

Modern Methods in Heterogeneous Catalysis Research



# **Adsorption and desorption**

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Literature:

R.I. <u>Masel</u>, Principles of adsorption and reaction on solid surfaces, Wiley, New York(1996).

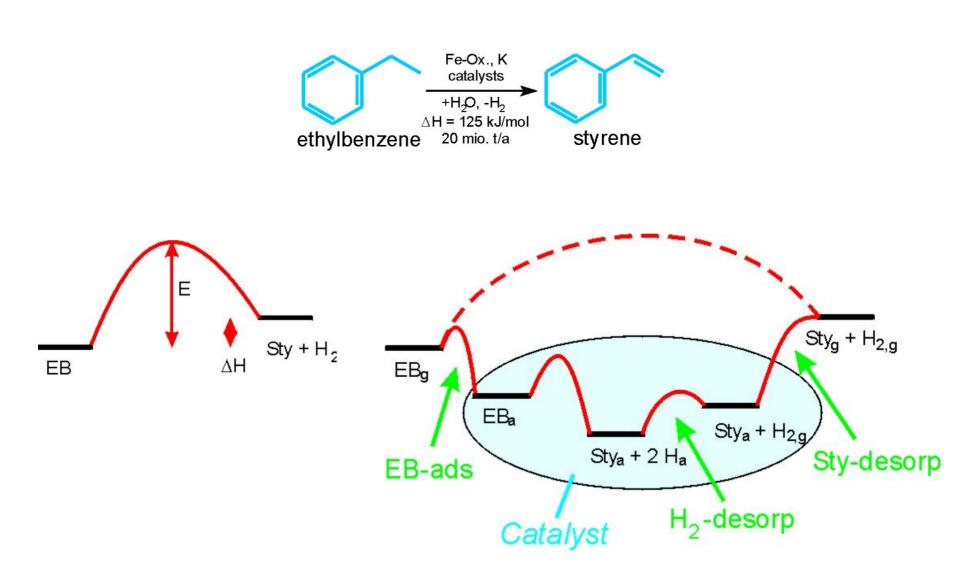
K. Christmann, Surface physical chemistry, Steinkopff, Darmstadt (1991).

G.A. Somorjai, Introduction to surface chemistry and catalysis, Wiley, New York (1994).

M. Henzler, W. Göpel, Oberflächenphysik des Festkörpers, Teubner, Stuttgart (1991).

W. <u>Ranke</u> and Y. <u>Joseph</u>, Determination of adsorption energies and kinetic parameters by isosteric methods, Phys. Chem. Chem. Phys. 4 (2002) 2483.

### Adsorption and catalysis



### Contents

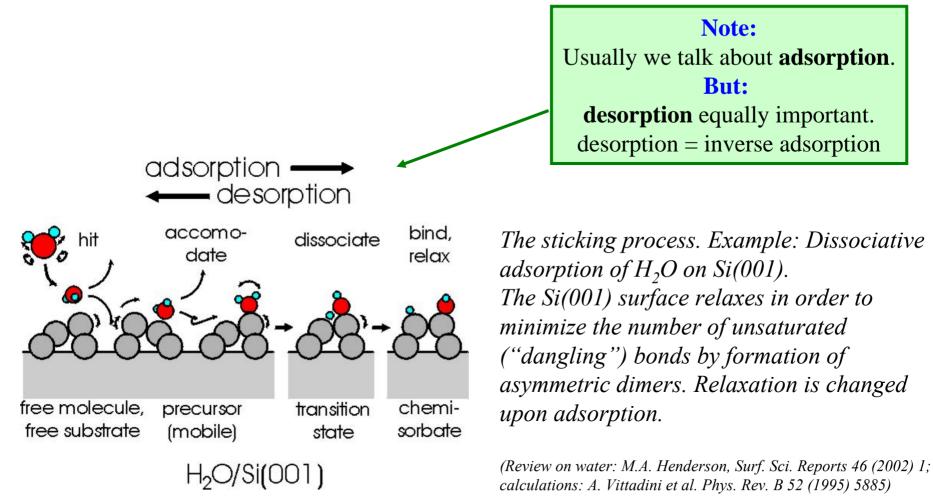
### Part I

- 1. Some definitions and concepts
- 2. Physisorption and chemisorption
- 3. Energetics: Potential curves

### Part II

- 4. Adsorption-desorption-equilibrium: Isotherms and isobars
- 5. A-D energetics: Isosteric heat of adsorption
- 6. A-D kinetics: The shape of isotherms/isobars
- 7. The magnitude of v: Entropy of the activated state
- 8. Further isotherms for mono- and multilayer adsorption

#### 1. Some definitions and concepts



#### **Adsorption:**

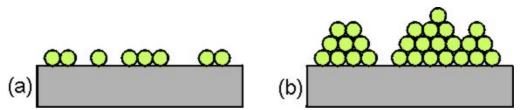
molecules from gas phase or solution bind in a layer of condensed phase on a solid or liquid surface. The molecules are called **adsorbate**, the substrate is called **adsorbent**. The process of binding is called **adsorption**. Removal of the molecules is called **desorption**. **Accomodation** and **Sticking**:

The incoming particle has to get rid of its **translational** and **rotational energy** and has to assume a suitable **spatial configuration**. The surface may have to rearrange: **Accomodation**. If this does not happen: **Reflection**.

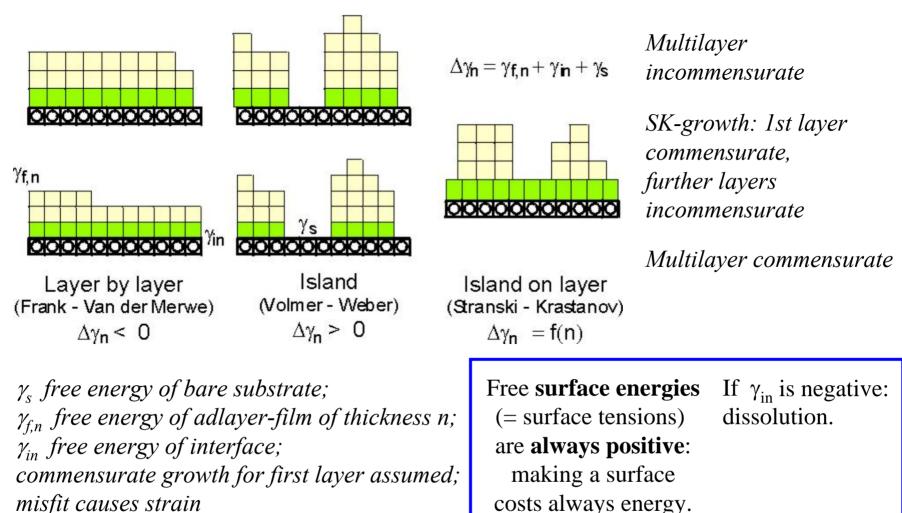
**Sticking**: The **sticking coefficient** *s* or **sticking probability** is the ratio of particles being bound compared to those hitting the surface.

The value of *s* often reaches 1 but may also be almost infinitely small.

Rate of molecules striking 1 cm <sup>2</sup> of surface:	N <sub>2</sub> at 273 K:	$j_{N}$ (cm <sup>-2</sup> s <sup>-1</sup> )	Time needed for 1 ML ( <i>s</i> =1)
	1000 mbar	3×10 <sup>23</sup>	5×10 <sup>-8</sup> s
$j_N = 2.63 \times 10^{22} \text{ p/(MT)}^{1/2} \text{ cm}^{-2} \text{ s}^{-1},$	1 mbar	3×10 <sup>20</sup>	5×10 <sup>-6</sup> s
p in mbar, M in g/mol.	10 <sup>-3</sup> mbar	3×10 <sup>17</sup>	5×10 <sup>-3</sup> s
	10 <sup>-6</sup> mbar	$3 \times 10^{14}$	5 s
For comparison:	10 <sup>-10</sup> mbar	3×10 <sup>10</sup>	$5 \times 10^3 \text{ s}$
Density of atoms on solid surfaces			
Pt(111): $1.5 \times 10^{15} \text{ cm}^{-2}$			
Si (001): 6.8×10 <sup>14</sup> cm <sup>-2</sup>			



Monolayer and multilayer adsorption



#### Not treated here:

Formation of new bulk components or bulk phases as result of

reactions like oxidation alloying dissolution etc.

Adsorption from liquid phase

- 2. **Physisorption and chemisorption**
- 2.1 **Physisorption**

Nonpolar gases on nonpolar surface, no chemical interaction: Van-der-Waals (induced-dipole – induced-dipole) interaction. Starting from the Lennard-Jones 6-12 potential (interaction between single atoms or molecules).

 $V(R) = 4\varepsilon \left[ (\sigma/R)^{12} - (\sigma/R)^6 \right],$  $\varepsilon$  depth of potential minimum;  $\sigma$ =0.891  $R_{e}$  (equilibrium distance). the London-equation,

$$E_{Ms}(z_M) = \rho_s \pi \ C_{Lond} \left( \frac{0.2(r_e^s)^6}{z_M^9} - \frac{1}{z_M^3} \right)$$

 $\rho_{\rm s}$ : density of surface atoms;  $\tilde{C}_{Lond}$ : London constant, depends on polarizabilities and ionization potentials;  $r_{a}^{s}$ : equ. distance when the adatom interacts with one atom only;

> **Compare**: Thermal E

> > (300 K)

 $z_M$ : distance of adatom from surface.

was derived for non-polar gases on non-polar non-conducting surfaces. on conducting surface: Same functional dependence.

> **Characteristic:**  $z_{\rm M}$  >2Å; RT~2.5 kJ/mol  $\Delta H_{ad}$ : 1.4 kJ/mol (He/graphite) – 42 kJ/mol (C<sub>5</sub>H<sub>12</sub>/graphite)

Even if  $\Delta H_{ad}$  is small,  $\Theta$  may be high at high pressure, see isotherms/isobars

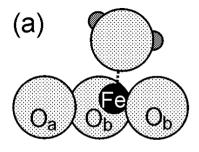
Van-der-Waals interaction is responsible for condensation of nonpolar gases like inert gases, alkanes, aromatics...

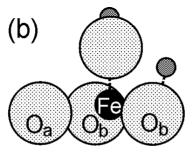
#### 2.2 Chemisorption

Electronic structure, bond lengths and bond angles strongly affected.
Substrate and adsorbate **not** separable. Simple models do not exist.
Theoretically, the whole system (substrate + adsorbate) has to be treated quantum mechanically which has become feasible during the last years.

Characteristic:  $z_{M} = 1 - 1.5 \text{ Å};$  $\Delta H_{ad}$  up to several hundred kJ/mol.

Covalent -	(Example: $H_2O/Si(001)$ ., see above.	
Ionic -	(Example: Alkali adsorption)	
Acid-base -	(Example: $H_2O/Fe_3O_4(111)$ )	
H-bond -	(Example: $H_2O$ -dimers and islands, lateral interaction)	
H-DONG -	(Example: $H_2O$ -dimers and islands, lateral interaction)	





Model for the interaction and dissociation of water on an acidic Fe-site on the surface of  $Fe_3O_4(111)$ ; the Fe is formally  $Fe^{3+}$  (Y. Joseph et al. Chem. Phys. Lett. 314 (1999) 195).

#### Hardness (Pearson) $\eta$ :

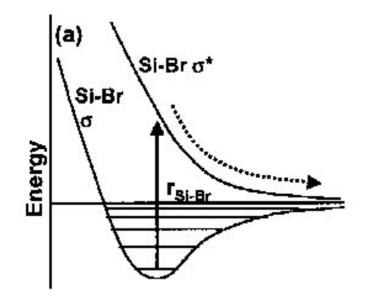
Measure of the capability to share electrons. Hard ( $\eta$  large): Species that form strong ionic bonds; Medium: Species that form mixed covalent and ionic bonds; Soft ( $\eta$  small): Species that form strong covalent bonds;

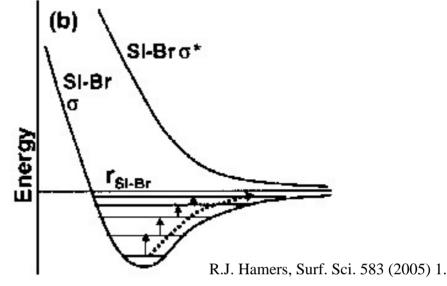
Charge acceptor: acid; charge donor: base. Strong interaction between a **hard** acid and a **hard** base or a **soft** acid and a **soft** base.

Weak interaction between a hard acid and a soft base.

Hard acids	H <sup>+</sup> , Li <sup>+</sup> , Mg <sup>2+</sup> , Cr <sup>3+</sup> , Co <sup>3+</sup> , Fe <sup>3+</sup> )Al <sup>3+</sup> , Al(CH <sub>3</sub> ) <sub>3</sub> , bulk Si
Borderline acids	Mn <sup>2+</sup> , Fe <sup>2+</sup> , Co <sup>2+</sup> , Ni <sup>2+</sup> , Cu <sup>2+</sup> , Ru <sup>2+</sup> , Os <sup>2+</sup> , Ir <sup>3+</sup> , Rh <sup>3+</sup>
Soft acids	Cu <sup>+</sup> , Ag <sup>+</sup> , Pd <sup>2+</sup> , Pt <sup>2+</sup> , Ga(CH <sub>3</sub> ) <sub>3</sub> , O, Cl, N, RO, ROO, metallic atoms,
	metal clusters
Soft bases	I <sup>-</sup> , CO, $C_2H_4$ , $P(C_6H_5)_3$ , $C_6H_6$ , H <sup>-</sup> , $H_2S$ , metal surfaces
Borderline bases	$C_5H_5N, NO_2^-, SO_3^{2-}, C_6H_5NH_2$
Hard bases	$F^{-}$ , $Cl^{-}$ , $H_2O$ , $NH_3$ , $OH^{-}$ , $CH_3COO^{-}$ , $RO^{-}$ , $ROH$ , $O^{-}$

**Energetics: Potential curves** 





recently observed "Ten or twenty phonons"transfer their energy collectively to an electron which is excited into an antibinding state from which dissociation occurs.

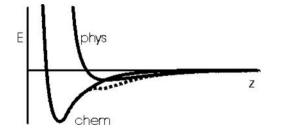
A number of phonons is sequentially absorbed. The adsorbate-substrate vibration gains amplitude until dissociation occurs.



Classical analog:

Interference-generated freak wave (Monsterwelle, Kaventsmann) Resonance-like amplitude build-up (Resonanzkatastrophe)

Adsorption, accomodation: invert all arrows



B+B  $E_{D}^{a_{1}}$   $B_{2}$   $B_{2}$ Molecular State
Distance From The Surface ---Dissociated State

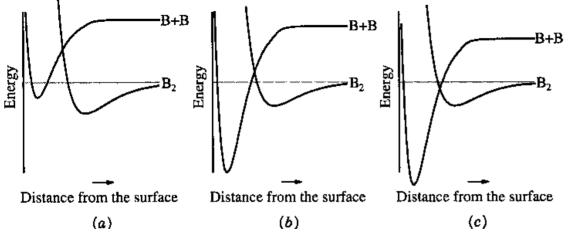
Potential curves for approach of a non-dissociated particle to a surface:

physisorption, chemisorption and sum curve (dotted).

 $\Delta H_{phys} \ll \Delta H_{chem}$ typically  $\Delta H_{phys} = 20 - 80$  kJ/mol,  $\Delta H_{chem} > 100$  kJ/mol.

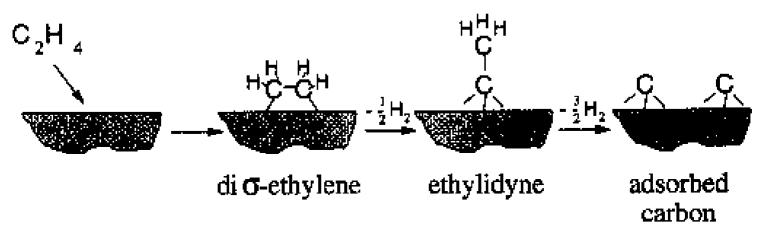
**But:** Magnitude of  $\Delta H_{ad}$  alone is **not** a sufficient criterion for the distinction: Change of geom. and electron. structure possible.

Potential curve for the dissociative adsorption of a  $B_2$  molecule on a surface (Masel fig. 3.8, p.119.)



Potential curves for (a) pure molecular adsorption, (b) activated dissociative adsorption (c) nonactivated dissociative adsorption (Masel fig. 3.9, p.119)

#### **3.2** Dissociative adsorption, several sequential steps



*The mechanism of ethylene decomposition on Pt(111).* (Proposed by Kesmodel et al. [1979] and confirmed by Ibach and Lehwald [1979], Masel fig. 3.10, p.121).

Example for **"irreversible"** adsorption.

Irreversible is not strictly the correct term:

It is irreversible under usual experimental conditions where a part of the formed species (here hydrogen in the form of gaseous  $H_2$ ) is irreversibly removed from the system.

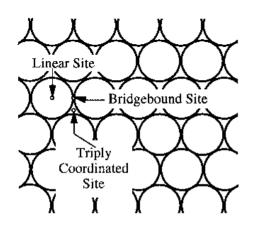
#### **3.3** Adsorption site

Surfaces are inhomogeneous

interaction with certain sites is stronger than with others.

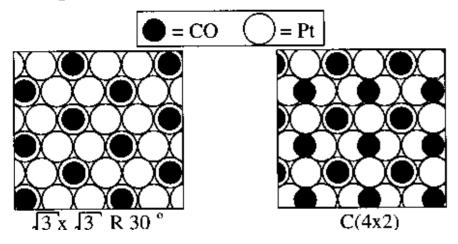
Intrinsic inhomogeneity from periodic atomic structure

✓ ordered adsorbate structures



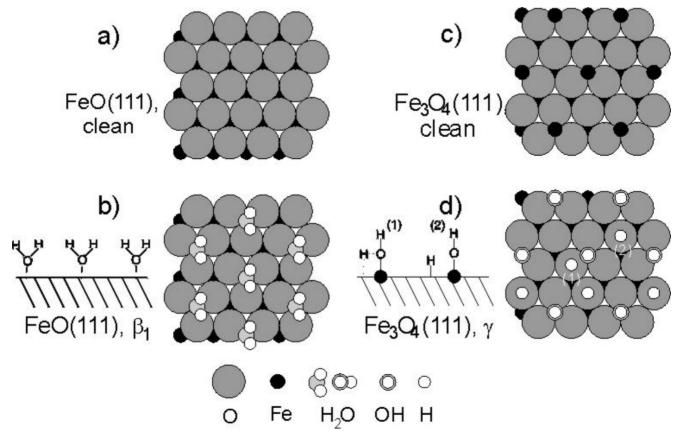
Different adsorption sites on a close-packed surface: Linear or on-top site, bridgebound site and triply coordinated or threefold hollow site (Masel fig. 3.13, p.123).

Example: CO/Pt(111)

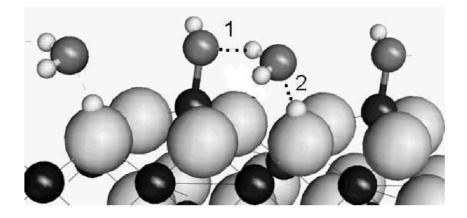


Binding sites for CO on Pt(111): Two ordered adsorption arrangements assumed sequentially with increasing coverage:  $\Theta_r = 0.33$  and 0.5 (Masel fig. 3.14, p.123).

Example:  $H_2O/Fe_3O_4(111)$ ,



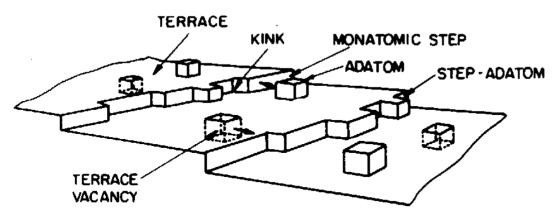
Clean and water covered surfaces of FeO(111) and  $Fe_3O_4(111)$ . Coverages and kind of adsorbed species deduced from UPS measurements. FeO(111) is O-terminated. The interaction is weak but still determined by the position of Fe in the second layer.  $Fe_3O_4$  is terminated by  $\frac{1}{4}$  ML of Fe. Water dissociates. The coverage is compatible with OH being adsorbed on the Fe sites and H binding to O-sites. Site (1) in (d) is the proposed transition state. (Y. Joseph et al. J. Phys. Chem. B 104 (2002) 3224;).



*Further water bridges OH and H via H-bonds. A H*<sub>3</sub>*O-OH like "dimer" is formed.* (*M.E. Grillo et al., Phys. Rev. B 77 (2008) 075407*)

#### Extrinsic inhomogeneity from defects

(steps, domain boundaries, kinks, adatoms, vacancies, contaminants...)



Model of a heterogeneous surface with different surface sites. They are distinguished by their number of nearest neighbours. (Somorjai, fig. 2.6, p. 41.)

 $\Delta H_{ad} = f(\text{substrate}, \Theta).$ 

- material dependence
- site dependence
- ads.-ads.-interaction: repulsive or attractive
- ads.-subs.-interaction may induce structural changes in substrate (reconstruction, relaxation)

**Remember these terms:** 

adsorption, desorption, adsorbent, adsorbate, substrate,

accomodation, sticking, sticking coefficient, monolayer, multilayer;

physisorption, chemisorption, hardness, hard/soft acid/base;

potential curve;

adsorption: molecular, dissociative, non-activated, activated reversible and irreversible;

> adsorption sites: linear, on-top, bridge, triply coordinated, threefold hollow site;

#### 4. Adsorption-Desorption-Equilibrium: Isotherms and isobars

Adsorption rate:

$$r_a = r_a(p, s_0, n, T, E_a, \Theta).$$

Desorption rate:

 $r_d = r_d(v, n, T, E_d, \Theta),$ increases with  $\Theta$ .

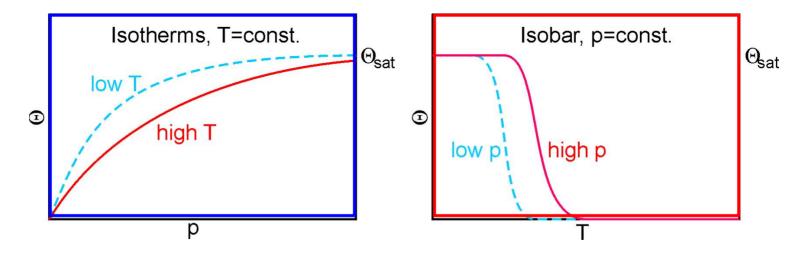
After a while, when

 $(s_0:$  initial sticking coefficient, n: reaction order for adsorption ;  $E_a:$  activation energy for adsorption,  $\Theta:$  coverage)

( $\nu$ : frequency- or prefactor for desorption; *n*: reaction order for desorption;  $E_d$ : activation energy for desorption)

 $r_a = -r_d$ , adsorption-desorption equilibrium condition an equilibrium coverage  $\Theta_{eq}$  will establish.

> $\Theta_{eq}(T)_p$  *T* variable, *p* constant: **Isobar**  $\Theta_{eq}(p)_T$  *p* variable, *T* constant: **Isotherm**



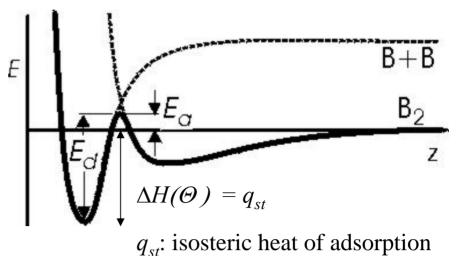
#### How to determine $\Theta$ ?

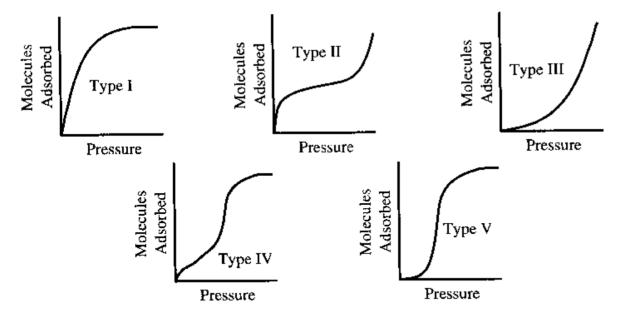
High surface area (powders, pellets...): pressure decrease by adsorption (e.g. BET), TDS.Low surface area (single crystals): $\Delta \phi$ , LEED, IRAS, UPS, TDS.

The difference and  $E_d - E_a$ determines the adsorption energetics

*p*,  $s_0$ , *n* and *v* determine the kinetics.

Aim: measure  $\Theta(p,T)$ , deduce energetic and kinetic parameters from it.





The fife types of adsorption isotherms described by Brunauer[1945]

(Masel fig. 4.4, p.238).

Туре	Where to be found
Ι	monolayer adsorption
II	multilayer adsorption, BET isotherm
III	island or droplet nucleation necessary for adsorption
IV	pore filling, followed by outer-surface adsorption
V	pore filling with nucleation (like III),
	followed by outer surface adsorption

#### 5. A-D energetics: Isosteric heat of adsorption

Measure a set of isotherms or isobars, deduce *p*-*T*-pairs for  $\Theta_{eq}$ =const, plot ln(*p*) vs(1/*T*), use Clausius-Clapeyron equation

> $d(\ln p)/d(1/T) = -\Delta H/R$  or  $\ln p = -\Delta H/RT + \text{const.}$

or, since  $\Delta H$  may be  $\Theta$ -dependent,  $\Delta H(\Theta) = q_{st}$  = isosteric heat of adsorption; isosteric means: for constant coverage.

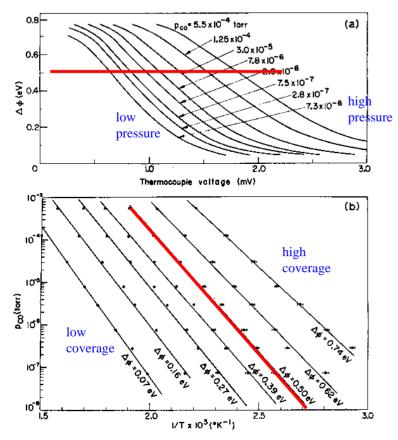
#### Isosteric methods

Advantage:

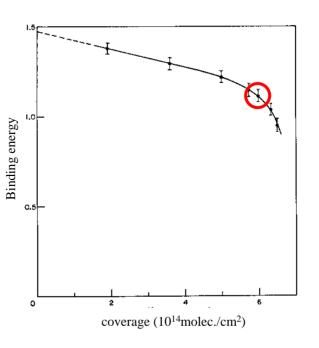
CC is deduced from equilibrium thermodynamics and independent of kinetics. Disadvantage:

adsorption must be reversible. (If not: => calorimetric methods). adsorption and desorption must be sufficiently fast (equilibrium established within reasonable time).

Strictly, CC is only applicable if adsorbate is separable from substrate (no strong interaction, physisorption). Kinetic considerations: Applicable if  $r=C \exp(-E_{acl}/RT)$ (Arrhenius law). (W. Weiss, W. Ranke, Prog. Surf. Sci. 70 (2002) 1). Examples:



a) Isobars for CO adsorption on Pd(111).
As a measure for the coverage, the work function change Δφ was used. In special cases, Δφ is proportional to Θ.
b) Arrhenius-plots ln p vs. 1/T for different coverages (in terms of Δφ).



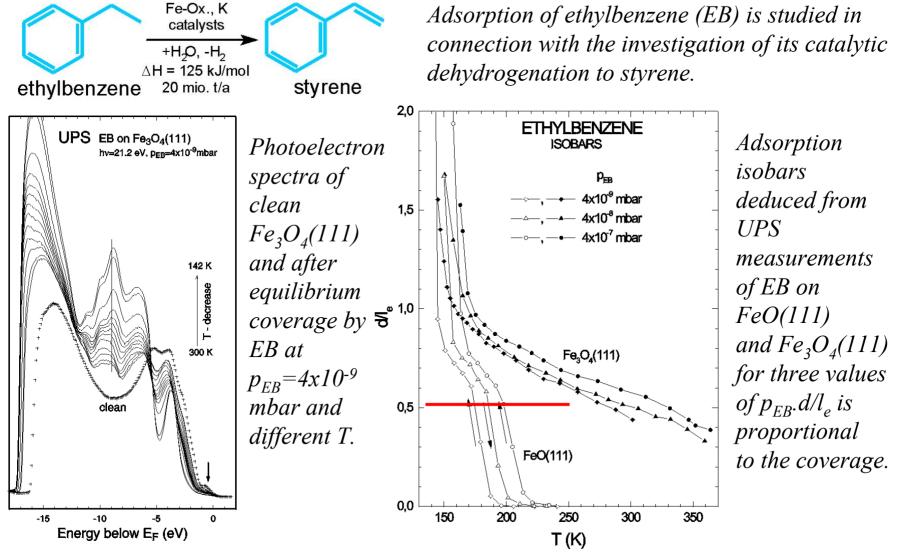
Coverage dependence of  $q_{st}$  deduced from the slopes of the curves in (b). (J.C. Tracy, P.W. Palmberg, Surf. Sci. 14 (1969) 274).

#### 2. Adsorption on a semiconductor: NH<sub>3</sub>/Ge

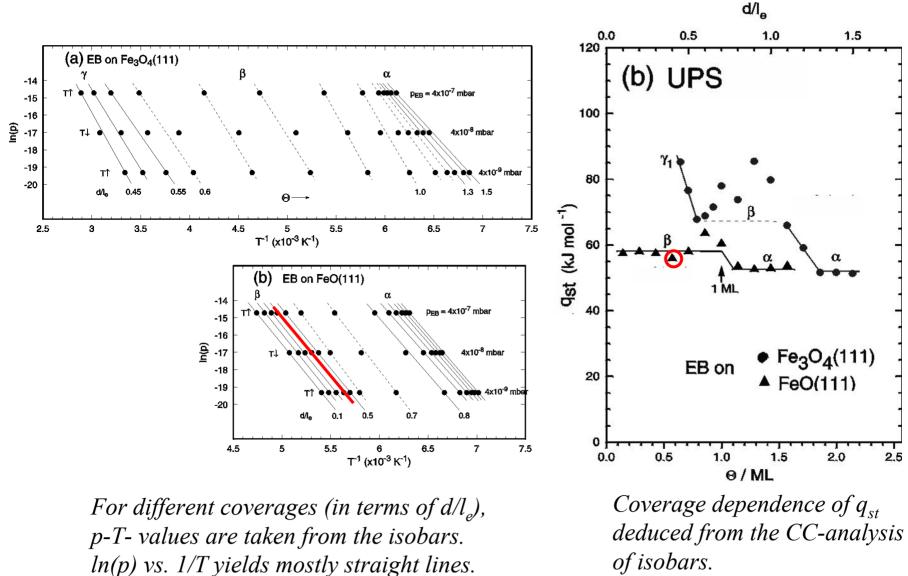
(W. Ranke, Surf. Sci. 342 (1995) 281; W. Ranke, J. Platen, Phys. Rev. B 54 (1996) 2873)

#### 3. Adsorption on an ionic material: Ethylbenzene/FeO(111), $Fe_3O_4(111)$

(W. Weiss, W. Ranke, Progr. Surf. Sci. 70(2002) 1, Figs. 54, 56, 57, 58b)



Intensity (rel. units)



*Their slope yields*  $q_{sr}$ 

==> TDS (TPD), a transient method, evaluation depends on kinetics.

#### 6. **A-D kinetics: The shape of isotherms/isobars**

Kinetics: rate of adsorption and desorption

$$r_a = \frac{\sigma_A \, d\Theta_r}{dt} = \frac{p}{\sqrt{2\pi \, mkT}} s_0 \, \exp(-E_a/kT) \, f(\Theta_r) = A \, f(\Theta_r)$$

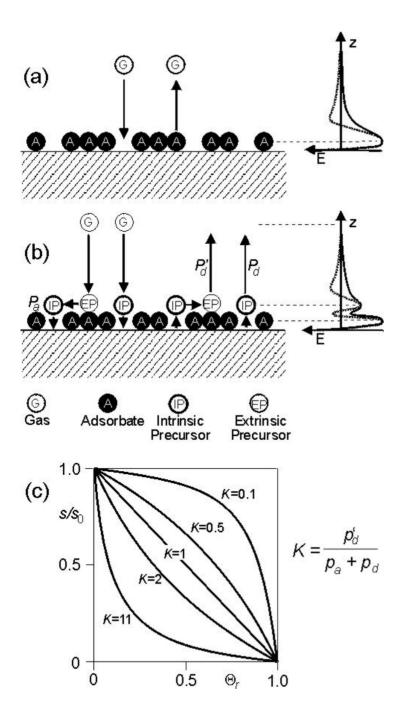
$$r_d = \frac{\sigma_A \, d\Theta_r}{dt} = -v_n \, \sigma_A^n \exp\left(-E_d / kT\right) \Theta_r^n \qquad = -B \, \Theta_r^n$$

 $E_{a}, E_{d}$ : activation energies for adsorption and desorption;  $s_0$ : initial sticking probability:  $\sigma_{4}$ : density of adsorption sites;  $\Theta_r = \Theta / \Theta_{sat}$ : relative coverage ( $0 < \Theta_r < 1$ );  $v_n$ : the frequency factor for reaction order

*n*: *n*: reaction order.

I follow the convention that  $r_d$  is negative since  $d\Theta/dt$  is negative for desorption.

 $f(\Theta_r)$  describes the dependence of the sticking probability s on the coverage,  $s = s_0 f(\Theta_r).$ 



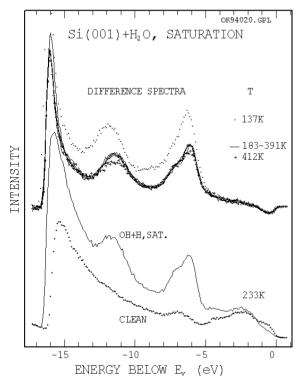
#### Langmuir:

Sticking only if unoccupied site is hit. 1<sup>st</sup> order (molecular):  $f(\Theta_r) = 1 - \Theta_r$ . 2<sup>nd</sup> order (dissociative):  $f(\Theta_r) = (1 - \Theta_r)^2$ .

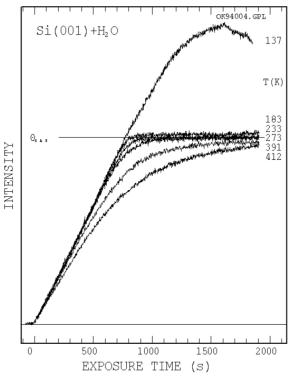
Mobile precursor kinetics: (G) may hit an empty site: *intrinsic precursor* state (IP) which may be mobile. It may transform into the adsorbate state (A) (probability  $P_a$ ) or desorb (probability  $p_d$ ). (G) may hit an occupied site: *extrinsic precursor* (EP) which may migrate to an empty site and transform into an (IP) and adsorb (see above) or desorb directly (probability  $p_d$ ').

**Kisliuk:** sticking behavior is determined *K* which depends on these probabilities only (P. Kisliuk, J. Phys, Chem. Solids, 3 (1957) 95). If  $p_d$ ' is small (stable (EP)), the sticking coefficient may be quite high almost until saturation ( $\Theta_r = 1$ ). 1st order:  $f(\Theta_r) = \frac{1 - \Theta_r}{1 + \Theta_r(K - 1)}$  with K>0. K=1 corresponds to Langmuir case. (W. Ranke, Y. Joseph, Phys. Chem. Chem. Phys. 4 (2002) 2483)

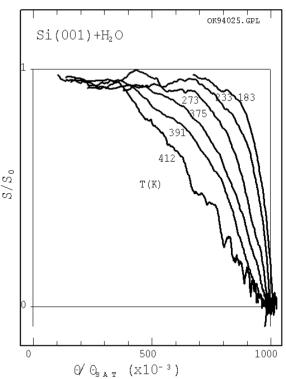
#### Check of kinetics: determine $s(\Theta)$ . Example: H<sub>2</sub>O/Si(001):



UP-spectrum (hv=21.2 eV) of Si(001), clean and after saturation with  $H_2O$  which adsorbs as OH+H. Equal species and saturation for  $183 \le T \le 391$  K. At 137 K, condensation of  $H_2O$  occurs, at 412 K, saturation is not yet reached.



Adsorption curves. Intensity increase of the -6.2 eV peak of OH with exposure time for  $p_{H2O}=2x10^{-9}$ mbar for the same values of T.



The first derivative of the adsorption curves represents the sticking coefficient  $s(\Theta)$ . The wiggles originate from pressure adjustments in the adsorption curves.

(W. Ranke, Surf. Sci. 369 (1996) 137)

In **adsorption-desorption equilibrium** we have  $r_a + r_d = 0$ . We had  $r_a = A(p,T) f(\Theta_r)$ ,  $r_d = -B(p,T) \Theta_r^n$ .

Inserting and resolving for  $\Theta_r$  yields the equations for isotherms/isobars:

Langmuir 1<sup>st</sup> order:  $\Theta_r(T,p) = \frac{b_1(T) p}{1 + b_1(T)}$ 

Langmuir 2<sup>nd</sup> order:  $\Theta_r(T,p) = \frac{\sqrt{b_2(T)} p}{1 + \sqrt{b_2(T)} p}$ 

with n = 1,  $f(\Theta_r) = 1 - \Theta_r$ ,

with 
$$n = 2$$
,  $f(\Theta_r) = (1 - \Theta_r)^2$ .

Holds if the adsorption sites for the dissociation products are equivalent and for inequivalent sites as long as their site concentrations are equal.

Langmuir 2<sup>nd</sup> order adsorption, quasi-1<sup>st</sup> order desorption (dissociated species immobile):

$$\Theta_r(T,p) = \frac{1+2b_1(T)p - \sqrt{1+4b_1(T)p}}{2b_1(T)p} \quad \text{with } n = 1, \ f(\Theta_r) = (1-\Theta_r)^2.$$

Kisliuk 1<sup>st</sup> order:

$$\Theta_r(T,p) = \frac{1 + b_1(T) \ p - \sqrt{[1 - b_1(T) \ p]^2 + 4 \ K \ b_1(T) \ p}}{2 \ (1 - K)} \quad \text{with} f(G)$$

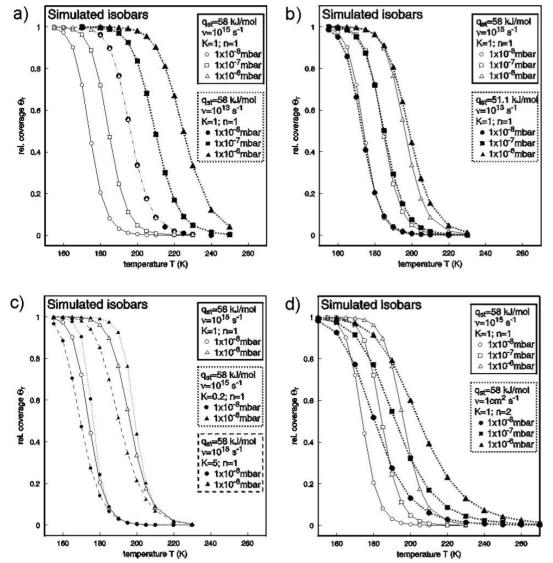
$$b_1(T) = \frac{s_0 \exp(q_{st} / kT)}{v_1 \sigma_A \sqrt{2\pi \ mkT}}$$

$$b_2(T) = \frac{s_0 \exp(q_{st} / kT)}{\nu_2 \sigma_A^2 \sqrt{2\pi \ mkT}}$$

with 
$$n = 1$$
,  
 $f(\Theta_r) = \frac{1 - \Theta_r}{1 + \Theta_r(K - 1)}$ 

$$q_{st} = E_d - E_a$$

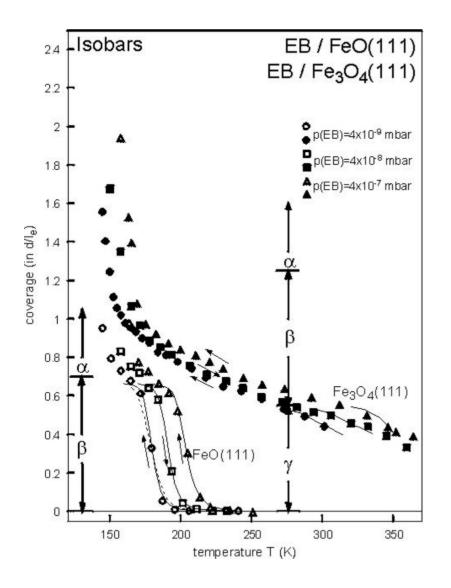
#### What should these isobars look like?



Assumption: Langmuir  $1^{st}$  order (n = 1),  $v_1 = 10^{15} \, s^{-1}$ ,  $q_{st} = 58 \text{ kJ/mol}$ (both independent of  $\Theta$ , data for ethylbenzene physisorbed on FeO(111)). Pressures as indicated a) comparison:  $v_1$  varied ( $\rightarrow 10^{13} \text{ s}^{-1}$ ); b) comparison:  $v_1 \rightarrow 10^{13} \text{ s}^{-1}$ , *q<sub>st</sub>* adjusted to yield best-possible agreement; c) comparison: Kisliuk kinetics (mobile precursor), K varied; d) comparison: Langmuir 2<sup>nd</sup> order.

(W. Ranke, Y. Joseph, Phys. Chem. Chem. Phys. 4 (2002) 2483).

#### Example: Ethylbenzene / FeO(111), Fe<sub>3</sub>O<sub>4</sub>(111)



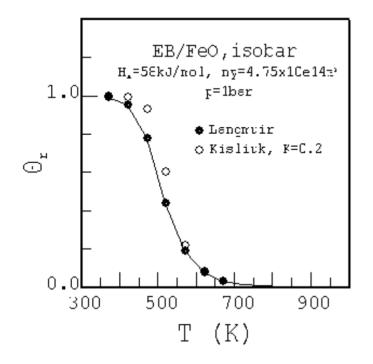
Same measurement as shown above, lines: fitted isobars.

FeO(111),  $\beta$ -range:  $q_{st} = 58 \text{ kJ/mol}$ ,  $v_1 = 4.75 \times 10^{14} \text{ s}^{-1}$ . Dotted: Langmuir 1<sup>st</sup> order gives too smooth curves. Lines: Kisliuk 1<sup>st</sup> order, K = 0.2.

Fe<sub>3</sub>O<sub>4</sub>(111),  $\gamma$ -range:  $q_{st} = 94...74$  kJ/mol,  $v_1 = 5x10^{12}...2x10^{10}$  s<sup>-1</sup> ( $\Theta$ -dependent); Langmuir or Kisliuk possible.  $\beta$ -range: Fit impossible reason: phase transitions in adlayer.

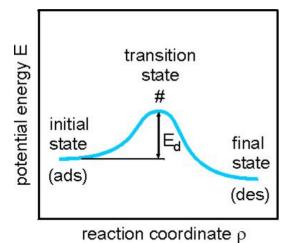
(W. Ranke, Y. Joseph, Phys. Chem. Chem. Phys. 4 (2002) 2483).

#### What should these isobars look like?

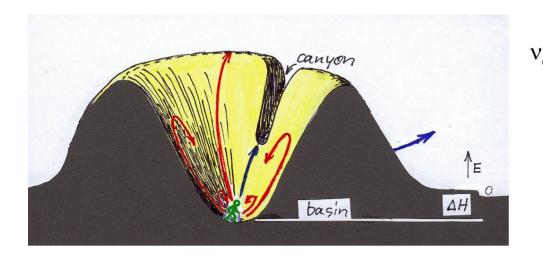


Isobars for EB/FeO(111) for high pressure conditions, extrapolated from data measured at low p, T: Although  $\Delta H_a = q_{st}$  is small,  $\Theta$  gets quite high if only p is high enough.

#### 7. The magnitude of *v*: Entropy of the activated state



The transition state theory starts from the assumption that the molecule is in thermodynamic equilibrium with its environment, even during the desorption reaction.



 $v_{des} \sim probability to realize$ the transition state dominating: Boltzmann factor  $exp(-E_d/kT)$ ; determining: configuration (here restricted direction, rotation).

probability to realize a state: **partition function** 

### It can be shown that (for 1<sup>st</sup> order desorption)

$$\gamma_{des} \approx \frac{kT}{h} \frac{q^{\#}}{q_{ad}} \approx \frac{kT}{h} \exp\left(\frac{\Delta S^{\#}}{R}\right)$$

partition functions  $q^{\#}$  in the transition state  $q_{ad}$  in the adsorbed state

(see e.g. K. Christmann, Surface physical chemistry, Steinkopff, Darmstadt (1991), p.27ff).

If 
$$q^{\#} = q_{ad}$$
,  $v_{des} \approx 10^{13} \text{ s}^{-1}$  (T=300K).

If the adsorbed state is immobile and the transition state is mobile and may even rotate,  $q^{\#}$  and consequently  $v_{des}$  may get much larger.

If the transition state requires a complicated or "demanding" configuration which has a low probability of realization  $q^{\#}$  and  $v_{des}$  may get much smaller.

The agreement between measured and calculated values of  $v_{des}$  is poor. Nevertheless, transition state theory gives an idea why  $v_{des}$  values vary so strongly.

Experimentally,  $v_{des}$  values between about  $10^{10}$  s<sup>-1</sup> and  $10^{19}$  s<sup>-1</sup> have been observed while values around  $10^{13}$  s<sup>-1</sup> do not appear to be especially probable.

Excursion: Partition function, probability and entropy			
Partition function of a system: $q = \Sigma_i \exp(-\varepsilon_i / kT)$ $=$ sum over all possible states, weighted by $\exp(-\varepsilon_i / kT)$ , the prob. of occupation (Boltzmann)ornumber of thermally accessible states at given T;orprobability to ,,materialize" in the suggested arrangement or configuration $(\rightarrow \text{entropy}, S=k \ln W).$			
Quantum mechanically: Discrete states, different contribution to q at 300 K: Translation: densely spaced, $\Delta \varepsilon \approx 10^{-16}$ kT dominating in gas phase, zero in immob. ads. Rotation: less densely spaced, $\Delta \varepsilon \approx 0.1$ kT small in gas, usually zero in ads. Vibration: not densely spaced, $\Delta \varepsilon \approx 10$ kT very small in gas, dominating in ads. Electronic: separated $\Delta \varepsilon \approx 40$ kT virtually zeroAdsorbed state: q much smaller than in gas $\rightarrow$ compensation by $E_{ad}$ necessary			

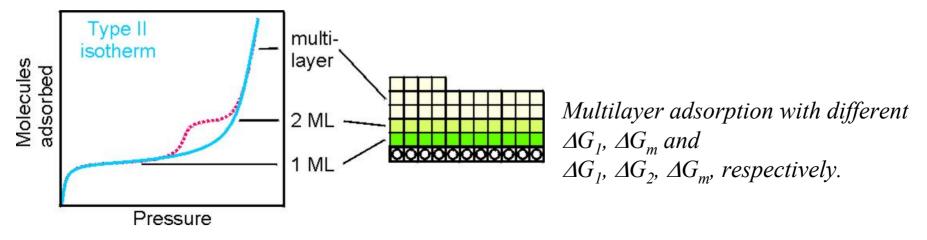
#### 8. Further isotherms for monolayer and multilayer adsorption

8.1 Other isotherms / isobars for monolayer adsorption:

Isotherm	Advantages	Disadvantages
Langmuir	Best one-parameter isotherm	Ignores adsorbate-adsorbate interactions
Freundlich, Toth	Two parameters	No physical basis for equation
Multisite	Many parameters	Good for inhomogeneous surfaces. Wrong physics for single crystals
Tempkin Fowler Slygin-Frumkin	Account for adsorbate-adsorbate interactions in an average sense	Does not consider how the adsorbate layer is arranged
Lattice gas	Complete description of adsorbate-adsorbate interactions for commensurate layers Predicts arrangement of adsorbed layer	Requires a computer to calculate isotherm Assumes commensurate adsorption Parameters used in the model are difficult to determine

(R.I. Masel, Principles of adsorption and reaction on solid surfaces, Wiley, New York (1996), table 4.1, p.246).

#### 8.2 Multilayer adsorption: The BET isotherm (Brunauer-Emmett-Teller)



#### Either:

Consider each step in the isotherm separately with its own *n*, *v*,  $s_0$ ,  $q_{st}$ , add isotherms (e.g. Langmuir) for 1<sup>st</sup>, 2<sup>nd</sup> ... layer. Problem: Multilayer condensation. Or:

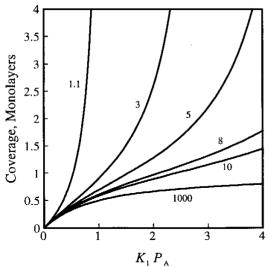
Develop an isotherm for multilayer adsorption: Brunauer-Emmett-Teller, BET-isotherm.

$$\frac{V}{V_M} = \frac{c_B x_B}{(1 - x_B) \left[1 + (c_B - 1) x_B\right]}$$

(V: total volume of adsorbed gas;  $V_M$ : volume gas in 1<sup>st</sup> layer (monolayer saturation));

 $x_B = p/p_{sat}$ (*p* the gas pressure and  $p_{sat}$  the saturation vapor pressure of the gas at the temperature of the measurement).  $c_B = K_l/K_m$  ratio of equilibrium constants for 1<sup>st</sup> layer and multilayer adsorption:

$$c_B = \frac{K_1}{K_m} = \frac{\exp(-\Delta G_1 / RT)}{\exp(-\Delta G_m / RT)} = \exp[-(\Delta G_1 - \Delta G_m) / RT]$$



A plot of adsorption isotherms predicted by the BET equation for various values of  $c_B$ . (Masel fig. 4.30, p.302)

In fact, the BET isotherm does in general not very well fit measured isotherms. However, in the environment of ML saturation it usually does.

If, however, ML.saturation can be unambiguously be identified in the isotherm by step structure, the ML capacity can directly be taken from the isotherm.





## **Remember these terms:**

adsorption-desorption equilibrium, isotherm, isobar, isosteric heat of adsorption; rates of adsorption and desorption, order of adsorption and desorption., frequency- or pre-factor; Langmuir adsorption, mobile precursor; transition state; **BET** isotherm.