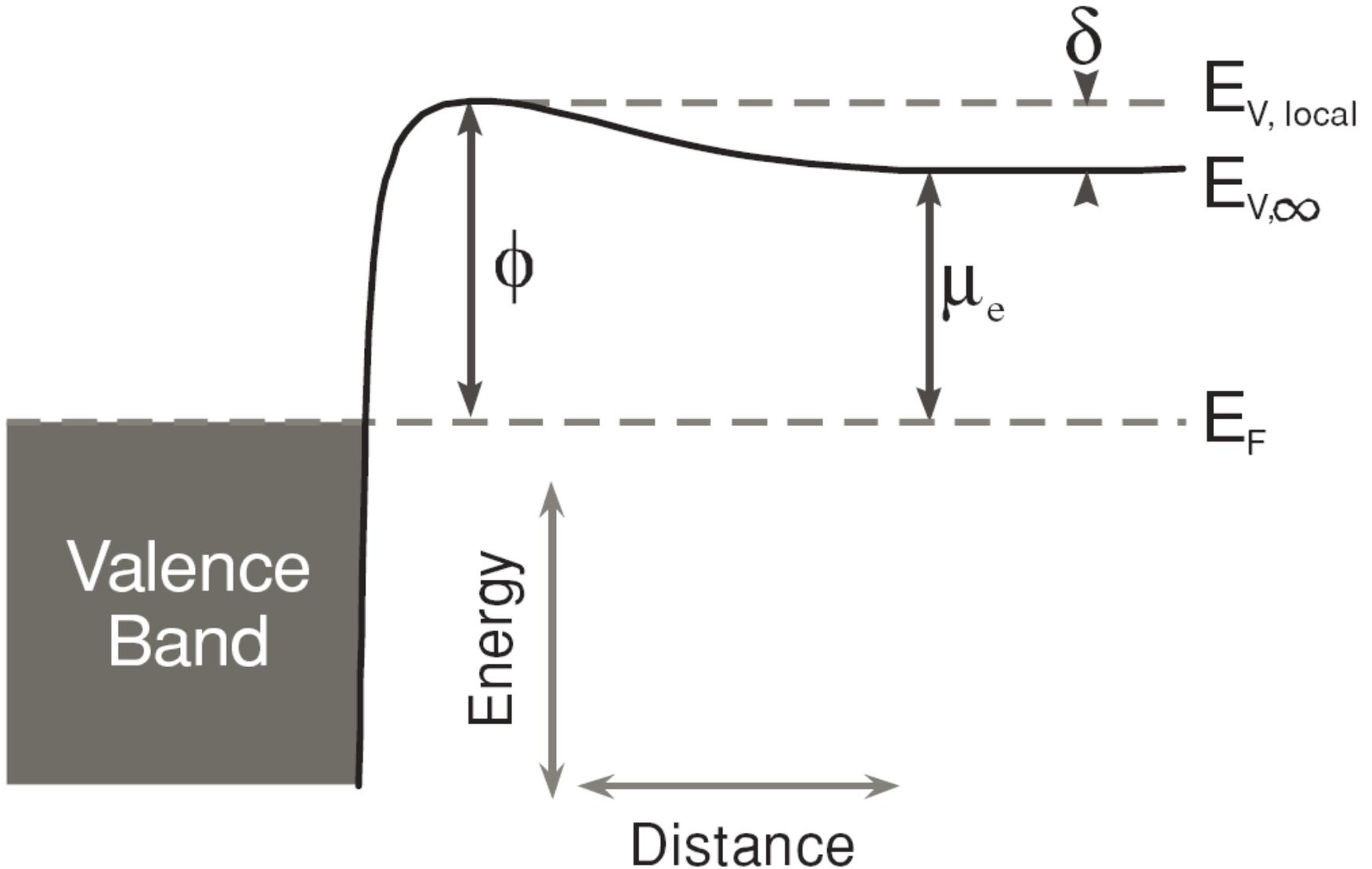
# 2.1 Work Function

- Work function is the minimum energy required to remove an electron from a solid.
- It is the energy difference between the “local” vacuum level and the Fermi level.
- Local is to be distinguished from absolute vacuum level
- Work function is a characteristic of surfaces.
- It is related to the Fermi level (which is a property of solids), but contains an additional term that is a property of the surface
- This term is the surface dipole.

# 2.1 Surface Dipole



# 2.1 Surface Dipole



## 2.1 Local vs. Absolute Vacuum Level

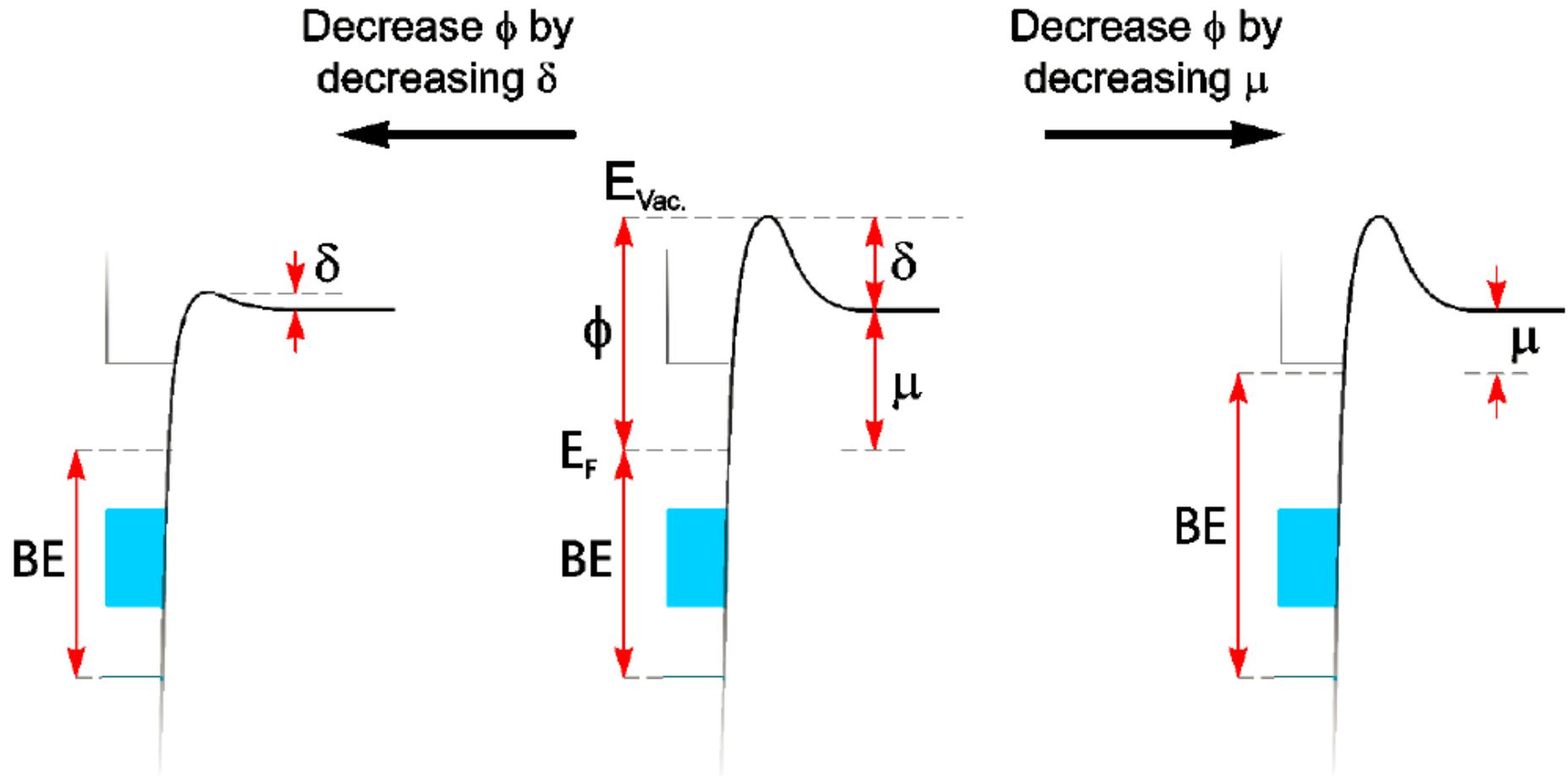
- Absolute vacuum is a universal constant, however, it is not directly measurable
- It is the rest energy of a test charge in absolute vacuum, at infinite distance away from all other particles.
- All we can measure is energy required to eject an electron.
- This is the sum of Fermi energy + surface dipole
- We cannot disentangle these two contributions
- Thus work function is not really Fermi energy or electron chemical potential, but it is the our best measurable approximation.

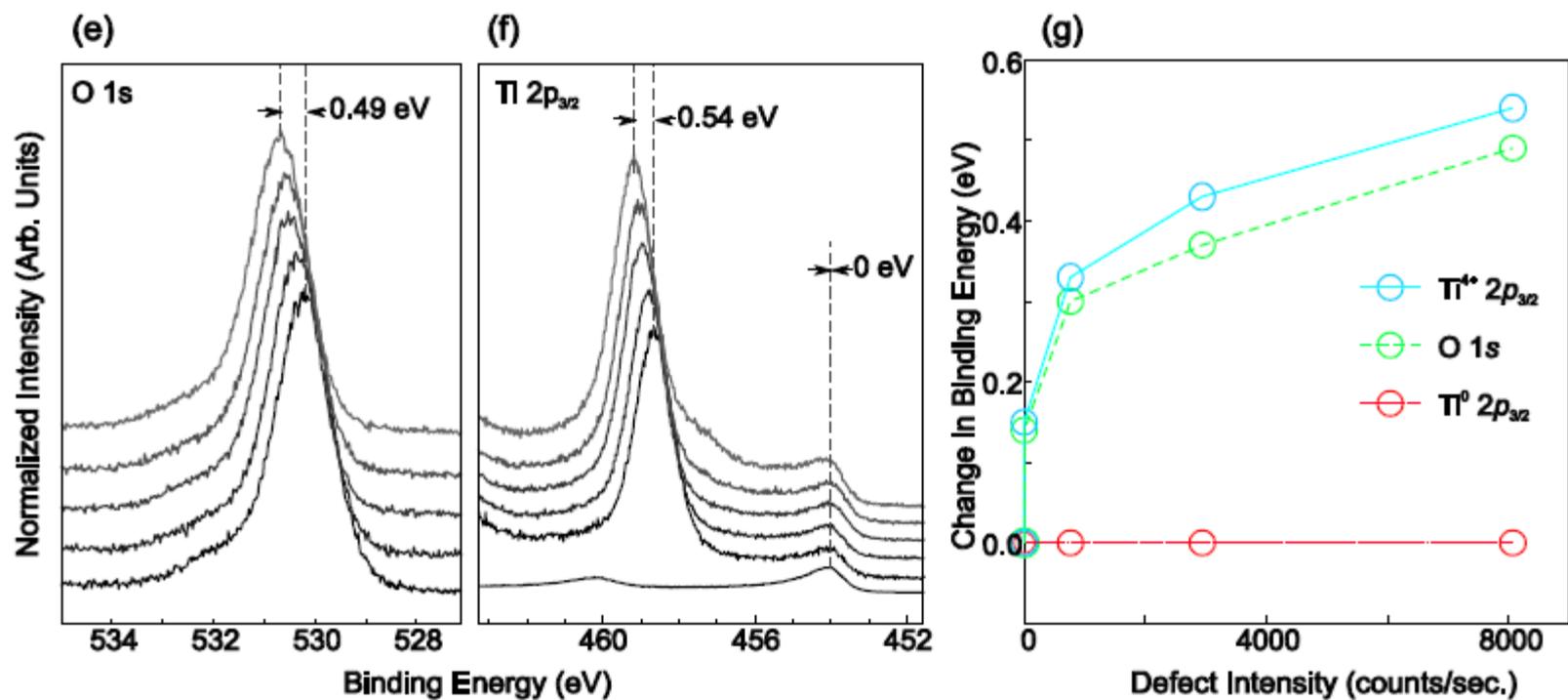
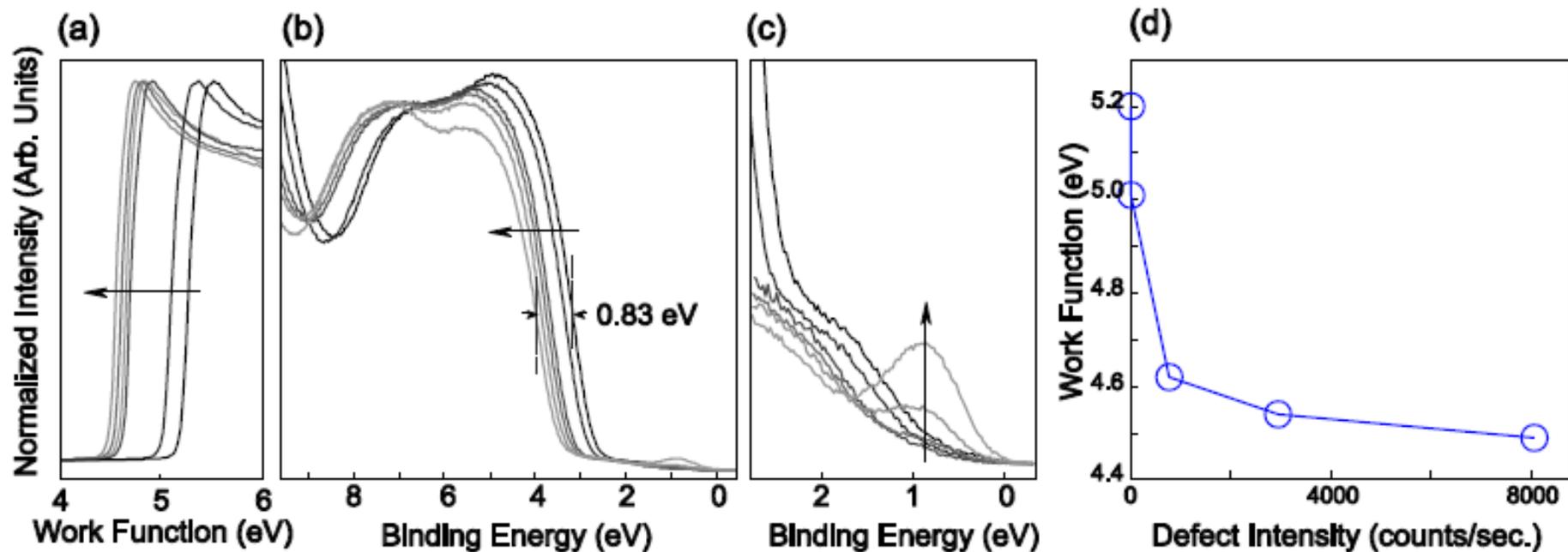
# 2.1 Work Function Anisotropy

- Different crystal faces have different surface dipoles due to differing atomic packing and resulting electron density
- Thus different crystal faces have different work functions.

	Work Function (eV)
Cu(111)	4.91
Cu(100)	4.81
Cu(110)	4.53

# 2.1 Work Function, Fermi Level and XPS Binding Energies

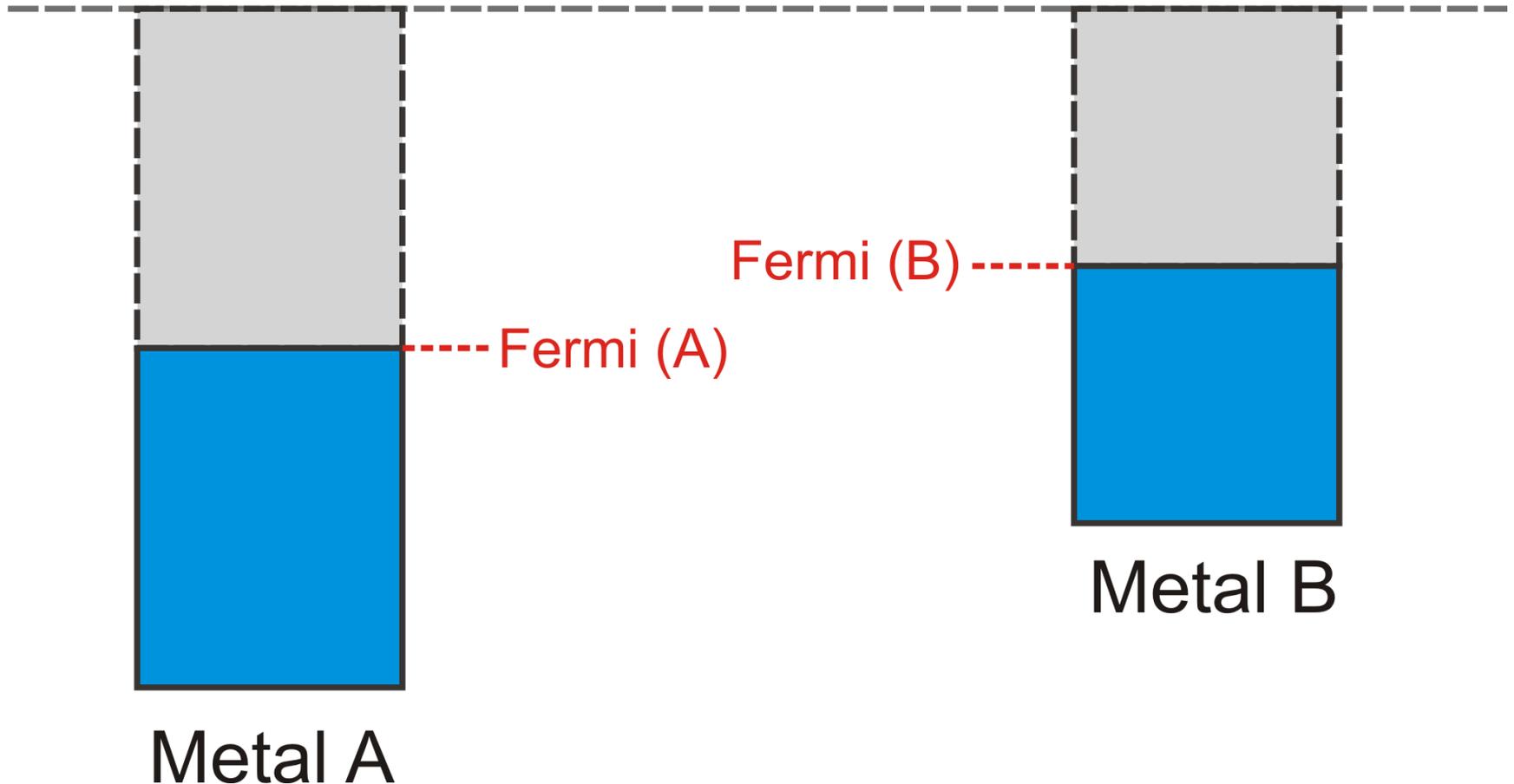




# 2.2 Equalizing Fermi Levels

Two metals with different Fermi levels at infinite distance

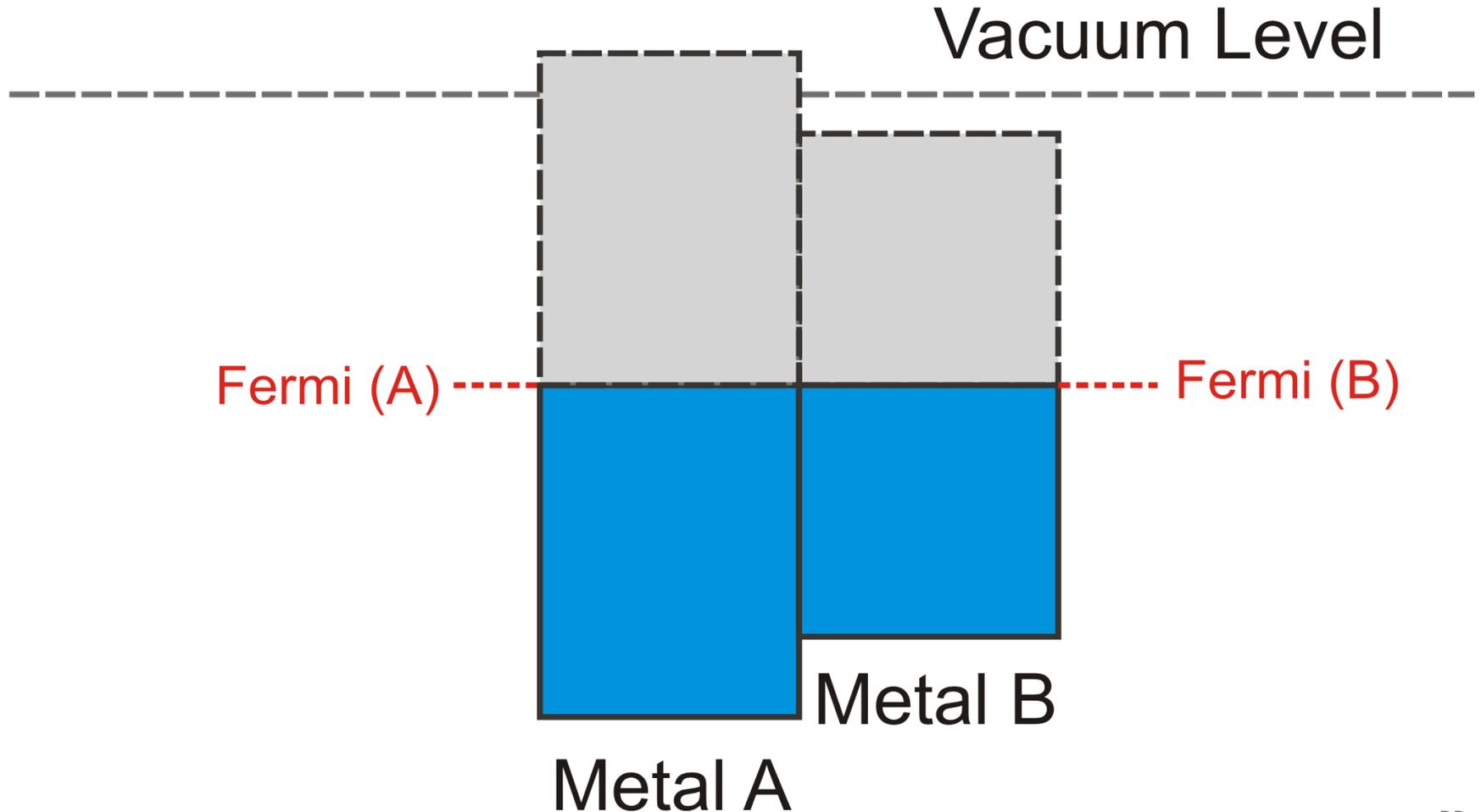
Vacuum Level



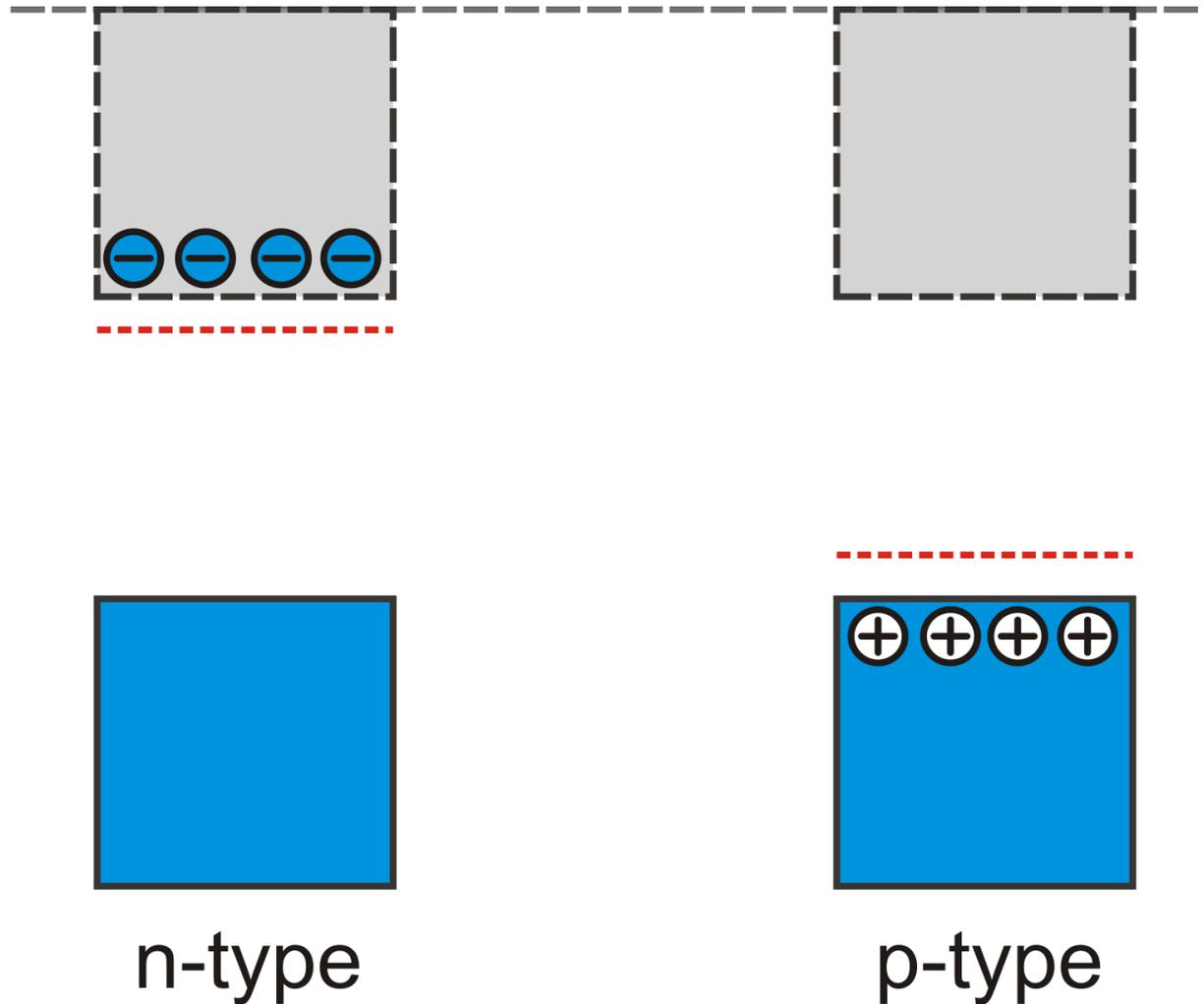
# 2.2 Equalizing Fermi Levels

Two metals with different Fermi levels in contact

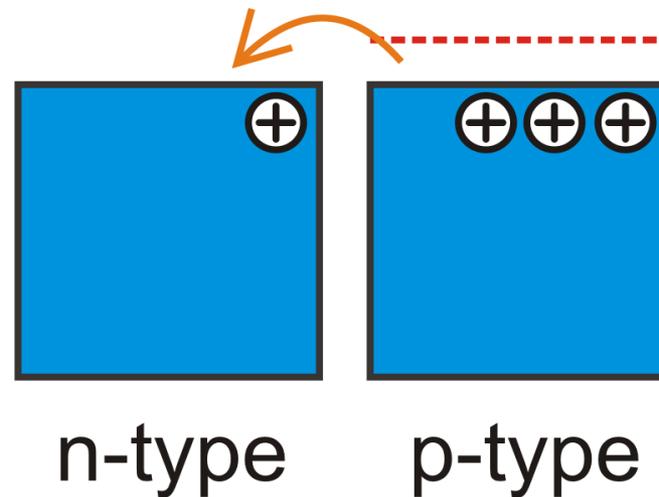
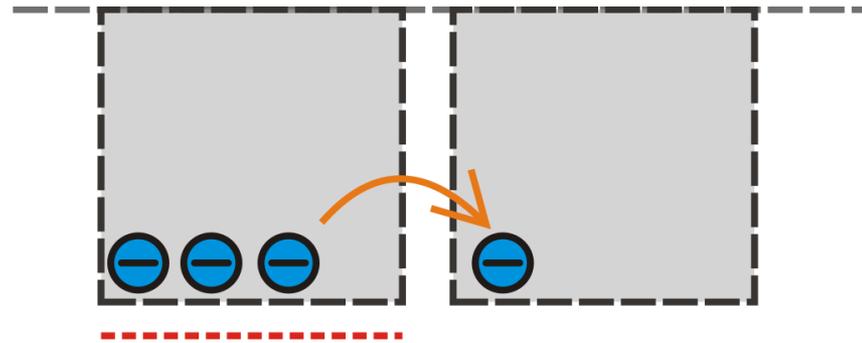
- Fermi levels equilibrate
- Electrons move from low work function to high work function



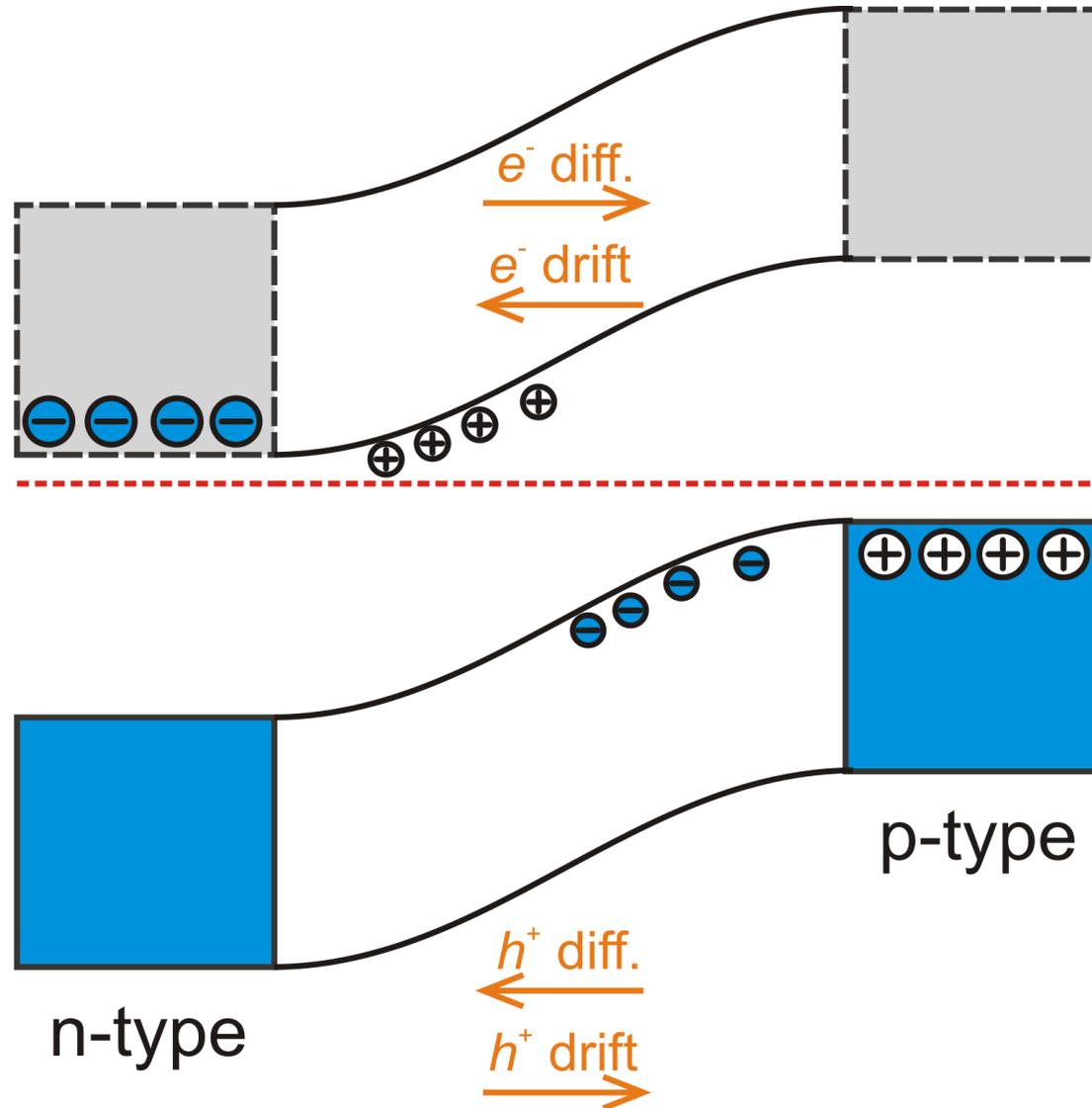
## 2.3 Band Bending at p-n Junction



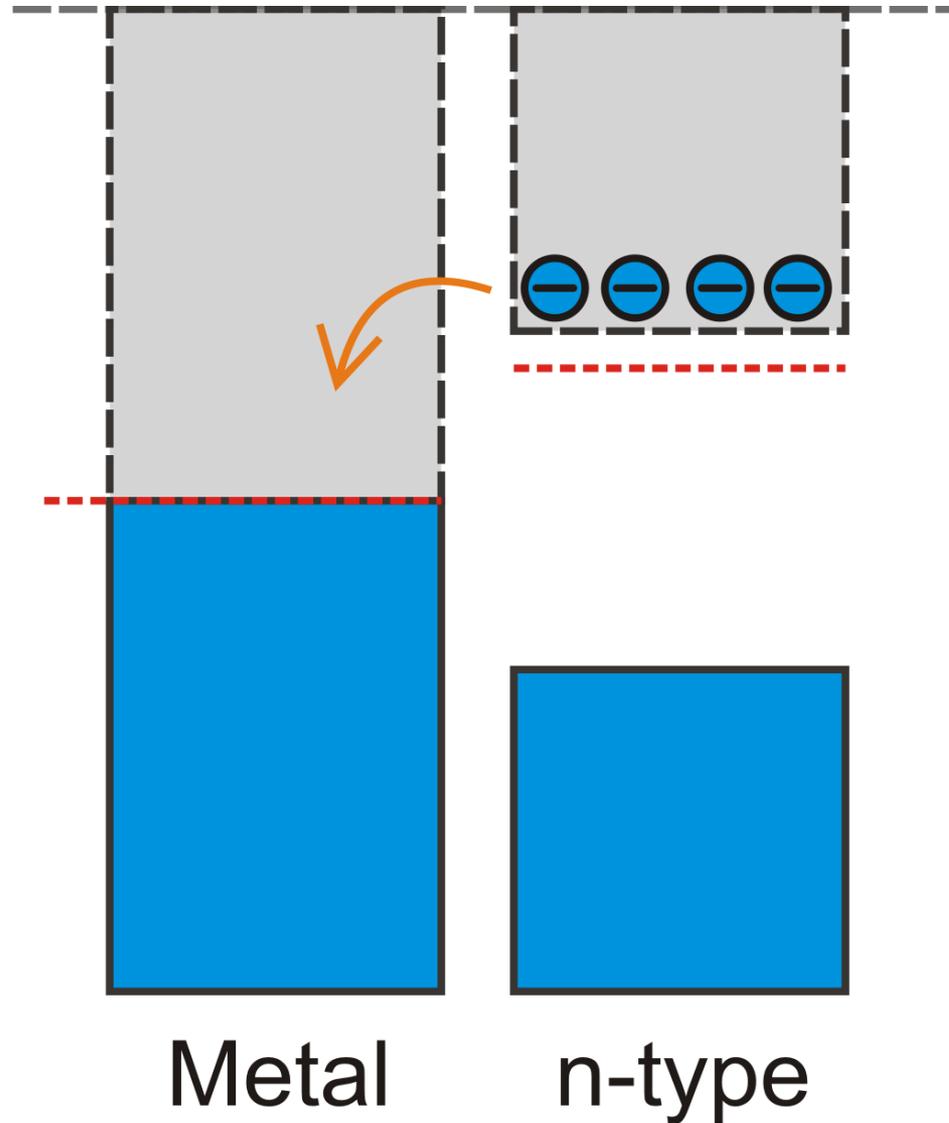
## 2.3 Band Bending at p-n Junction



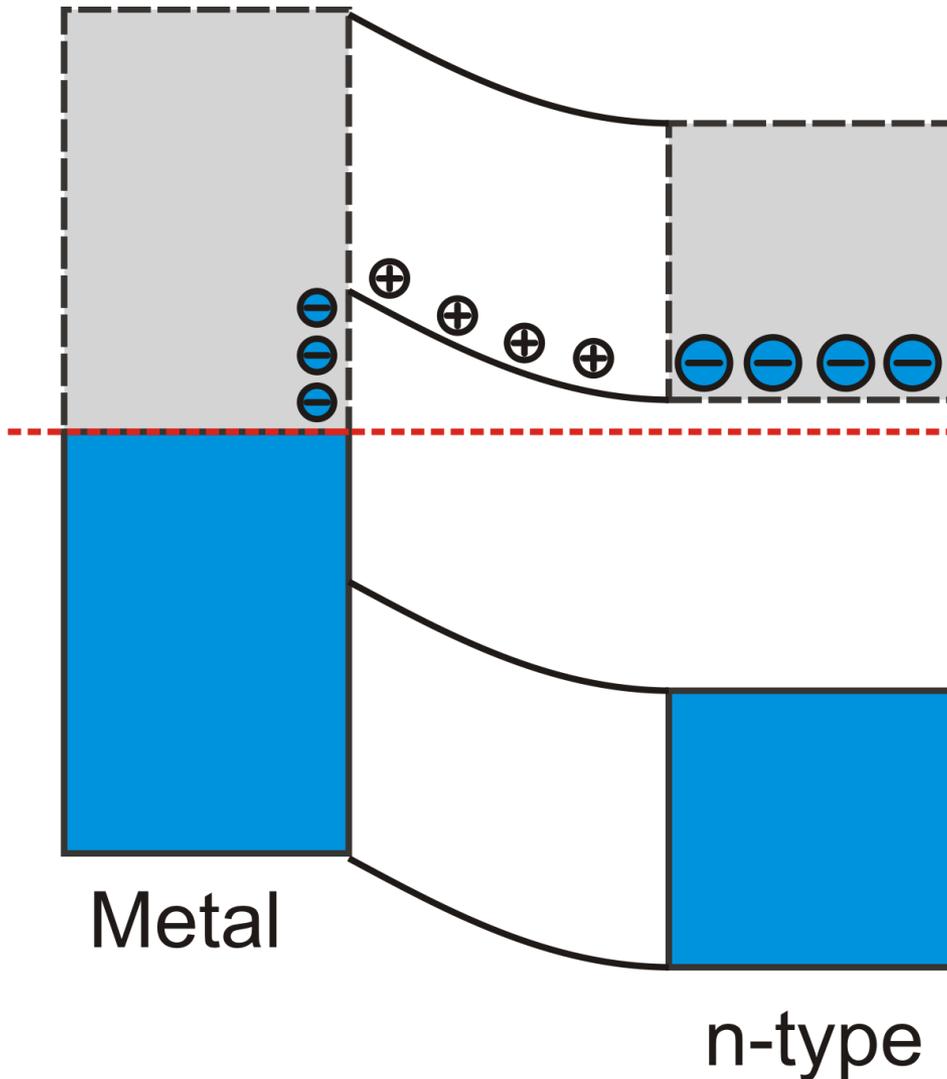
## 2.3 Band Bending at p-n Junction



## 2.4 Band Bending at Metal-SC Junction



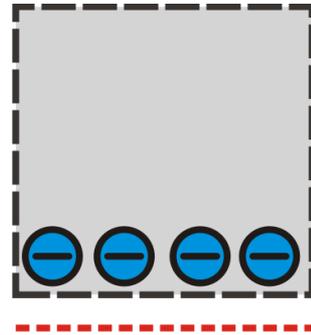
## 2.4 Band Bending at Metal-SC Junction



### Note:

- Metals have carrier densities of  $\sim 10^{22} \text{ cm}^{-3}$
- Screening lengths of ca. 0.1 nm
- Semi-cond. have carrier densities  $\sim 10^{17} \text{ cm}^{-3}$
- Screening lengths of ca. 10 nm

# 2.5 Band Bending at Surfaces

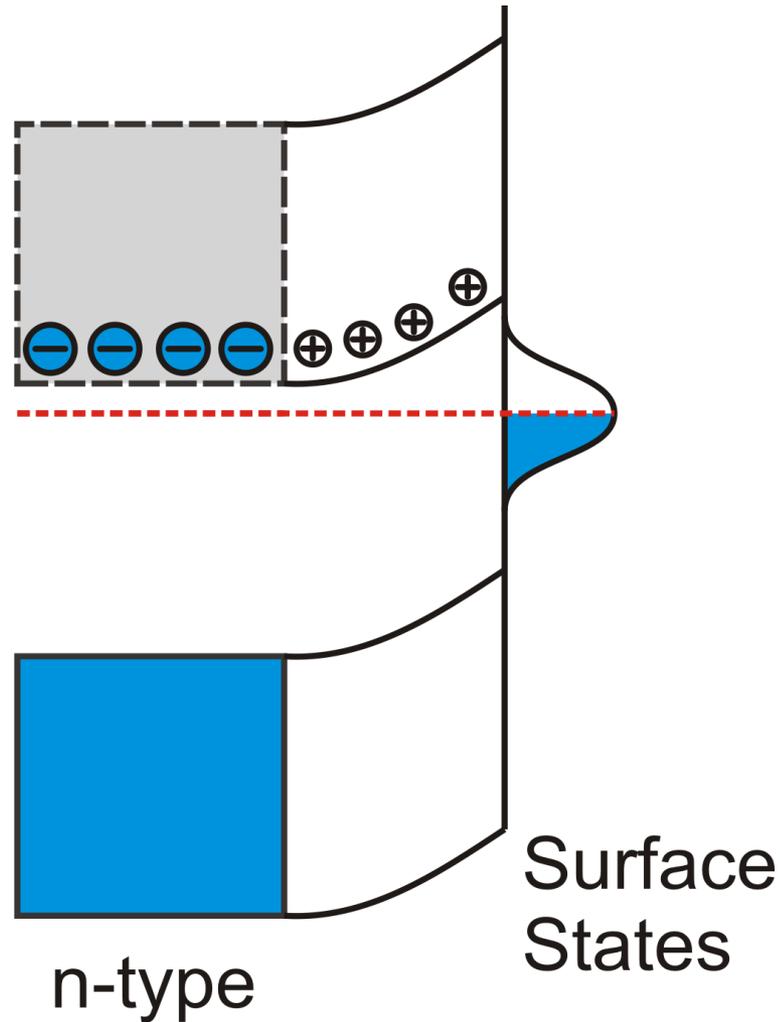


n-type



Surface  
States

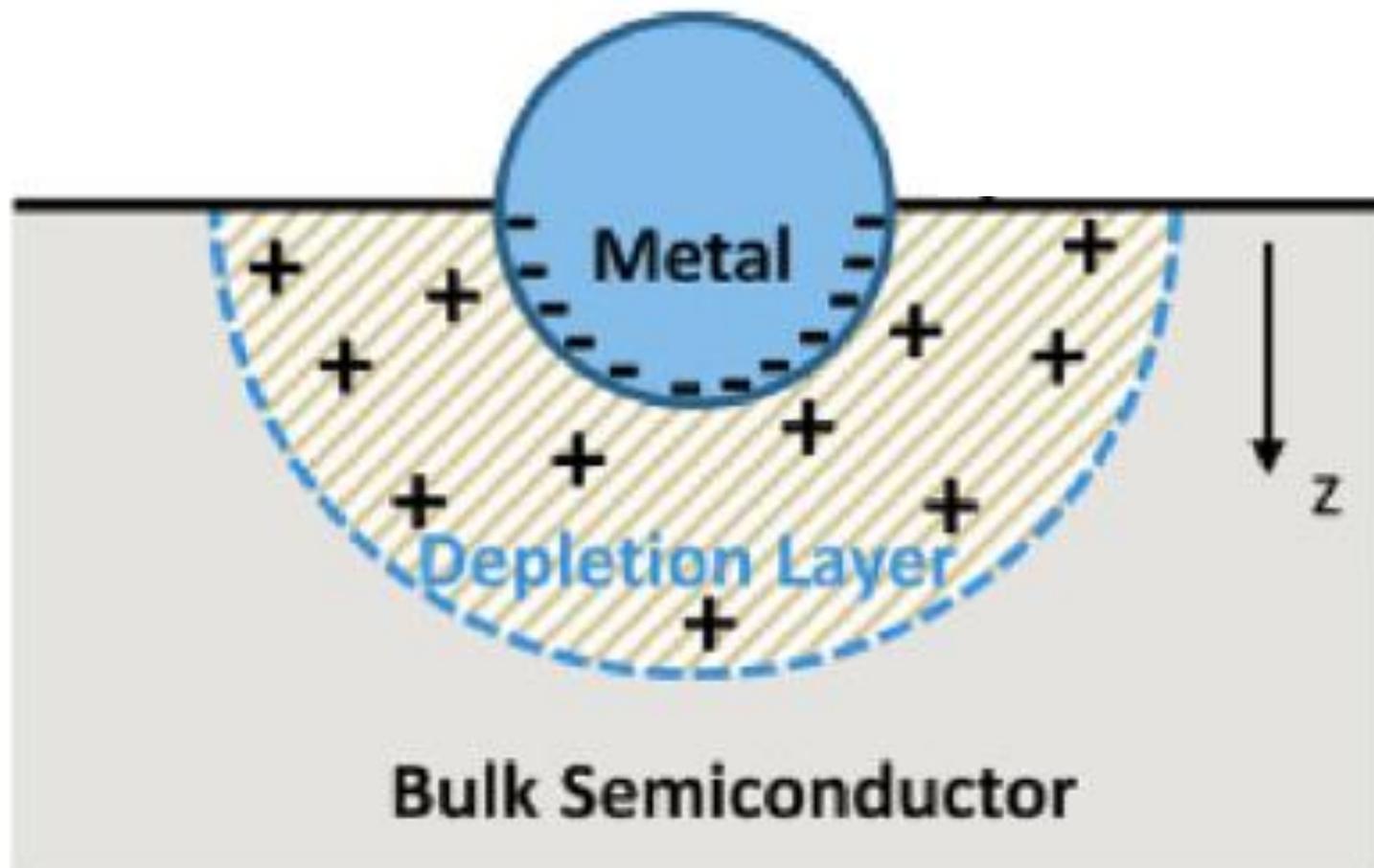
# 2.5 Band Bending at Surfaces

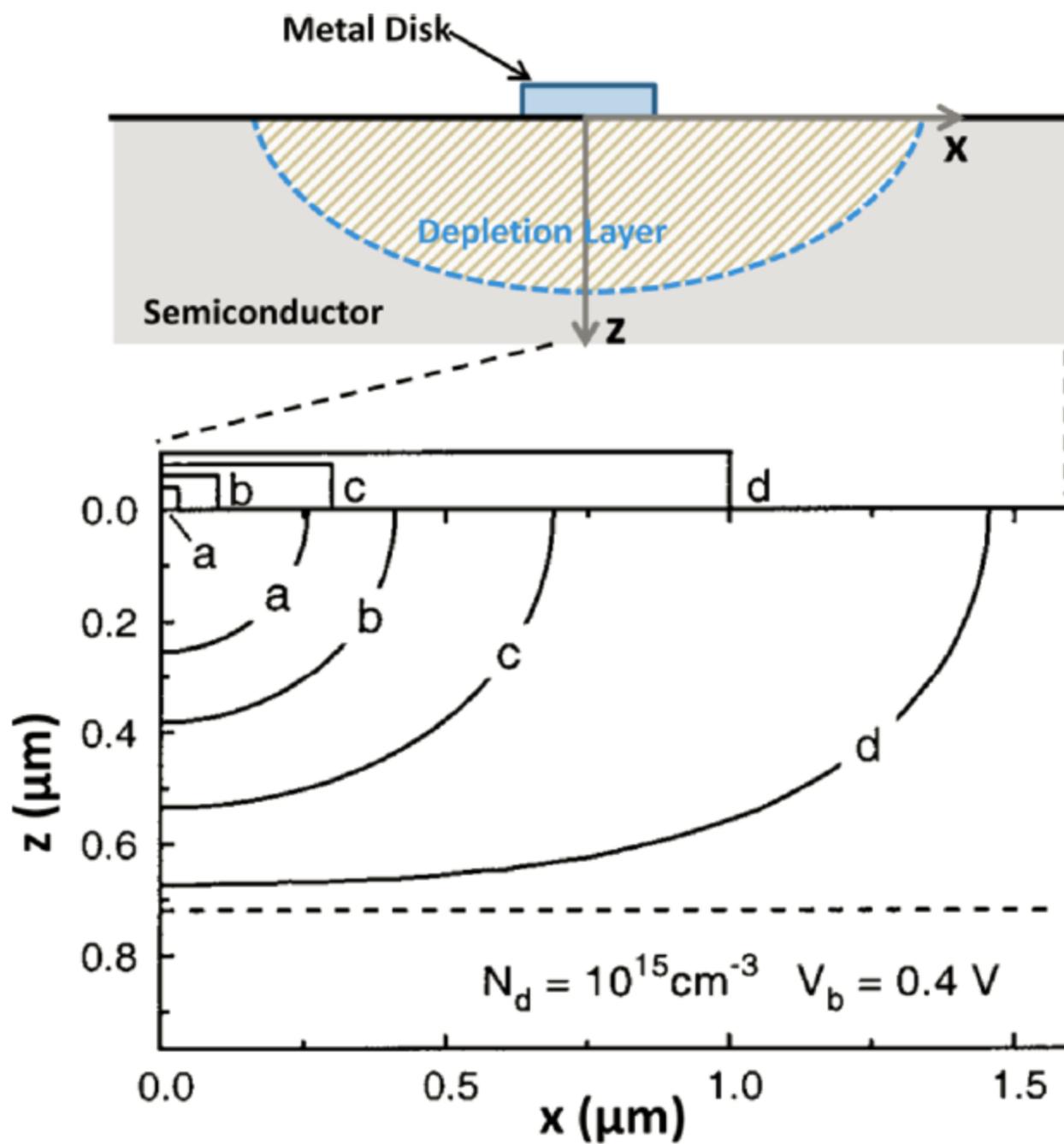


# What does all this have to do with catalysis?

- Surfaces can have significantly different energy levels than expected from bulk considerations.
- This will affect interactions and charge transfer with adsorbed molecules.
- Adsorbed molecules can also change the surface band bending.

# Influence on Metal-Support Interaction

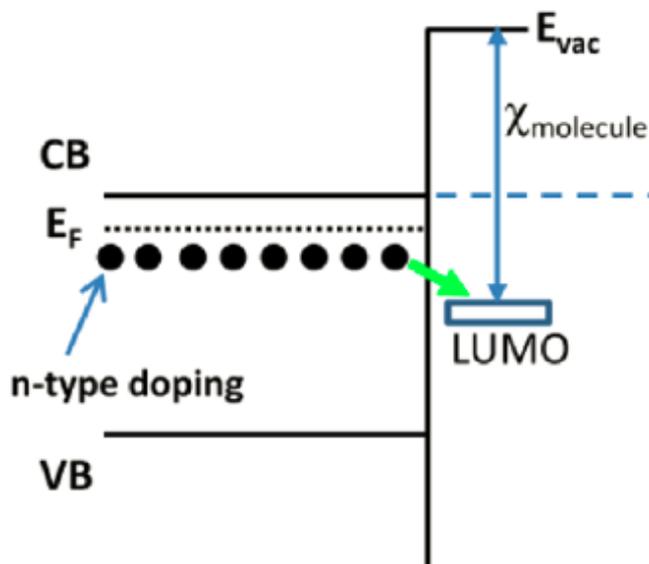




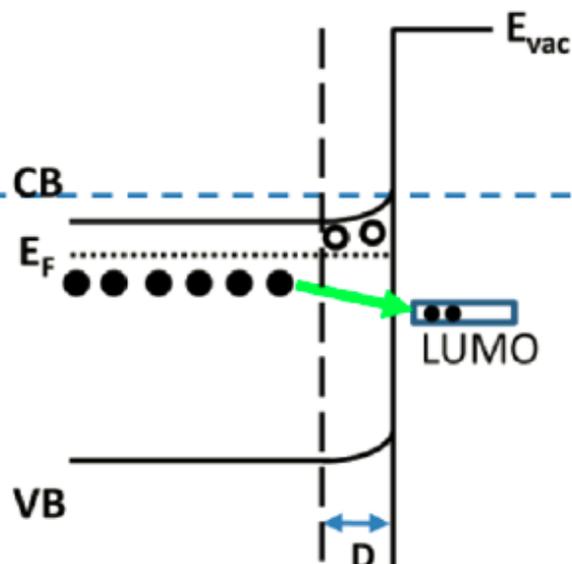
# Band Bending and Adsorption

## n-type Semiconductor

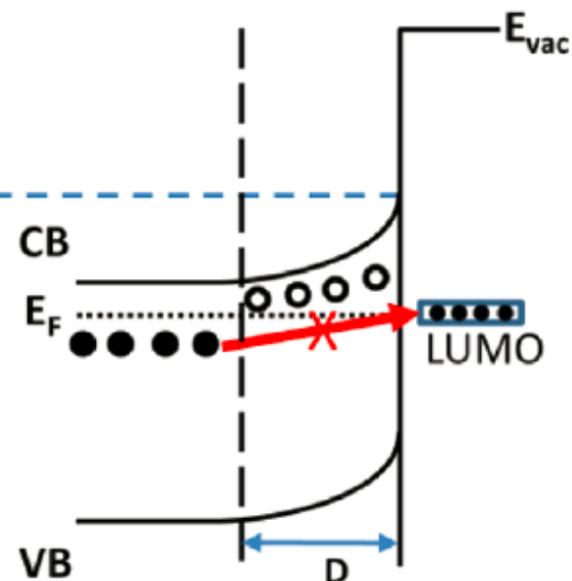
(a)



(b)

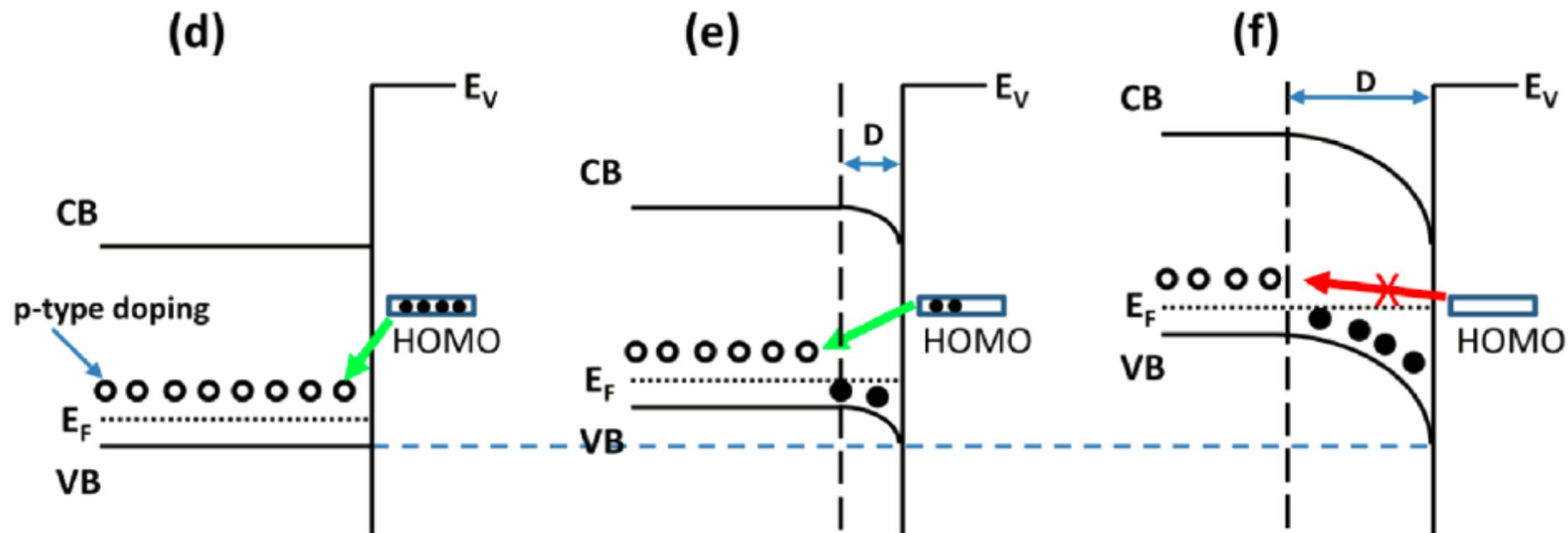


(c)

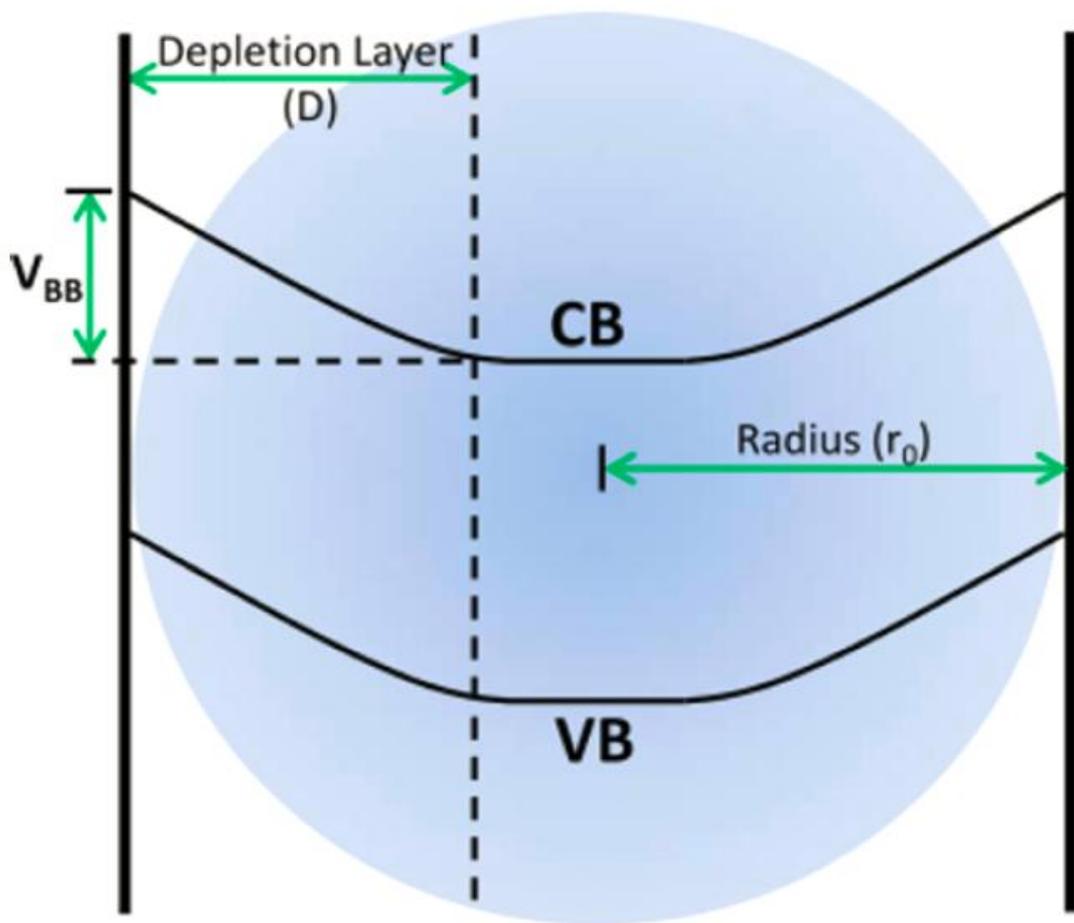


# Band Bending and Adsorption

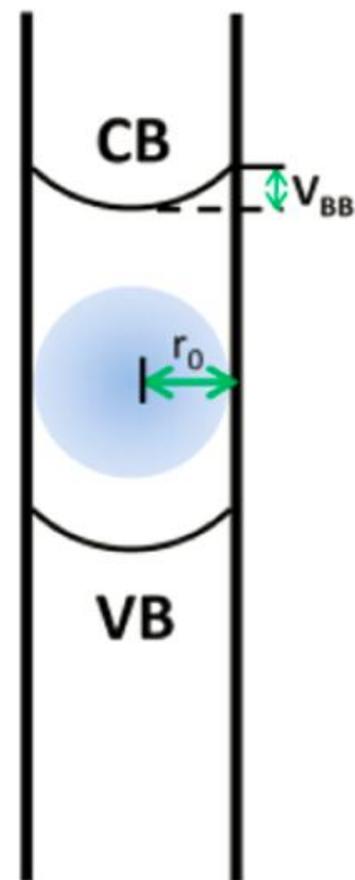
## p-type Semiconductor



# Band Bending and Particle Size

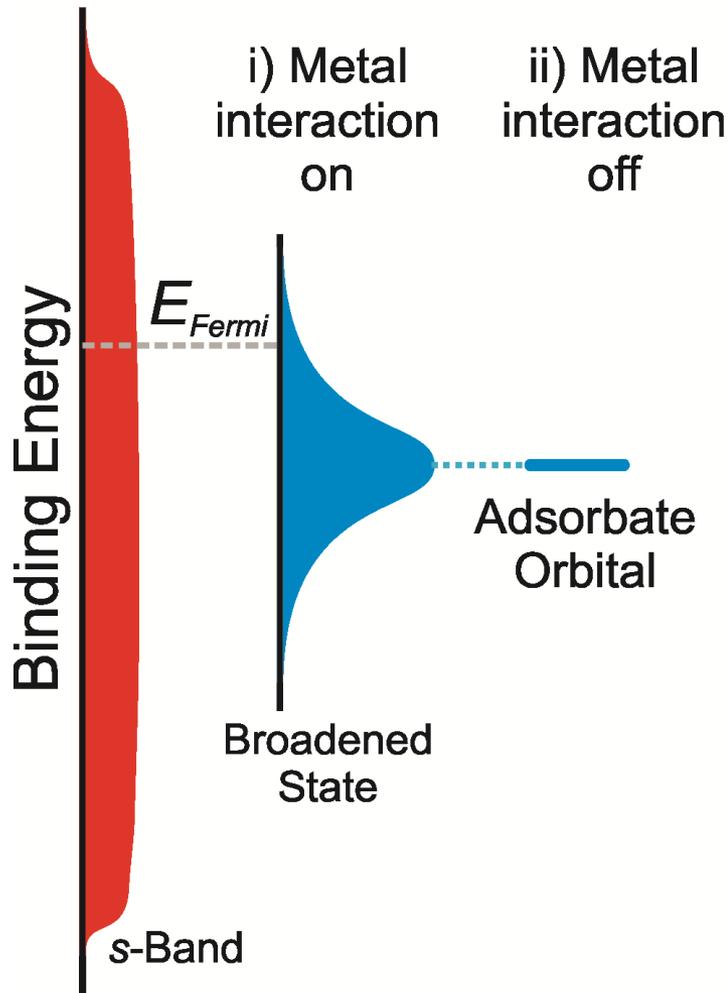


Big Particle

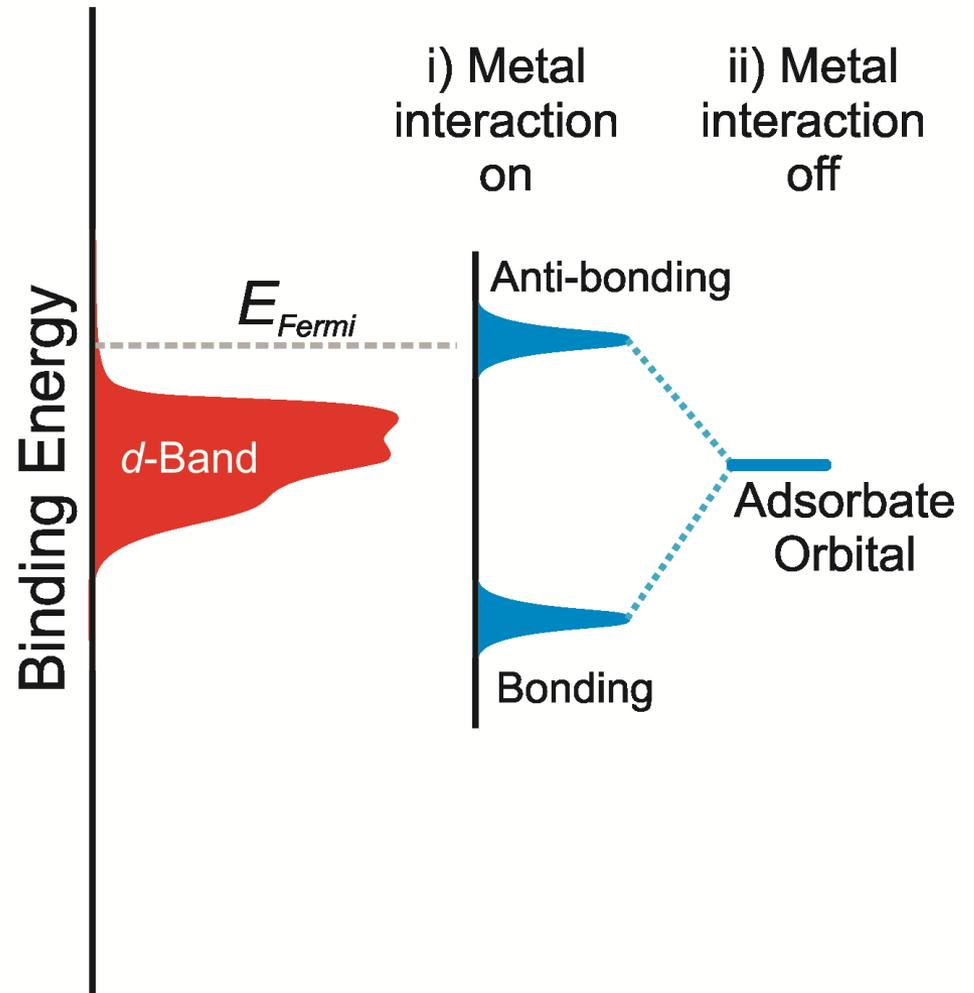


Small Particle

### a) Case 1: Broad s-Band



### b) Case 2: Narrow d-Band



# Further Reading

- “Solids and Surfaces: A Chemist’s View of Bonding in Extended Structures” R. Hoffmann
- “Electronic Properties of Materials” D. Jiles
- “Solid Surfaces, Interfaces and Thin Films” H. Lüth
- “Band Bending in Semiconductors: Chemical and Physical Consequences at Surfaces and Interfaces” Zhang, Yates: *Chem. Reviews* (2012)

Any Questions?

# Doping with defects

