



Analysis of Local Structure by Atomic Pair Distribution Function

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Lecture series: Modern Methods in Heterogeneous Catalysis Research







- What for do we need PDF?
- Scattering intensity
- Mathematical basis of the PDF
- Direct structural information from PDF
- PDF experiments: X-rays, neutrons, electrons
- Catalysis related examples



What is a PDF?





Thomas Proffen: "Total Scattering. The Key to Understanding disordered, nano-crystalline and amorphous materials" Tutorial 9th Canadian Powder Diffraction Workshop





The challenge : Knowing the local structure

- Traditional crystallographic approach to structure determination is insufficient or fails for
 - * Non crystalline materials
 - Disordered materials: The interesting properties are often governed by the defects or local structure !
 - Nanostructures: Well defined local structure, but long-range order limited to few nanometers (-> poorly defined Bragg peaks)
- A new approach to determine local and nano-scale structures is needed.



S.J.L. Billinge and I. Levin, **The Problem with Determining Atomic Structure at the Nanoscale**, *Science* **316**, 561 (2007).



What for do we need PDF



Conventional XRD analysis of Bragg intensities yields the **average** structure of materials which can be deceiving !

Considering going to a party where all you know is the average age is 40..



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What exactly is meant by the difference between the local structure and the average crystallographic structure?









Neutron powder diffraction data

fdoo'

T. Egami & S. Billinge "Underneath the Bragg Peaks"







T. Egami & S. Billinge "Underneath the Bragg Peaks"



Pair Distribution Function from total scattering experiments



PDF from total scattering experiments (X-ray, neutron or electron sources)



Emil S. Bozin, (Brookhaven National Laboratory) School and Conference on Analysis of Diffraction data in Real Space 2013



Pair Distribution Function from total scattering experiments



Reduced structure function

2D diffraction pattern









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the source and detector are far away from the object the object is small



Elastic - no energy loss coherent - no random phase shift scattering







 $Q = 2|k|\sin\theta$ scattering vector

 $\Delta \phi = \mathbf{Q} \bullet \mathbf{R}$ phase shift

Scattering intensity

$$I(\mathbf{Q}) = \Sigma A_m e^{-i (\mathbf{Q} \cdot \mathbf{R}_m)} \cdot \Sigma A_n e^{i (\mathbf{Q} \cdot \mathbf{R}_n)} = \Sigma \Sigma A_m A_n e^{-i (\mathbf{Q} \cdot (\mathbf{R}_m - \mathbf{R}_n))}$$

 $\{\mathbf{R}_{m}^{-}\mathbf{R}_{n}\}$ – Object property (unique set)





Atomic scattering amplitude E(Q)

E(Q) is the sum of all the electrons scattering amplitudes

 $A_{at}(Q) = \sum_{i=0}^{i=0} 1z$ -1 ##Ae(Q) e1-i(Qr)

Atomic scattering factor f(Q)



 $f(Q) = A_{at} (Q) / A_e (Q)$





Bohr radius $a_0 \approx 0.5292$ Å

 $r_j=0$, f(Q)=Z and doesn't depend on Q

Since electrons are not concentrated in one point f(Q) depends on Q = 4 sin θ/λ



Atomic scattering factor











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Mathematical basis of the PDF





Q (Å⁻¹)



Mathematical basis of the PDF



$$F(Q) = Q[S(Q) - 1] \xrightarrow{\mathsf{FT}} G(r) = \frac{2}{\pi} \int_{0}^{\infty} F(Q) \sin(Qr) dQ$$

In practice, Q_{min} and Q_{max} are limited

 Q_{max} = 30-50 Å⁻¹ (X-ray Synchrotron) Q_{max} = 10-23 Å⁻¹ (electrons)

$$G(r) = \left(\frac{2}{\pi}\right) \int_{Q_{\min}}^{Q_{\max}} F(Q) \sin(Qr) dQ$$



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Truncation of Fourier series







Mathematical basis of the PDF









 $N_c = \int r 1 f r^2 m R(r) dr$ the number of neighbors



Mathematical basis of the PDF











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Direct information from the PDF



PDF contains the following structural/ configurational information:

- Average interatomic distances
- A structural disorder

COD)

- Average coordination properties —
- Particle size effect

- PDF peak position
- PDF peak width (FWHM)
- Integral intensity of PDF peaks
- PDF peak cut-off





Direct information from the PDF



• PDF peak cut-off



K.L. Page, Th. Proffen, H. Terrones, M. Terrones, L. Lee, Y. Yang, S. Stemmer, R. Seshadri and A.K. Cheetham, Direct Observation of the Structure of Gold Nanoparticles by Total Scattering Powder Neutron Diffraction, *Chem. Phys. Lett.* **393**, 385-388 (2004).





• The PDF peak position yields bond length directly

 $In_{1-x}Ga_xAs$

The lattice constant changes linearly with x, following Vegard's low

The crystallographic (Ga,In)-As distance represents only the average distance between the atoms at the (Ga,In) and As sites and corresponds to neither the actual Ga-As nor In-As distances.



Petkov, V., Jeong, I.-K., Chung, J.S., Thorpe, M.F., Kycia, S. & Billinge, S.J.L. (1999) Phys. Rev.Lett., 83, 4089.

Direct information from the PDF



• The integral intensity of PDF peaks yields the coordination number of the atom-atom correlation



Petkov et al., PHILOSOPHICAL MAGAZINE B, 1999, VOL. 79, No. 10, 1519-1530

Direct information from the PDF









The PDF peak width reveals information about static and dynamic disorder of atomic pairs







PDF peak height and PDF peak width are inversely proportional

IrO₂ crystalline phase (Ir⁴⁺) IrO_x hydroxide disordered phase (Ir⁴+Ir³⁺)



Direct information from the PDF

Mn



The PDF peak width as a function of doping and temperature

La_{0.79}Ca_{0.21}MnO₃

shows a phase transition from insulating paramagnetic phase to a metallic ferromagnetic phase

Insulating phase has localized charges: Mn^{3+} and Mn^{4+} Metallic phase has delocalized charges: $Mn^{(3+x)+}$



Th. Proffen *et al.* Z. Kristallogr. 218 (2003) 132 – 143



Direct information from the PDF



The PDF peak width as a function of temperature



J. Phys.: Condens. Matter 12 (2000) L723–L730







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- In house X-ray experiments:
 - Low r resolution, slow (and/or poor statistics)
 - $Mo: Q_{max} \sim 17 \text{Å}^{-1}, Ag: Q_{max} \sim 20 \text{Å}^{-1}$
 - Easy
 - Flat plate transmission or reflection geometry



- Synchrotron experiments:
 - Little precious beamtime, but measurements are quick(er)
 - High *r* resolution using high energy X-rays ($Q_{max} > 60 \text{Å}^{-1}$)
 - High intensity (x 10,000 times stronger)
 - Parallel beam optics
 - Generally flat plate or cylindrical transmission geometry

Where can we measure X-ray PDF?



European Synchrotron Radiation Facility (ESRF) Location: Grenoble (France)



Advantages:

High brightness → small samples High collimation → high resolution of 20 Continuous spread of wavelengths

X-ray beam Injection point Storage ring Electron beam Bending magnet

Disadvantage:

the synchrotron sources are not available on a daily basis

Where can we measure X-ray PDF?

micro

fdoo'







X-rays or neutrons?



Highly complementary, depends on the nature of the problem at hand

Synchrotron X-rays

Neutrons

b – fluctuates with Z

sensitive to light elements

b – constant with Q

sample amount $\sim 5-10~g$

acquisition time: ~2-5 hours, possibly longer (much faster at the latest neutron sources, e.g., Oak Ridge Laboratory, Tennessee)

Choice could depend also on issues related to the underlying structure, or sample characteristics such as size or absorption



H C O TI Fe NI Neutrons 1 0 0 0 0 0 0 1 2 46 54 58 47 0 58 48 0 60 49 a50 0 62

X-ravs

f – decreases with $Q = 4\pi \sin \theta / \lambda$

sample amount $\sim 100 \text{ mg}$

acquisition time:

~5 hours for point detector ~10 seconds for area detector ~0.1 seconds a-Si area detector

f – proportional to Z (atomic number)

poor sensitivity to light elements





X-rays or electrons?



Pros

- 1. A low amount of a sample is needed (fractions of μ g).
- 2. Synchrotron beam time is not needed!!! TEM operator can easily switch from the image- to diffraction-mode.
- 3. It is a local method in the TEM.



that of X-rays









X-rays or electrons?



10 nm Au nanoparticles ($Q_{max} = 10 \text{\AA}^{-1}$)



2D ED pattern







Software for PDF modelling



The Software to abstract PDF from diffraction data

PDFgetX3 for XRD pattern PDFgetN for NPD pattern

SUePDF for electron diffraction pattern



The software to model PDF



DiffPy-CMI advanced modelling

THE BILLINGE GROUP: COLUMBIA UNIVERSITY Department of Applied Physics and Applied Mathematics https://thebillingegroup.com/simon-billinge/





Advanced modeling for nanoscale materials with SrFit and SrReal using Debye equation



Christopher L. Farrow et al. J. Am. Chem. Soc. 2013, 135, 6403-6406







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Outline

• Catalysis related examples





PDF of an Operating Fuel Cell In Operando

Pt/C catalyst for Data were first measured with this device at 11IDB at APS



Redmond, E.L., Setzler, B.P., Juhas, P., Billinge, S.J.L. & Fuller, T.F. (2012) Electrochem. Solid State, 15, B72.



Catalysis related examples





V. Petkov et al. ACS Nano 3 441-445 (2009)



Catalysis related examples



IrOx (FHI-made) shows much higher OER catalytic performance than IrOx (Alfa Aesar)



Except for the presence of metal impurities, XRD data does not show essential difference between commercial and FHI-made IrO_x phases



HAADF



IrO_x -commercial







Atomic resolution HAADF



IrO_x -commercial

IrO_x-FHI



EDX analysis shows O/Ir \sim 2 for both samples IrO_x-FHI shows in addition a trace amount of K

Crystallographic approach





Ir⁴⁺

micro

fdoo

 $Ir^{4+}Ir^{3+}$

 $Ir^{4+}Ir^{3+}$



ePDF analysis







ePDF modeling



IrO_x-commercial











Ir $O_{2,3}$ -edge (5p \rightarrow 5d, 6s electron transitions)

micr

02



EELS analysis







E-beam effect







E-beam effect







Summary



The PDF allows to reveal the details of the local structure which is not available in the case of conventional structural analysis which deals with the average structure









Figure 1 Background modeling for electron powder diffraction data (solid-black) of nanoporous carbon: power-law model (dash-red) compared with Eq. (4) Laurent-type model (dot-blue) with N = 7.















Finite crystal



 $\sim \sin t^2 (N_1 \varphi_1) / \sin t^2 \varphi_1 \cdot \sin t^2 (N_2 \varphi_2) / \sin t^2 \varphi_2$ ·sin

$$\varphi_1 = \pi (\mathbf{S} \cdot \mathbf{a}), \varphi_2 = \pi (\mathbf{S} \cdot \mathbf{b}), \mathbf{W} = \mathbf{W} (\mathbf{S} \cdot \mathbf{c})$$



 $N_1 = 10$

There are high peaks in I(S) when $(S \cdot a^*)$ - integer

 $N \longrightarrow \infty$ $I \sim N_1^2 \cdot \delta(\mathbf{S} \cdot \mathbf{n} \cdot \mathbf{a}^*)$



micro

(d00)

OT





Anode	Ка	Κβ
Cu	1.54184 Å	1.39222 Å
Мо	0.71073 Å	0.63229 Å





Scattering intensity distribution



Thomson scattering formula

 $I = I \downarrow 0 \ e^{\uparrow} 4 \ /r^{\uparrow} 2 \ m^{\uparrow} 2 \ c^{\uparrow} 4 \ (1 + \cos^{\uparrow} 2 \ 2 \ M)/2)$





X-ray scattering by crystalline (periodic) ensemble of atoms

 $R_{mnp} = ma + nb + pc$

Periodic arrangement of atoms is given by lattice with the basis vectors **a**,**b**,**c**

Scattering amplitude $A(S) = A_e \Sigma f_j e^{-i2\pi(S \cdot R_j)} = -A_e f \Sigma e^{i2\pi(S \cdot R_mnp)}$

In general A(S) ~ 0 for infinite crystal, i.e. m, n and p $\rightarrow \infty$

But if $(\mathbf{S} \cdot \mathbf{R}_{mnp}) = (\mathbf{S} \cdot \mathbf{a})m + (\mathbf{S} \cdot \mathbf{b})n + (\mathbf{S} \cdot \mathbf{c})p = q - integer$

then $e^{i2\pi(S \cdot R_{mnp})} = 1$ and $|A(S)| = N A_e f(s)$, N – number of atoms,

There is the only solution: $S = \{G_{hkl}\}=\{ha^* + kb^* + lc^*\}$ with $a^* = [bxc]/V$, $b^* = [cxa]/V$, $c^* = [axb]/V$, V = (a[bxc])



Ewald construction



Space of wave vectors - Ewald construction





Atomic scattering factor





O, z=8, atomic radius =0.66 Cl, z=17, atomic radius=1 Cl-, z=18, radius= 1.81 K, z = 19, atomic radius=2.02 K+ =18, radius= 1.33

X-ray atomic factors of O, Cl , Cl⁻ and K⁺ ; smaller charge distributions have a wider atomic factor.



Structural flexibility





HAADF of **I** −MnO₂



Grangeon et al. Geochem Trans (2015) 16:12



Daria Mikhailova et al., Inorg. Chem. 2016, 55,







